

Group Classification and Exact Solutions of Nonlinear Wave Equations

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Abstract We perform complete group classification of the general class of quasi linear wave equations in two variables. This class may be seen as a broad generalization of the nonlinear d'Alembert, Liouville, sin/sinh-Gordon and Tzitzeica equations. In this way we derived a number of new genuinely nonlinear invariant models with high symmetry properties. In particular, we obtain four classes of nonlinear wave equations admitting five-dimensional invariance groups. Applying the symmetry reduction technique we construct multi-parameter families of exact solutions of these equations.

1. Introduction

It was Sophus Lie who was the first to utilize group properties of differential equations in order to actually solve them (to construct their exact solutions). He computed the maximal invariance group of the one-dimensional heat conductivity equation and utilized this symmetry to construct its explicit solutions. Saying it the modern way, he performed symmetry reduction of the heat equation. Nowadays symmetry reduction is one of the most powerful tools for solving nonlinear partial differential equations (PDEs).

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The majority of fundamental equations of mathematical and theoretical physics admit wide symmetry groups. It is the rich symmetry that enables developing a variety of efficient methods for mathematical analysis of linear differential equations. However, linear equations give mathematical description of physical, chemical or biological processes in a first approximation only. To provide a more detailed and precise description a mathematical model has to incorporate nonlinear terms. What is more, some important mathematical physics equations have no linearized versions at all. A well-known example is the system of Yang–Mills equations.

Hyperbolic type second-order nonlinear PDEs in two independent variables are utilized to describe different types of wave propagation. They are also used in differential geometry, in various fields of hydro- and gas dynamics, chemical technology, super conductivity, crystal dislocation to mention only a few applications areas. However the list of models (equations) is surprisingly narrow. By the most part it is comprised by the Liouville, sine/sinh-Gordon, Goursat, d'Alembert and Tzitzeica equations and a couple of others. From the group-theoretical viewpoint the popularity of these very models is due to the fact that they have non-trivial Lie or Lie–Bäcklund symmetry. By this very reason they are either integrable by the inverse problem methods (see, e.g., [1–3]) or are linearizable [4–6] and completely integrable [7, 8].

In this connection it seems to be a very important problem to select from the reasonably extensive class of nonlinear hyperbolic type PDEs those enjoying the best symmetry properties. Saying 'reasonably extensive' we mean that this class should contain the above enumerated equations as particular cases, on the one hand, and it should contain a wide variety of new invariant models of potential interest for applications, on the other. The list of the so obtained invariant equations will contain candidates for realistic nonlinear mathematical models of the physical and chemical processes enumerated above.

The history of group classification methods goes back to Lie. The first paper on this subject is [9], where Lie proves that a linear two-dimensional second-order PDE may admit at most a three-parameter invariance group (apart from the trivial infinite-parameter symmetry group, which is due to linearity).

The modern formulation of the problem of group classification of PDEs was suggested by Ovsyannikov in [10]. He developed the regular method (we will refer to it as to the Lie–Ovsyannikov method) for classifying differential equations with non-trivial symmetry and performed complete group classification of a class of nonlinear heat conductivity equations. In a number of subsequent publications more general types of nonlinear heat equations were classified (a review of these results can be found in [11]).

Analysis of the papers on group classification of PDEs reveals that the majority of them deal with equations whose arbitrary elements (functions) depend on one variable only. The reason for this fact is that application of Lie–Ovsyannikov method to PDEs, which contain arbitrary functions of several variables, is not always possible. To achieve a complete classification one either needs to specify the transformation group realization or somehow restrict arbitrariness of functions contained in the equation under study.

Recently, we developed the efficient approach enabling to overcome this difficulty for low dimensional PDEs [12, 13]. Utilizing it we have obtained the ultimate solution of the problem of group classification of the general quasi-linear heat conductivity equations in two independent variables.

In this paper we apply the approach mentioned above to perform group classification of the most general quasi-linear hyperbolic type PDE in two independent variables.

2. Group Classification Algorithm

We begin this section by formulating the problem to be solved. Then we briefly review the already known results. Finally we give a short description of our approach to group classification of PDEs (for the detailed account of the necessary facts, see ([13])).

While classifying a given class of differential equations into subclasses, one can use different classifying features, like linearity, order, the number of independent or dependent variables, etc. In group analysis of differential equations the principal classifying features are symmetry properties of equations under study. This means that classification objects are equations together with their symmetry groups. This point of view is based on the well-known fact that any PDE admits a (possibly trivial) Lie transformation group. And what is more, any transformation group corresponds to a class of PDEs, which are invariant under this group. So that performing group classification of a class of PDEs means describing all possible (inequivalent) pairs (PDE, maximal invariance group), where PDE should belong to the class of equations under consideration.

We perform group classification of the following class of quasi-linear wave equations:

$$u_{tt} = u_{xx} + F(t, x, u, u_x). \quad (1.1)$$

Here F is an arbitrary smooth function, $u = u(t, x)$. Hereafter we adopt notations $u_t = \frac{\partial u}{\partial t}$, $u_x = \frac{\partial u}{\partial x}$, $u_{tt} = \frac{\partial^2 u}{\partial t^2}$, ...

Our aim is describing all equations of the form (1.1) that admit non-trivial symmetry groups. The challenge of this task is in the word *all*. If, for example, we fix the form of a desired invariance group, then the classification problem simplifies immensely. A slightly more cumbersome (but still tractable with the standard Lie–Ovsyannikov approach) is the problem of group classification of equation with arbitrary functions of at most one variable.

As equations invariant under similar Lie groups are identical within the group-theoretic framework, it makes sense to consider non-similar transformation groups [14, 15] only. The important example of similar Lie groups is provided by Lie transformation groups obtained one from another by a suitable change of variables. Consequently, equations obtained one from another by a change of variables have similar symmetry groups and cannot be distinguished within the group-theoretical viewpoint. That is why, we perform group classification of (1.1) within a (locally invertible) change of variables preserving the class of PDEs (1.1).

The problem of group classification of linear hyperbolic type equation

$$u_{tx} + A(t, x)u_t + B(t, x)u_x + C(t, x)u = 0 \quad (1.2)$$

with $u = u(t, x)$, was solved by Lie [9] (see, also, [16]). In view of this fact, we consider only those equations of the form (1.1) which are not (locally) equivalent to the linear equation (1.2).

As we have already mentioned in Introduction, the Lie–Ovsyannikov method of group classification of differential equations has been suggested in [10]. Utilizing this method enabled solving group classification problem for a number of important one-dimensional nonlinear wave equations:

$$\begin{aligned}
 u_{tt} &= u_{xx} + F(u); & [17] - [19] \\
 u_{tt} &= [f(u)u_x]_x; & [20] - [22] \\
 u_{tt} &= f(u_x)u_{xx}; & [22, 23] \\
 u_{tt} &= F(u_x)u_{xx} + H(u_x); & [24] \\
 u_{tt} &= F(u_{xx}); & [22] \\
 u_{tt} &= u_x^m u_{xx} + f(u); & [25] \\
 u_{tt} + f(u)u_t &= (g(u)u_x)_x + h(u)u_x; & [26] \\
 u_{tt} &= (f(x, u)u_x)_x. & [27]
 \end{aligned}$$

Analysis of the above list shows that the most of arbitrary elements (= arbitrary functions) depend on one variable. This is not coincidental, indeed, the Lie–Ovsyannikov approach works most efficiently for the case when arbitrary elements are functions of one variable only. The reason for this fact is that the obtained system of determining equations is still over-determined. So that it can be effectively solved using the same technique that is applied to compute maximal symmetry group of PDEs containing no arbitrary elements.

The matter becomes much more complicated for the case when arbitrary elements are functions of two (or more) arguments. By this very reason group classifications of nonlinear wave equations

$$\begin{aligned}
 u_{tt} + \lambda u_{xx} &= g(u, u_x); & [28, 29] \\
 u_{tt} &= [f(u)u_x + g(x, u)]_x; & [30] \\
 u_{tt} &= f(x, u_x)u_{xx} + g(x, u_x) & [31]
 \end{aligned}$$

are not complete.

We suggest the new approach to problem of group classification of low dimensional PDEs in [12, 13]. This approach is based on Lie–Ovsyannikov infinitesimal method and classification results for abstract finite-dimensional Lie algebras. It enabled obtaining the complete solution of group classification problem for the general heat equation with a nonlinear source

$$u_t = u_{xx} + F(t, x, u, u_x).$$

Later on, we performed complete group classification of the most general quasi-linear evolution equation [32–34]

$$u_t = f(t, x, u, u_x)u_{xx} + g(t, x, u, u_x).$$

A starting point of our analysis is a well-known fact that solutions $\mathbf{v}_a = (\tau_a, \xi_a, \eta_a)$, $a = 1, \dots, n$ of the determining equations (we denote them as DE) for symmetry operators

$$Q = \tau(t, x, u)\partial_t + \xi(t, x, u)\partial_x + \eta(t, x, u)\partial_u$$

admitted by (1.1) span a Lie algebra ℓ . So without any loss of generality we can replace DE with the (possibly infinite) set of systems of PDEs

$$\begin{cases} \text{DE,} \\ [Q_i, Q_j] = C_{ij}^k Q_k, \end{cases}$$

or, equivalently,

$$\begin{cases} \text{DE,} \\ Q_i \tau_j - Q_j \tau_i = C_{ij}^k \tau_k, \\ Q_i \xi_j - Q_j \xi_i = C_{ij}^k \xi_k, \\ Q_i \eta_j - Q_j \eta_i = C_{ij}^k \eta_k. \end{cases}$$

In the above formulas the indices i, j, k take the values $1, \dots, n$ ($n \geq 1$ is a dimension of the corresponding Lie algebra), C_{ij}^k are structure constants of the Lie algebra ℓ and $Q_a = \tau_a(t, x, u)\partial_t + \xi_a(t, x, u)\partial_x + \eta_a(t, x, u)\partial_u$.

If we solve the (over-determined) system of PDEs given above for all possible dimensions $n \geq 1$ of all admissible Lie algebras, ℓ , then the problem of group classification of Equation (1.1) is completely solved. In other words the problem of group classification of the general wave equation (1.1) reduces to integrating over-determined systems of PDEs in question for all $n = 1, 2, \dots, n_0$, where n_0 is the maximal dimension of the Lie algebra admitted by the equation under study. More formally, our algorithm for group classification of the class of PDEs (1.1) consists of the following steps (for more details, see [34]):

- I. Using the infinitesimal Lie method we derive the system of determining equations for coefficients of the first-order operator that generates symmetry group of Equation (1.1) (Note that the determining equations which explicitly depend on the function F and its derivatives are called classifying equations). Integrating equations that do not depend on F we obtain the form of the most general infinitesimal operator admitted by Equation (1.1) under arbitrary F . Another task of this step is calculating the equivalence group \mathcal{E} of the class of PDEs (1.1).
- II. We construct all realizations of Lie algebras A_n of the dimension $n \leq 3$ in the class of operators obtained at the first step within the equivalence relation defined by transformations from the equivalence group \mathcal{E} . Inserting the so obtained operators into classifying equations we select those realizations that can be symmetry algebras of a differential equation of the form (1.1).
- III. We compute extensions of the realizations constructed at the previous step to realizations of higher dimensional ($n > 3$) Lie algebras. Since extending symmetry algebras results in reducing arbitrariness of the function F , at some point this function will contain either arbitrary functions of one variable or arbitrary constants. At this point, we apply Lie–Ovsyannikov classification method to derive the maximal symmetry group of the equation under study thus completing its group classification.

Performing the above enumerated steps yields the complete list of inequivalent equations of the form (1.1) together with their maximal (in Lie's sense) symmetry algebras.

We say that the group classification problem is completely solved if it has been proved that

- 1) The constructed symmetry algebras are maximal invariance algebras of the equations under consideration;
- 2) The list of invariant equations contains only inequivalent ones, namely, no equation can be transformed into another one from the list by a transformation from the equivalence group \mathcal{E} .

3. Preliminary Group Classification of Equation (1.1)

The first step of the algorithm is looking for the infinitesimal operator of symmetry group of Equation (1.1) in the form

$$Q = \tau(t, x, u)\partial_t + \xi(t, x, u)\partial_x + \eta(t, x, u)\partial_u, \tag{2.1}$$

where τ, ξ, η are smooth functions defined on an open domain Ω of the space $V = \mathbb{R}^2 \times \mathbb{R}^1$ of independent $\mathbb{R}^2 = \langle t, x \rangle$ and dependent $\mathbb{R}^1 = \langle u \rangle = u(t, x)$ variables.

Operator (2.1) generates one-parameter invariance group of Equation (1.1) iff its coefficients $\tau, \xi, \eta, \epsilon$ satisfy the equation (Lie’s invariance criterion)

$$\varphi^{tt} - \varphi^{xx} - \tau F_t - \xi F_x - \eta F_u - \varphi^x F_{u_x} \Big|_{(1.1)} = 0, \tag{2.2}$$

where

$$\begin{aligned} \varphi^t &= D_t(\eta) - u_t D_t(\tau) - u_x D_t(\xi), \\ \varphi^x &= D_x(\eta) - u_t D_x(\tau) - u_x D_x(\xi), \\ \varphi^{tt} &= D_t(\varphi^t) - u_{tt} D_t(\tau) - u_{tx} D_t(\xi), \\ \varphi^{xx} &= D_x(\varphi^x) - u_{tx} D_x(\tau) - u_{xx} D_x(\xi) \end{aligned}$$

and D_t, D_x are operators of total differentiation with respect to the variables t, x . As customary, by writing $\Big|_{(1.1)}$ we mean that one needs to replace u_{tt} and its differential consequences with the expression $u_{xx} + F$ and its differential consequences.

After a simple algebra we represent (2.2) in the form of system of four PDEs:

$$\begin{aligned} (1) \quad & \xi_u = \tau_u = \eta_{uu} = 0, \\ (2) \quad & \tau_t - \xi_x = 0, \quad \xi_t - \tau_x = 0, \\ (3) \quad & 2\eta_{tu} + \tau_x F_{u_x} = 0, \\ (4) \quad & \eta_{tt} - \eta_{xx} - 2u_x \eta_{xu} + [\eta_u - 2\tau_t]F - \tau F_t - \xi F_x \\ & - \eta F_u - [\eta_x + u_x(\eta_u - \xi_x)]F_{u_x} = 0. \end{aligned} \tag{2.3}$$

It follows from Equation (1) that $\tau = \tau(t, x), \xi = \xi(t, x), \eta = h(t, x)u + r(t, x)$. In the sequel we differentiate between the cases $F_{u_x u_x} \neq 0$ and $F_{u_x u_x} = 0$.

Case $F_{u_x u_x} \neq 0$. It follows from (3) that $\tau_x = h_t = 0$. Taking into account these equations and also Equation (2), we obtain $\tau = \lambda t + \lambda_1, \xi = \lambda x + \lambda_2, h = h(x)$, where $\lambda, \lambda_1, \lambda_2$ are arbitrary real constants.

Case $F_{u_x u_x} = 0$. If this is the case, then $F = g(t, x, u)u_x + f(t, x, u)$, where f and g are arbitrary smooth functions.

Given the condition $g_u \neq 0$, it follows from (3) that $\tau_x = h_t = 0$. So that taking into account Equation (2) we arrive at the already known expressions for τ, ξ, h .

If $g_u = 0$, then $f_{uu} \neq 0$, since otherwise (1.1) becomes linear.

Let $g \equiv 0$. It follows from (3) that $\eta = h(x)u + r(t, x)$. Equation (4) now reads as

$$r_{tt} - r_{xx} - 2h'u_x + [h - 2\tau_t]f - \tau f_t - \xi f_x - [hu + r]f_u = 0.$$

As functions τ, ξ, h, r, f do not depend on u_x , we have $h' = 0$. Hence $\eta = mu + r(t, x)$, where m is an arbitrary real constant. Furthermore, if $g = g(t, x) \neq 0$, then it is straightforward to verify that system of Equations (3), (4) is equivalent to the following equations:

$$2h_t = -\tau_x g, \quad 2h_x = -\tau_t g - \tau g_t - \xi g_x,$$

$$(h_{tt} - h_{xx})u + r_{tt} - r_{xx} + f[h - 2\tau_t] - \tau f_t - \xi f_x - [hu + r]f_u - (h_x u + r_x)g = 0.$$

Integrating Equation (2) yields $\tau = \varphi(\theta) + \psi(\sigma)$, $\xi = -\varphi(\theta) + \psi(\sigma)$, where φ, ψ are arbitrary smooth functions of $\theta = t - x, \sigma = t + x$. So we prove the following assertion.

THEOREM 1. *Provided $F_{u_x u_x} \neq 0$, the maximal invariance group of Equation (1.1) is generated by the following infinitesimal operator:*

$$Q = (\lambda t + \lambda_1)\partial_t + (\lambda x + \lambda_2)\partial_x + [h(x)u + r(t, x)]\partial_u, \quad (2.4)$$

where $\lambda, \lambda_1, \lambda_2$ are real constants and $h = h(x), r = r(t, x), F = F(t, x, u, u_x)$ are functions obeying the constraint

$$\begin{aligned} r_{tt} - r_{xx} - \frac{d^2 h}{dx^2} u - 2 \frac{dh}{dx} u_x + (h - 2\lambda) F \\ - (\lambda t + \lambda_1) F_t - (\lambda x + \lambda_2) F_x - (hu + r) F_u \\ - (r_x + \frac{dh}{dx} u + (h - \lambda)u_x) F_{u_x} = 0. \end{aligned} \quad (2.5)$$

If $F = g(t, x, u)u_x + f(t, x, u), g_u \neq 0$, then the maximal invariance group of Equation (1.1) is generated by infinitesimal operator (2.4), where $\lambda, \lambda_1, \lambda_2$ are real constants, and h, r, g, f are functions satisfying the system of two equations

$$\begin{aligned} -2h' - \lambda g &= (\lambda t + \lambda_1)g_t + (\lambda x + \lambda_2)g_x + (hu + r)g_u, \\ -h'' u + r_{tt} - r_{xx} + (h - 2\lambda)f &= (\lambda t + \lambda_1)f_t + (\lambda x + \lambda_2)f_x \\ &+ (hu + r)f_u + g(h'u + r_x). \end{aligned} \quad (2.6)$$

Next, if $F = g(t, x)u_x + f(t, x, u), q \neq 0, f_{uu} \neq 0$, then the infinitesimal operator of invariance group of Equation (1.1) reads as

$$Q = \tau(t, x)\partial_t + \xi(t, x)\partial_x + (h(t, x)u + r(t, x))\partial_u,$$

where τ, ξ, h, r, g, f are functions satisfying the system of PDEs

$$\begin{aligned} \tau_t - \xi_x &= 0, \quad \xi_t - \tau_x = 0, \\ 2h_t &= -\tau_x g, \quad 2h_x = -\tau_t g - \tau g_t - \xi g_x, \\ (h_{tt} - h_{xx})u + r_{tt} - r_{xx} + f(h - 2\tau_t) - \tau f_t \\ &- \xi f_x - (hu + r) f_u - (h_x u + r_x) g = 0. \end{aligned}$$

Finally, if $F = f(t, x, u)$, $f_{uu} \neq 0$, then the maximal invariance group of Equation (1.1) is generated by the infinitesimal operator

$$Q = [\varphi(\theta) + \psi(\sigma)]\partial_t - [\varphi(\theta) - \psi(\sigma)]\partial_x + [ku + r(t, x)]\partial_u,$$

where $k \in \mathbb{R}$, $\theta = t - x$, $\sigma = t + x$ and functions φ, ψ, r, f and constant k satisfy the following equation:

$$\begin{aligned} r_{tt} - r_{xx} + [k - 2\varphi' - 2\psi']f - (\varphi + \psi) f_t + \\ + (\varphi - \psi) f_x - (ku + r) f_u = 0, \quad \varphi' = \frac{d\varphi}{d\theta}, \quad \psi' = \frac{d\psi}{d\theta}. \end{aligned}$$

Summing up the above considerations we conclude that the problem of group classification of Equation (1.1) reduces to the one of classifying equations of the more specific forms

$$u_{tt} = u_{xx} + F(t, x, u, u_x), \quad F_{u_x u_x} \neq 0; \tag{2.7}$$

$$u_{tt} = u_{xx} + g(t, x, u)u_x + f(t, x, u), \quad g_u \neq 0; \tag{2.8}$$

$$u_{tt} = u_{xx} + g(t, x)u_x + f(t, x, u), \quad g \neq 0, \quad f_{uu} \neq 0; \tag{2.9}$$

$$u_{tt} = u_{xx} + f(t, x, u), \quad f_{uu} \neq 0. \tag{2.10}$$

Consider the last two equations. Using the change of variables

$$\bar{t} = t - x, \quad \bar{x} = t + x, \quad u = v(\bar{t}, \bar{x})$$

we reduce them to the equations

$$\begin{aligned} v_{\bar{t}\bar{x}} &= \frac{1}{4} f(\bar{t}, \bar{x}, v), \\ v_{\bar{t}\bar{x}} &= -\frac{1}{4} g(\bar{t}, \bar{x})(v_{\bar{t}} - v_{\bar{x}}) + \frac{1}{4} f(\bar{t}, \bar{x}, v). \end{aligned} \tag{2.11}$$

Now making the change of variables

$$\tilde{t} = \bar{t}, \quad \tilde{x} = \bar{x}, \quad \tilde{v}(\tilde{t}, \tilde{x}) = \Lambda(\bar{t}, \bar{x})v,$$

where $\Lambda = \exp\left[-\frac{1}{4} \int g(\bar{t}, \bar{x})d\bar{x}\right]$, we transform (2.11) to become

$$\tilde{v}_{\tilde{t}\tilde{x}} = \left(\frac{1}{4}g - \Lambda^{-1}\Lambda_{\tilde{t}}\right)\tilde{v}_{\tilde{x}} - \frac{1}{4}g\Lambda^{-1}\Lambda_{\tilde{t}}\tilde{v} + \frac{1}{4}\Lambda^{-1}\Lambda_{\tilde{x}}g\tilde{v} + \Lambda^{-1}f.$$

Hence we conclude that the following assertion holds true.

ASSERTION 1. The problem of group classification of Equations (2.9), (2.9) is equivalent to the one of classifying equations

$$u_{tx} = g(t, x)u_x + f(t, x, u), \quad g_x \neq 0, \quad f_{uu} \neq 0; \quad (2.12)$$

$$u_{tx} = f(t, x, u), \quad f_{uu} \neq 0. \quad (2.13)$$

Note that the constraint $g_x \neq 0$ is essential, since otherwise (2.12) is locally equivalent to (2.13).

Summing up we conclude that the problem of group classification of (1.1) reduces to classifying the more particular classes of PDEs (2.7), (2.8), (2.12), (2.13). In what follows, we provide full calculation details for Equations (2.9) and (2.10) only. The reason is just to save space and still be able to present all details of the algorithm.

First, we consider Equations (2.8), (2.12), (2.13).

4. Group Classification of Equation (2.8)

According to Theorem 1 the invariance group of Equation (2.8) is generated by infinitesimal operator (2.4). And what is more, the real constants $\lambda, \lambda_1, \lambda_2$ and functions h, r, g, f satisfy Equation (2.6). System (2.6) is to be used to specify both the form of nonlinear term (2.8) (i.e., the functions f, g) and the functions h, r and constants $\lambda, \lambda_1, \lambda_2$ in (2.4). It is called the system of classifying equations.

Efficiency of the Lie method for calculation of maximal invariance group of PDE relies essentially on the fact that routinely the corresponding system of determining equations is over-determined. This is clearly not the case, since we have only one equation for four (!) arbitrary functions and three of the latter depend on two variables. By this reason the straightforward application of Lie–Ovsyannikov algorithm is not possible.

Next, we compute the equivalence group \mathcal{E} of Equation (2.8). This group is generated by invertible transformations of the space V preserving the differential structure of Equation (2.8) (see, e.g., [14]). Saying it another way, a group transformation from \mathcal{E}

$$\bar{t} = \alpha(t, x, u), \quad \bar{x} = \beta(t, x, u), \quad v = U(t, x, u), \quad \frac{D(\bar{t}, \bar{x}, v)}{D(t, x, u)} \neq 0,$$

should reduce (2.8) to an equation of the same form

$$v_{\bar{t}\bar{t}} = v_{\bar{x}\bar{x}} + \tilde{g}(\bar{t}, \bar{x}, v)v_{\bar{x}} + \tilde{f}(\bar{t}, \bar{x}, v), \quad \tilde{g}_v \neq 0$$

with possibly different \tilde{f}, \tilde{g} .

Ovsyannikov prove [14] that it is possible to modify the Lie's infinitesimal approach to calculate the equivalence group in the essentially same way as the invariance group. We omit simple intermediate calculations and present the final result.

ASSERTION 2. The maximal equivalence group \mathcal{E} of Equation (2.8) is

$$\bar{t} = kt + k_1, \quad \bar{x} = \epsilon kx + k_2, \quad v = X(x)u + Y(t, x), \quad (3.1)$$

where $k \neq 0, X \neq 0, \epsilon = \pm 1, k, k_1, k_2 \in \mathbb{R}$, and X, Y are arbitrary smooth functions.

This completes the first step of the algorithm.

4.1. Preliminary Group Classification of Equation (2.8)

First, we describe equations of the form (2.8) which admit one-parameter invariance groups.

LEMMA 1. *There exist transformations (3.1) that reduce operator (2.4) to one of the six forms:*

$$\begin{aligned}
 Q &= m(t\partial_t + x\partial_x), \quad m \neq 0; \quad Q = \partial_t + \beta\partial_x, \quad \beta \geq 0; \\
 Q &= \partial_t + \sigma(x)u\partial_u, \quad \sigma \neq 0; \quad Q = \partial_x; \\
 Q &= \sigma(x)u\partial_u, \quad \sigma \neq 0; \quad Q = \theta(t, x)\partial_u, \quad \theta \neq 0.
 \end{aligned}
 \tag{3.2}$$

Proof can be found in [35].

THEOREM 2. *There are exactly five inequivalent equations of the form (2.8) that admit one-parameter transformation groups. They are listed below together with one-dimensional Lie algebras generating their invariance groups (note that we do not present the full form of invariant PDEs and just give the functions f and g)*

$$\begin{aligned}
 A_1^1 &= \langle t\partial_t + x\partial_x \rangle : g = x^{-1}\tilde{g}(\psi, u), \\
 &\quad f = x^{-2}\tilde{f}(\psi, u), \psi = tx^{-1}, \tilde{g}_u \neq 0; \\
 A_1^2 &= \langle \partial_t + \beta\partial_x \rangle : g = \tilde{g}(\eta, u), \quad f = \tilde{f}(\eta, u), \\
 &\quad \eta = x - \beta t, \quad \beta \geq 0, \quad \tilde{g}_u \neq 0; \\
 A_1^3 &= \langle \partial_t + \sigma(x)u\partial_u \rangle : g = -2\sigma'\sigma^{-1} \ln |u| + \tilde{g}(\rho, x), \\
 &\quad f = (\sigma'\sigma^{-1})^2 u \ln^2 |u| - \sigma'\sigma^{-1}\tilde{g}(\rho, x)u \ln |u| - \sigma^{-1}\sigma''u \ln |u| + u\tilde{f}(\rho, x), \\
 &\quad \rho = u \exp(-t\sigma), \quad \sigma \neq 0; \\
 A_1^4 &= \langle \partial_x \rangle : g = \tilde{g}(t, u), \quad f = \tilde{f}(t, u), \quad \tilde{g}_u \neq 0; \\
 A_1^5 &= \langle \sigma(x)u\partial_u \rangle : g = -2\sigma'\sigma^{-1} \ln |u| + \tilde{g}(t, x), \quad f = (\sigma'\sigma^{-1})^2 u \ln^2 |u| \\
 &\quad -(\sigma^{-1}\sigma'' + \sigma^{-1}\sigma'\tilde{g}(t, x))u \ln |u| + u\tilde{f}(t, x), \quad \sigma' \neq 0.
 \end{aligned}$$

Proof can be found in [35].

In a sequel we will need the following technical lemma.

LEMMA 2. *There are no realizations of semi-simple Lie algebras by operators of the form (2.4).*

Proof. To prove the lemma it suffices to check that there are no realizations of the lowest order simple Lie algebras by operators (2.4). The commutation relations defining these algebras are [36]:

$$\begin{aligned}
 so(3) &= \langle e_1, e_2, e_3 \rangle : [e_1, e_2] = e_3, \quad [e_1, e_3] = -e_2, \quad [e_2, e_3] = e_1; \\
 sl(2, \mathbb{R}) &= \langle e_1, e_2, e_3 \rangle : [e_1, e_2] = 2e_2, \quad [e_1, e_3] = -2e_3, \quad [e_2, e_3] = e_1.
 \end{aligned}$$

We start by noting that one of the basis operators e_1, e_2, e_3 can be reduced to one of the five operators (3.2) (see, Lemma 1). We consider in detail the case of operator

$$t\partial_t + x\partial_x \quad (3.3)$$

only, since the remaining cases are handled in a similar way.

Let the basis operator e_1 of the algebras $so(3)$ and $sl(2, \mathbb{R})$ be of the form (3.3). Computing the commutator of e_1 and Q of the form (2.4) yields the relation

$$[e_1, Q] = -\lambda_1\partial_t - \lambda_2\partial_x + [xh'u + xr_x + tr_t]\partial_u.$$

To satisfy the first two commutation relations for each of the algebras under study, the basis operators e_2, e_3 have to be of the form

$$\alpha_1\partial_t + \alpha_2\partial_x + (\gamma(x)u + \mu(t, x))\partial_u,$$

where $\alpha_1, \alpha_2 \in \mathbb{R}$, γ and μ are smooth functions. It is straightforward to verify that these operators cannot satisfy the third commutation relation for either algebra $sl(2, \mathbb{R})$ and $so(3)$.

The lemma is proved.

THEOREM 3. *There are no nonlinear equation (2.8) invariance algebras of which are isomorphic to semi-simple Lie algebras or contain them as subalgebras.*

Proof. Suppose the inverse. Let (2.8) be an equation the invariance algebra of which contain a subalgebra that is semi-simple Lie algebra L . Then by properties of semi-simple Lie algebras there exist linear combinations of the basis elements of L forming the basis of either $so(3)$ or $sl(2, \mathbb{R})$. However, due to Lemma 2 there are no realizations of the algebras $so(3)$, $sl(2, \mathbb{R})$ by operators (2.4). We arrive at the contradiction which proves the theorem.

It follows from Theorem 3 and Levi–Maltsev theorem (see, e.g., [36, 37]) that nonlinear equation (2.8) can admit invariance algebras of the dimension higher than one provided, (i) these algebras are isomorphic to real solvable Lie algebras, or (ii) their finite dimensional subalgebras are real and solvable. Using this fact and also the concept of compositional row for solvable Lie algebras we can perform hierarchical classification of invariant equations starting from the lowest dimensional solvable Lie algebras and increasing dimension by one till we exhaust all possible invariant equations. We start by considering two-dimensional solvable Lie algebras.

There exist two inequivalent two-dimensional solvable Lie algebras [37, 38]

$$A_{2,1} = \langle e_1, e_2 \rangle : [e_1, e_2] = 0;$$

$$A_{2,2} = \langle e_1, e_2 \rangle : [e_1, e_2] = e_2.$$

To construct all possible realizations of the above algebras we take as the first basis element one of the realizations of one-dimensional invariance algebras obtained above. The second operator is looked for in the form (2.4). In the case of commutative algebra $A_{2,1}$ there is no difference between operators e_1 and e_2 , while for the algebra $A_{2,2}$ these operators require separate analysis. We give full computation details for the case when one of the basis elements is of the form A_1^1 given in Theorem 2.

Algebra $A_{2,1}$. Let the operator e_1 be of the form (3.3) and the operator e_2 read as (2.4). It follows from the relation $[e_1, e_2] = 0$ that $\lambda_1 = \lambda_2 = xh' = 0$, $tr_t + xr_x = 0$.

Consequently, we can choose the basis elements of the algebra under study in the form $\langle t\partial_t + x\partial_x, (mu + r(\psi))\partial_u \rangle$, where $m \in \mathbb{R}$, $\psi = tx^{-1}$. Provided $m = 0$ the operator e_2 becomes $r(\psi)\partial_u$. As established earlier, this realization does not satisfy the determining equations. Hence, $m \neq 0$. Making the change of variables

$$\bar{t} = t, \quad \bar{x} = x, \quad v = u + m^{-1}r(\psi)$$

reduces the basis operators in question to the form $\bar{t}\partial_{\bar{t}} + \bar{x}\partial_{\bar{x}}, mv\partial_v$. That is why we can restrict our considerations to the realization $\langle t\partial_t + x\partial_x, u\partial_u \rangle$.

The second determining equation from (2.6) after being written for the operator $u\partial_u$ takes the form $ug_u = 0$, whence it follows that this realization does not satisfy the determining equations. So the realization A_1^1 cannot be extended to a realization of the two-dimensional algebra $A_{2,1}$.

Algebra $A_{2,2}$. If operator e_1 is of the form (3.3), then it follows from $[e_1, e_2] = e_2$ that $\lambda = \lambda_1 = \lambda_2 = 0, xh' = h, tr_t + xr_x = r$.

Next, if the operator e_2 reads as (3.3), then we get from $[e_1, e_2] = e_2$ the erroneous equality $1 = 0$.

So the only possible case is when $e_2 = (mxu + xr(\psi))\partial_u, m \neq 0, \psi = tx^{-1}$, which gives rise to the following realization of the algebra $A_{2,2}$: $\langle t\partial_t + x\partial_x, xu\partial_u \rangle$. This algebra is indeed invariance algebra of an equation from the class (2.8) and the functions f and g read as

$$g = -2x^{-1} \ln |u| + x^{-1} \tilde{g}(\psi), \quad f = x^{-2}u \ln^2 |u| - x^{-2} \tilde{g}(\psi)u \ln |u| + x^{-2}u \tilde{f}(\psi), \quad \psi = tx^{-1}.$$

Analysis of the remaining realizations of one-dimensional Lie algebras yields 10 inequivalent $A_{2,1}$ - and $A_{2,2}$ -invariant equations (see the assertions below). What is more, the obtained (two-dimensional) algebras are maximal symmetry algebras of the corresponding equations.

THEOREM 4. *There are at most four inequivalent $A_{2,1}$ -invariant nonlinear equation (2.8). Below we list the realizations of $A_{2,1}$ and the corresponding expressions for f and g .*

- 1) $\langle \partial_t, \sigma(x)u\partial_u, \rangle : g = -2\sigma'\sigma^{-1} \ln |u|,$
 $f = (\sigma'\sigma^{-1})^2u \ln^2 |u| - \sigma^{-1}\sigma'u \ln |u| + u \tilde{f}(x), \quad \sigma' \neq 0;$
- 2) $\langle \partial_t, \partial_x \rangle : g = \tilde{g}(u), \quad f = \tilde{f}(u), \quad \tilde{g}_u \neq 0;$
- 3) $\langle \partial_x, \partial_t + u\partial_u \rangle : g = \tilde{g}(\omega), \quad f = \exp(t) \tilde{f}(\omega), \quad \omega = \exp(-t), \quad \tilde{g}_\omega \neq 0;$
- 4) $\langle \sigma(x)u\partial_u, \partial_t - \frac{1}{2}k\sigma(x)\psi(x)u\partial_u \rangle : g = -2\sigma'\sigma^{-1} \ln |u| + kt + \tilde{g}(x),$
 $f = (\sigma'\sigma^{-1})^2u \ln^2 |u| - \sigma^{-1}\sigma''u \ln |u| - \sigma^{-1}\sigma'(kt + \tilde{g}(x))u \ln |u|$
 $+ u \left[\frac{1}{2}k\sigma'\sigma^{-1}t + \frac{1}{4}k^2t^2 + \frac{1}{2}k\tilde{g}(x) + \tilde{f}(x) \right],$
 $k \neq 0, \sigma' \neq 0, \psi = \int \sigma^{-1} dx.$

THEOREM 5. *There exist at most six inequivalent $A_{2,2}$ -invariant nonlinear equation (2.8). Below we list the realizations of $A_{2,1}$ and the corresponding expressions for f and g .*

- 1) $\langle t\partial_t + x\partial_x, k^{-1}|x|^k u\partial_u \rangle : g = x^{-1}(-2k \ln |u| + \tilde{g}(\psi)),$
 $f = x^{-2}u(-k^2 \ln^2 |u| + k\tilde{g}(\psi) \ln |u| + k(k-1) \ln |u| + \tilde{f}(\psi)),$
 $k \neq 0, \psi = tx^{-1};$
- 2) $\langle \partial_t + \beta\partial_x, \exp(\beta^{-1}x)u\partial_u \rangle : g = -2\beta^{-1} \ln |u| + \tilde{g}(\eta),$
 $f = \beta^{-2}u \ln^2 |u| - (\beta^{-2} + \beta^{-1}\tilde{g}(\eta))u \ln |u| + u\tilde{f}(\eta),$
 $\beta > 0, \eta = x - \beta t;$
- 3) $\langle -t\partial_t - x\partial_x, \partial_t + \beta\partial_x \rangle : g = \eta^{-1}\tilde{g}(u), \quad f = \eta^{-2}\tilde{f}(u), \quad \beta \geq 0,$
 $\eta = x - \beta t, \quad \tilde{g}_u \neq 0;$
- 4) $\langle -t\partial_t - x\partial_x, \partial_t + mx^{-1}u\partial_u \rangle : g = x^{-1}(2m\psi + \tilde{g}(\omega)),$
 $f = x^{-1}[-2m\psi u - 2m\psi - 2 - \tilde{g}(\omega) + \exp(m\psi)\tilde{g}(\omega)],$
 $m > 0, \omega = u \exp(-m\psi), \psi = tx^{-1}, \tilde{g}_\omega \neq 0;$
- 5) $\langle \partial_x, e^x u\partial_u \rangle : g = -2 \ln |u| + \tilde{g}(t), \quad f = u \ln^2 |u|$
 $-u \ln |u|(1 + \tilde{g}(t)) + u\tilde{f}(t);$
- 6) $\langle -t\partial_t - x\partial_x, \partial_x \rangle : g = t^{-1}\tilde{g}(u), \quad f = t^{-2}\tilde{f}(u), \quad \tilde{g}_u \neq 0.$

4.2. Completing Group Classification of (2.8)

As the invariant equations obtained in the previous subsection contain arbitrary functions of at most one variable, we can use the standard Lie–Ovsyannikov approach to complete group classification of (2.8). We give the computation details for the case of the first $A_{2,1}$ -invariant equation. The remaining cases are handled in a similar way.

Putting $g = -2\sigma'\sigma^{-1} \ln |u|$, $f = (\sigma'\sigma^{-1})u \ln^2 |u| - \sigma^{-1}\sigma''u \ln |u| + u\tilde{f}(x)$, $\sigma = \sigma(x)$, $\sigma' \neq 0$ we rewrite the first determining equation to become:

$$-2h' + 2\lambda\sigma'\sigma^{-1} \ln |u| = -2(\lambda x + \lambda_2)(\sigma'\sigma^{-1})'_x \ln |u| - 2h\sigma'\sigma^{-1} - 2r\sigma'\sigma^{-1}u^{-1}.$$

As $h = f(x)$, $\sigma = \sigma(x)$, $r = r(t, x)$, $\lambda, \lambda_2 \in \mathbb{R}$, the above relation is equivalent to the following ones:

$$h' = \sigma'\sigma^{-1}h, \quad r = 0, \quad \lambda\sigma'\sigma^{-1} = -(\lambda x + \lambda_2)(\sigma'\sigma^{-1})'.$$

If σ is an arbitrary function, then $\lambda = \lambda_2 = r = 0$, $h = C\sigma$, $C \in \mathbb{R}$ and $\langle \partial_t, \sigma(x)u\partial_u \rangle$ is the maximal symmetry algebra of the equation under study. Hence extension of symmetry algebra is only possible when the function $\psi = \sigma'\sigma^{-1}$ is a (non-vanishing identically) solution of equation

$$(\alpha x + \beta)\psi' + \alpha\psi = 0, \quad \alpha, \beta \in \mathbb{R}, \quad |\alpha| + |\beta| \neq 0.$$

If $\alpha \neq 0$, then utilizing displacements by x we can get $\beta = 0$, so that $\psi = mx^{-1}$, $m \neq 0$. Integrating the remaining determining equations we get

$$g = -2mx^{-1} \ln |u|, \quad f = mx^{-2}[mu \ln^2 |u| - (m-1)u \ln |u| + nu], \quad m \neq 0, \quad m, n \in \mathbb{R}.$$

The maximal invariance algebra of the obtained equation is the three-dimensional Lie algebra $\langle \partial_t, |x|^m u \partial_u, t \partial_t + x \partial_x \rangle$ isomorphic to $A_{3,7}$.

Next, if $\alpha = 0$, then $\beta \neq 0$ and what is more $\psi = m$, $m \neq 0$. If this is the case we have

$$g = \ln |u|, \quad f = \frac{1}{4} u \ln^2 |u| - \frac{1}{4} u \ln |u| + nu, \quad n \in \mathbb{R}.$$

The maximal invariance algebra of the above equation reads as

$$\langle \partial_t, \partial_x, \exp\left(-\frac{1}{2}x\right) u \partial_u \rangle.$$

It is isomorphic to $A_{3,2}$.

Similarly we prove that the list of inequivalent equations of the form (2.8) admitting three-dimensional symmetry algebras is exhausted by the equations given below. Note that the presented algebras are maximal. This means, in particular, that the maximal symmetry algebra of Equation (2.8) is at most three-dimensional.

$A_{3,2}$ -invariant equations

- 1) $u_{tt} = u_{xx} + u_x \ln |u| + \frac{1}{4} u \ln^2 |u| - \frac{1}{4} u \ln |u| + nu$ ($n \in \mathbb{R}$) : $\langle \partial_t, \partial_x, \exp\left(-\frac{1}{2}x\right) u \partial_u \rangle$;
- 2) $u_{tt} = u_{xx} + m[\ln |u| - t]u_x + \frac{m^2}{4} u[(\ln |u| - t)(\ln |u| - t - 1)] + nu$ ($m > 0$, $n \in \mathbb{R}$) : $\langle \partial_x, \partial_t + u \partial_u, \exp\left(-\frac{1}{2}mx\right) u \partial_u \rangle$.

$A_{3,4}$ -invariant equations

- 1) $u_{tt} = u_{xx} + x^{-1}[2 \ln |u| + mx^{-1}t + n]u_x + x^{-2}u \ln |u| + (mx^{-1}t + n - 2)x^{-2}u \ln |u| + \frac{1}{4}m^2x^{-4}t^2u + \frac{1}{2}m(n - 3)x^{-3}tu + px^{-2}u$ ($m \neq 0, n, p \in \mathbb{R}$) : $\langle t \partial_t + x \partial_x, x^{-1}u \partial_u, \partial_t - \frac{m}{2}x^{-1} \ln |x| u \partial_u \rangle$.

$A_{3,5}$ -invariant equations

- 1) $u_{tt} = u_{xx} + |u|^m u_x + n|u|^{1+2m}$ ($m \neq 0, n \in \mathbb{R}$) : $\langle \partial_t, \partial_x, t \partial_t + x \partial_x - m^{-1}u \partial_u \rangle$;
- 2) $u_{tt} = u_{xx} + e^u u_x + ne^{2u}$ ($n \in \mathbb{R}$) : $\langle \partial_t, \partial_x, t \partial_t + x \partial_x - \partial_u \rangle$;
- 3) $u_{tt} = u_{xx} - x^{-1}[2 \ln |u| - mx^{-1}t - n]u_x + x^{-2}u \ln^2 |u| - x^{-2}(mx^{-1}t + n)u \ln |u| + ux^{-2} \left[\frac{m}{4}x^{-2}t^2 + \frac{m}{2}(n - 1)x^{-1}t + p \right]$ ($m, n, p \in \mathbb{R}$) : $\langle t \partial_t + x \partial_x, xu \partial_u, \partial_t + \frac{m}{4}x^{-1}u \partial_u \rangle$.

$A_{3,7}$ -invariant equations

- 1) $u_{tt} = u_{xx} - 2mx^{-1}u_x \ln |u| + mx^{-2}[mu \ln^2 |u| - (m - 1)u \ln |u| + nu]$ ($m \neq 0, 1; n \in \mathbb{R}$) : $\langle \partial_t, |x|^m u \partial_u, t \partial_t + x \partial_x \rangle$;
- 2) $u_{tt} = u_{xx} - x^{-1}[2k + \ln |u| - mx^{-1}t - n]u_x + k^2x^{-2}u \ln^2 |u| - kx^{-2}[m x^{-1} + k + n - 1]u \ln |u| + \frac{1}{2}m(k - 2 + n)tx^{-3}u + \frac{1}{4}m^2t^2x^{-4}u + px^{-2}u$ ($|k| \neq 0, 1; m \neq 0, n, p \in \mathbb{R}$) : $\langle t \partial_t + x \partial_x, |x|^k u \partial_u, \partial_t + \frac{m}{2(1+k)}x^{-1}u \partial_u \rangle$.

This completes group classification of nonlinear equation (2.8).

5. Group Classification of Equation (2.12)

Omitting the intermediate calculation details we give the determining equations for symmetry operators admitted by Equation (2.12).

ASSERTION 3. The maximal invariance group of PDE (2.12) is generated by the infinitesimal operator

$$Q = \tau(t)\partial_t + \xi(x)\partial_x + [h(t)u + r(t, x)]\partial_u, \quad (4.1)$$

where τ, ξ, h, r, f, g are smooth functions satisfying the conditions

$$\begin{aligned} r_{tx} + f[h - \tau_t - \xi_x] &= gr_x + \tau f_t + \xi f_x + [hu + r]f_u, \\ h_t &= \tau_t g + \tau g_t + \xi g_x. \end{aligned} \quad (4.2)$$

ASSERTION 4. The equivalence group \mathcal{E} of (2.12) is formed by the following transformations of the space V :

$$\begin{aligned} (1) \quad \bar{t} &= T(t), \quad \bar{x} = X(x), \quad v = U(t)u + Y(t, x), \quad T'X'U \neq 0; \\ (2) \quad \bar{t} &= T(x), \quad \bar{x} = X(t), \quad v = \Psi(x)\Phi(t, x)u + Y(t, x), \quad T'X'\Psi \neq 0, \\ \Phi(t, x) &= \exp[-\int g(t, x)dt], \quad g_x \neq 0. \end{aligned} \quad (4.3)$$

Given the arbitrary functions g and f Equation (4.2) holds only when $\tau = h = \xi = r = 0$. This means that in the generic case the maximal invariance group of Equation (2.12) is the trivial group of identical transformations.

We begin symmetry classification of (2.12) by constructing equations that admit one-dimensional invariance algebras.

LEMMA 3. *There exist transformations (4.3) reducing operator (4.1) to one of the seven canonical forms*

$$\begin{aligned} Q &= t\partial_t + x\partial_x; \quad Q = \partial_t; \quad Q = \partial_x + tu\partial_u; \\ Q &= \partial_x + \epsilon u\partial_u, \quad \epsilon = 0, 1; \quad Q = tu\partial_u, \\ Q &= u\partial_u, \quad Q = r(t, x)\partial_u, \quad r \neq 0. \end{aligned} \quad (4.4)$$

Proof. Transformations (4.3) reduce operator Q (4.1) to become

$$\tilde{Q} = \tau T'\partial_{\bar{t}} + \xi X'\partial_{\bar{x}} + [(\tau U' + Uh)u + \tau Y_t + \xi Y_x + Ur]\partial_v. \quad (4.5)$$

Provided $\sigma \cdot \xi \neq 0$, we can choose non-vanishing identically solutions of the equations

$$\tau T' = T, \quad \xi X' = X, \quad \tau U' + hU = 0, \quad \tau Y_t + \xi Y_x + Ur = 0$$

as T, X, U, Y thus getting operator \tilde{Q} (4.5) in the form $\tilde{Q} = \bar{t}\partial_{\bar{t}} + \bar{x}\partial_{\bar{x}}$. If $\tau \neq 0$, and $\xi = 0$, then taking the solutions of the equations

$$\tau T' = 1, \tau U' + hU = 0 (U \neq 0), \tau Y_t + Ur = 0$$

as T, U, Y reduces operator (4.1) to the form $\tilde{Q} = \partial_{\bar{t}}$. If $\tau = 0, \xi \neq 0$, then under $h' \neq 0$ we get the operator $\tilde{Q} = \partial_{\bar{x}} + \bar{t}v\partial_v$. Next, if $h' = 0$, we arrive at the operator $\tilde{Q} = \partial_{\bar{x}} + \epsilon v\partial_v$, where either $\epsilon = 0$ or $\epsilon = 1$.

Finally, the case $\tau = \xi = 0$, gives rise to the operators $\tilde{Q} = \bar{t}v\partial_v, \tilde{Q} = v\partial_v, \tilde{Q} = r(\bar{t}, \bar{x})\partial_v$. Rewriting them in the initial variables we get the operators listed in the statement of lemma. The lemma is proved.

THEOREM 6. *There exist at most three inequivalent nonlinear equation (2.12) that admit one-dimensional invariance algebras. The form of functions f, g and the corresponding symmetry algebras are given below.*

$$\begin{aligned} A_1^1 &= \langle t\partial_t + x\partial_x \rangle : g = t^{-1}\tilde{g}(\omega), f = t^{-2}f(u, \omega), \omega = tx^{-1}, \tilde{g}'_\omega \neq 0, f_{uu} \neq 0; \\ A_1^2 &= \langle \partial_t \rangle : g = \tilde{g}(x), f = \tilde{f}(x, u), \tilde{g}' \neq 0, \tilde{f}_{uu} \neq 0; \\ A_1^3 &= \langle \partial_x + tu\partial_u \rangle : g = x + \tilde{g}(t), f = e^{tx}\tilde{f}(t, \omega), \omega = e^{-tx}u, \tilde{f}_{\omega\omega} \neq 0. \end{aligned}$$

Proof. If Equation (2.12) admits one-parameter transformation group, then the latter is generated by infinitesimal operator (4.1). According to Lemma 3 there exist equivalence transformations (4.3) reducing this operator to one of the seven canonical operators (4.4). Now we need to solve determining Equation (4.2) for each of these operators. The first three operators yield invariant equations and corresponding symmetry algebras given in the statement of theorem. The next two operators give rise to inconsistent equations.

Finally, the remaining operators yield that the functions f and g are linear in u , which means that the corresponding invariant equations are linear.

It is straightforward to verify that for the case of arbitrary functions \tilde{f}, \tilde{g} the corresponding one-dimensional algebras are maximal in Lie’s sense.

The theorem is proved.

We proceed now to analyzing Equation (2.12) which admit two-dimensional symmetry algebras.

THEOREM 7. *There exist at most three inequivalent nonlinear equation (2.12) that admit two-dimensional symmetry algebras, all of them being $A_{2,2}$ -invariant equations. The forms of functions f and g and the corresponding realizations of the Lie algebra $A_{2,2}$ are given below*

$$\begin{aligned} A_{2,2}^1 &= \langle t\partial_t + x\partial_x, t^2\partial_t + x^2\partial_x + mut\partial_u \rangle (m \in \mathbb{R}) : \\ g &= [mt + (k - m)x]t^{-1}(t - x)^{-1}, k \neq 0, \\ f &= |t - x|^{m-2}|x|^{-m}\tilde{f}(\omega), \\ \omega &= u|t - x|^{-m}|x|^m, \tilde{f}_{\omega\omega} \neq 0; \end{aligned}$$

$$\begin{aligned}
 A_{2,2}^2 &= \langle t\partial_t + x\partial_x, t^2\partial_t + mtu\partial_u \rangle \ (m \in \mathbb{R}) : \\
 &g = t^{-2}[kx + mt], \ k \neq 0, \ f = |t|^{m-2}|x|^{-m}\tilde{f}(\omega), \\
 &\omega = |t|^{-m}|x|^m u, \ \tilde{f}_{\omega\omega} \neq 0; \\
 A_{2,2}^3 &= \langle t\partial_t + x\partial_x, x^2\partial_x + tu\partial_u \rangle : \\
 &g = (tx)^{-1}(mx - t) \ (m \in \mathbb{R}), \ f = x^{-2}\exp(-tx^{-1})\tilde{f}(\omega), \\
 &\omega = u\exp(tx^{-1}), \ \tilde{f}_{\omega\omega} \neq 0.
 \end{aligned}$$

To prove the theorem we need to extend realizations A_i^i ($i = 1, 2, 3$) to realizations of the algebras $A_{2,1}$, $A_{2,2}$ by operators (4.1). We skip the calculation details.

Note that if the functions \tilde{f} are arbitrary, then the invariance algebras given in the statement of Theorem 7 are maximal.

Now we can complete the group classification presented in Theorem 7 utilizing the Lie–Ovsyannikov classification routine. Consider in some detail the case of $A_{2,2}^1$ -invariant equations (the remaining cases are treated in a similar way). The second determining equation from Equation (4.2) reads now as

$$(t - x)^2 h_t = t^{-1}\tau_t[m(t - x)^2 + kx(t - x)] + \tau[-t^{-2}m(t - x)^2 - 2kt^{-1}x + kt^{-2}x^2] + k\xi. \tag{4.6}$$

Differentiating right- and left-hand sides of (4.6) twice by x yields

$$h_t = (m - k)(t^{-1}\tau_t - t^{-2}\tau) + k\xi''.$$

Hence we get $\xi''' = 0$ and

$$\begin{aligned}
 \xi &= \lambda_1 x^2 + \lambda_2 x + \lambda_3, \ \lambda_1, \lambda_2, \lambda_3 \in \mathbb{R}, \\
 h &= (m - k)t^{-1}\tau + \lambda_1 kt + \lambda_4, \ \lambda_4 \in \mathbb{R}.
 \end{aligned}$$

With account of the above facts we obtain from (4.6) that $\tau = \lambda_1 t^2 + \lambda_2 t + \lambda_3$. So it follows from (4.6) that the coefficients of infinitesimal operator (4.1), which generates symmetry group of $A_{2,2}^1$ -invariant equation, read necessarily as

$$\begin{aligned}
 \tau &= \lambda_1 t^2 + \lambda_2 t + \lambda_3, \\
 \xi &= \lambda_1 x^2 + \lambda_2 x + \lambda_3, \\
 h &= m\lambda_1 t + (m - k)\lambda_3 t^{-1} + (m - k)\lambda_2 + \lambda_4, \ \lambda_1, \lambda_2, \lambda_3, \lambda_4 \in \mathbb{R}.
 \end{aligned}$$

Consequently, the first determining equation from (4.2) takes the form

$$(a\omega + b)\tilde{f}_{\omega} - a\tilde{f} = c, \tag{4.8}$$

where $a, b, c \in \mathbb{R}$, $|a| + |b| \neq 0$. On the other hand, it follows from (4.8)

$$(a\omega + b)\tilde{f}_{\omega\omega} = 0,$$

whence $f_{\omega\omega} = 0$. We arrive at the contradiction which proves that there is no extension of the realization $A_{2,2}^1$ in question to a higher dimensional invariance algebra of the Equation (2.12). Analyzing $A_{2,2}^2$ - and $A_{2,3}^2$ -invariant equations we arrive at the same conclusion.

Consequently, there are no nonlinear equations of the form (2.12) whose maximal invariance algebras are solvable Lie algebras of the dimension higher than two. Next, as the algebra $sl(2, \mathbb{R})$ contains two-dimensional subalgebra isomorphic to $A_{2,2}$, there

are no nonlinear equation (2.12), whose invariance algebras are either isomorphic to $sl(2, \mathbb{R})$ or contain it as a subalgebra. Finally, we verify that there are no realizations of the algebra $so(3)$ by operators (4.1).

Summing up the above reasonings we arrive at the following assertion.

THEOREM 8. *A nonlinear equation (2.12) having non-trivial symmetry properties is equivalent to one of the equations listed in Theorems 6 and 7.*

This completes group classification of the class of nonlinear PDEs (2.12).

6. Group Classification of Equation (2.13)

As earlier, we present the results of the first step of our group classification algorithm skipping derivation details.

ASSERTION 5. Invariance group of Equation (2.13) is generated by the infinitesimal operator

$$Q = \tau(t)\partial_t + \xi(x)\partial_x + (ku + r(t, x))\partial_u, \tag{5.1}$$

where k is a constant and τ, ξ, r, f are functions satisfying the relation

$$r_{tx} + [k - \tau' - \xi']f = \tau f_t + \xi f_x + [ku + r]f_u. \tag{5.2}$$

ASSERTION 6. Equivalence group \mathcal{E} of the class of Equation (2.13) is formed by the following transformations:

$$\begin{aligned} (1) \quad & \bar{t} = T(t), \quad \bar{x} = X(x), \quad v = mu + Y(t, x), \\ (2) \quad & \bar{t} = T(x), \quad \bar{x} = X(t), \quad v = mu + Y(t, x), \quad T'X'm \neq 0. \end{aligned} \tag{5.3}$$

Note that given an arbitrary f it follows from (5.2) that $\tau = \xi = k = r = 0$, i.e., the group admitted is trivial. To obtain equations with non-trivial symmetry we need to specify properly the function f . To this end we perform classification of equations under study admitting one-dimensional invariance algebras.

LEMMA 4. *There exist transformations from the group \mathcal{E} (5.3) which reduce (5.1) to one of the four canonical forms:*

$$\begin{aligned} Q &= \partial_t + \partial_x + \epsilon u \partial_u \quad (\epsilon = 0, 1); \\ Q &= \partial_t + \epsilon u \partial_u \quad (\epsilon = 0, 1); \\ Q &= u \partial_u, \quad Q = g(t, x) \partial_u \quad (g \neq 0). \end{aligned}$$

Proof. Utilizing transformations (1) from (5.3) we reduce the operator Q to one of the following forms:

$$Q = \partial_{\tilde{t}} + \partial_{\tilde{x}} + \epsilon v \partial_v \quad (\epsilon = 0, 1) :$$

$$Q = \partial_{\tilde{t}} + \epsilon v \partial_v \quad (\epsilon = 0, 1);$$

$$Q = \partial_{\tilde{x}} + \epsilon v \partial_v \quad (\epsilon = 0, 1);$$

$$Q = v \partial_v, \quad Q = g(\tilde{t}, \tilde{x}) \partial_v \quad (g \neq 0).$$

Next we note that the change of variables $\tilde{t} = \bar{x}$, $\tilde{x} = \bar{t}$, $\tilde{v} = v$ which is of the form (2) from (5.3) transforms the second operator into the third one. Rewriting the obtained operators in the initial variables completes the proof.

THEOREM 9. *There exist exactly two nonlinear equations of the form (2.13) admitting one-dimensional invariance algebras. The corresponding expressions for function f and invariance algebras are given below.*

$$A_1^1 = \langle \partial_t + \partial_x + \epsilon u \partial_u \rangle \quad (\epsilon = 0, 1) : f = e^{\epsilon t} \tilde{f}(\theta, \omega), \quad \theta = t - x, \quad \omega = e^{-\epsilon t} u; \quad \tilde{f}_{\omega\omega} \neq 0;$$

$$A_1^2 = \langle \partial_t + \epsilon u \partial_u \rangle \quad (\epsilon = 0, 1) : f = e^{\epsilon t} \tilde{f}(x, \omega), \quad \omega = e^{-\epsilon t} u, \quad \tilde{f}_{\omega\omega} \neq 0.$$

To prove the theorem, it suffices to select those operators from the list given in Lemma 4 that can be invariance algebra of nonlinear equation of the form (2.13). To this end we need to solve Equation (5.2) for each of the operators in question.

The first two operators yield A_1^1 - and A_1^2 -invariant equations. The last two operators gives rise to linear invariant equation (2.13), which are not taken into account.

What is more, if the function \tilde{f} is arbitrary, then the algebras A_1^1 and A_1^2 are maximal invariance algebras of the corresponding equations.

Next, we classify nonlinear equations admitting symmetry algebras of the dimension higher than one. We begin by considering equations whose invariance algebras contain semi-simple subalgebras. It turns out that the class of operators (5.1) contain no realizations of the algebra $so(3)$. Furthermore it contains the four inequivalent realizations of the algebra $sl(2, \mathbb{R})$ given below.

$$(1) \left\langle \partial_t, \frac{1}{2} e^{2t} \partial_t, -\frac{1}{2} e^{-2t} \partial_t \right\rangle;$$

$$(2) \left\langle \partial_t, \frac{1}{2} e^{2t} (\partial_t + \partial_u), -\frac{1}{2} e^{-2t} (\partial_t - \partial_u) \right\rangle;$$

$$(3) \left\langle \partial_t, \frac{1}{2} e^{2t} (\partial_t + x \partial_u), -\frac{1}{2} e^{-2t} (\partial_t - x \partial_u) \right\rangle;$$

$$(4) \left\langle \partial_t + \partial_x, \frac{1}{2} e^{2t} \partial_t + \frac{1}{2} e^{2x} \partial_x, -\frac{1}{2} e^{-2t} \partial_t - \frac{1}{2} e^{-2x} \partial_x + \epsilon [e^{-2x} - e^{-2t}] \partial_u \right\rangle, \quad \epsilon = 0, 1.$$

Before starting analysis of $sl(2, \mathbb{R})$ -invariant equations we briefly review the group properties of the Liouville equation

$$u_{tx} = \lambda e^u, \quad \lambda \neq 0. \tag{5.4}$$

It is a common knowledge that the maximal invariance group of this equation is the infinite-parameter group generated by the following infinitesimal operator [39]:

$$Q = h(t)\partial_t + g(x)\partial_x - (h' + g')\partial_u,$$

where h and g are arbitrary smooth functions. Note that due to this fact the Liouville equation can be linearized by a non-local change of variables (see, e.g., [11, 40, 41]).

After a simple algebra we obtain that realizations (1), (3), (4) with $\epsilon = 1$ cannot be invariance algebras of a nonlinear equation of the form (2.13). Realization (2) is the invariance algebra of equation

$$u_{tx} = \tilde{f}(x)e^{-2u}, \quad \tilde{f} \neq 0,$$

which reduces to Equation (5.4) via the change of variables

$$t = t, \quad x = x, \quad u = -\frac{1}{2}(v - \ln|\tilde{f}|), \quad v = v(t, x).$$

Finally making use of the change of variables

$$\tilde{t} = e^{-2t}, \quad \tilde{x} = e^{-2x}, \quad v = u$$

we rewrite (4) under $\epsilon = 0$ to become

$$\langle \partial_t + \partial_x, t\partial_t + x\partial_x, t^2\partial_t + x^2\partial_x \rangle.$$

The corresponding invariant equation reads as

$$u_{tx} = (t - x)^{-2} \tilde{f}(u), \quad \tilde{f}_{uu} \neq 0. \tag{5.5}$$

If the function \tilde{f} is arbitrary, then the above presented realization is the maximal invariance algebra of the equation under study. Using the Lie–Ovsyannikov algorithm we establish that extension of symmetry is only possible when $\tilde{f} = \lambda e^u + 2$. However the corresponding equation is reduced to the Liouville equation by the change of variables

$$t = t, \quad x = x, \quad u = v(t, x) + 2 \ln|t - x|.$$

Thus the only inequivalent nonlinear equation (2.13) whose invariance algebras contain semi-simple subalgebras are given in (5.4) and (5.5), where \tilde{f} is an arbitrary smooth function of u .

To complete group classification of Equation (2.13) we need to describe equations whose invariance algebras are solvable Lie algebras of the dimension higher than one. We begin with those realizations of two-dimensional Lie algebras $A_{2,1}$, $A_{2,2}$, which can be admitted by nonlinear equation (2.13).

It turns out that the class of operators (5.1) contains one inequivalent realization of the algebra $A_{2,1}$ which meets the invariance requirements, namely,

$$\langle \partial_t + \epsilon_1 u \partial_u, \partial_x + \epsilon_2 u \partial_u \rangle \quad (\epsilon_1 = 0, 1; \quad \epsilon_2 = 0, 1).$$

The corresponding invariant equation reads as

$$u_{tx} = \exp(\epsilon_1 t + \epsilon_2 x) \tilde{f}(\omega), \quad \omega = u \exp(-\epsilon_1 t - \epsilon_2 x). \tag{5.6}$$

Analysis of Equation (5.6) with arbitrary $f(\omega)$ shows that under $\epsilon_1 + \epsilon_2 \neq 0$ the above realization is its maximal invariance algebra. Provided $\epsilon_1 = \epsilon_2 = 0$ the equation takes the form

$$u_{tx} = f(u) \tag{5.7}$$

and its maximal invariance algebra is the three-dimensional Lie algebra of the operators

$$\langle \partial_t, \partial_x, t\partial_t - x\partial_x \rangle,$$

which is isomorphic to $A_{3,6}$.

It is a common knowledge (see, e.g., [17–19]) that (5.7) admits extension of its symmetry if it is equivalent either to the Liouville equation (5.4) or to the nonlinear d'Alembert equation

$$u_{tx} = \lambda|u|^{n+1}, \quad \lambda \neq 0, \quad n \neq 0, -1. \quad (5.8)$$

The maximal invariance algebra of (5.8) is the four-dimensional Lie algebra

$$\langle t\partial_t - \frac{1}{n}u\partial_u, x\partial_x - \frac{1}{n}u\partial_u, \partial_t, \partial_x \rangle.$$

It is isomorphic to the algebra $A_{2,2} \oplus A_{2,2}$.

Extension of symmetry algebra of Equation (5.6) with $\epsilon_1 = 1, \epsilon_2 = 0$, is only possible when:

$$u_{tx} = \lambda e^{-mt} |u|^{m+1}, \quad \lambda \neq 0, \quad m \neq 0, -1; \quad (5.9)$$

$$u_{tx} = \lambda e^t \exp(ue^{-t}), \quad \lambda \neq 0. \quad (5.10)$$

The maximal invariance algebra of (5.9) is the four-dimensional Lie algebra of operators

$$\langle \partial_t + u\partial_u, e^{mt}\partial_t, \partial_x, x\partial_x - \frac{1}{m}u\partial_u \rangle,$$

which is isomorphic to $A_{2,2} \oplus A_{2,2}$. Note that the change of variables

$$\bar{t} = e^{-mt}, \quad \bar{x} = x, \quad u = v(\bar{t}, \bar{x})$$

reduces the above equation to the form (5.8).

The maximal invariance algebra of (5.10) is spanned by the operators

$$\langle \partial_t + u\partial_u, \partial_x, x\partial_x - e^t\partial_u \rangle$$

and is isomorphic to $A_1 \oplus A_{2,2}$.

Analyzing $A_{2,2}$ -invariant equations we arrive at the following conclusion. The class of operators (5.1) contains six inequivalent realizations of the algebra $A_{2,2}$ which meet the invariance requirements

- (1) $\langle -t\partial_t + x\partial_x, \partial_t \rangle;$
- (2) $\langle -t\partial_t - x\partial_x, \partial_t + \partial_x \rangle;$
- (3) $\langle -t\partial_t - x\partial_x + u\partial_u, \partial_t + \partial_x \rangle;$
- (4) $\langle -t\partial_t + \partial_u, \partial_t \rangle;$
- (5) $\langle -t\partial_t - x\partial_x - u\partial_u, \partial_t \rangle;$
- (6) $\langle -t\partial_t - x\partial_x, \partial_t \rangle.$

(5.11)

Equation invariant under realization (1) reads as

$$u_{tx} = \exp(x^{-1}u). \quad (5.12)$$

Its maximal symmetry algebra is the three-dimensional Lie algebra

$$\langle -t\partial_t + x\partial_x, \partial_t, x\partial_x + u\partial_u \rangle$$

isomorphic to $A_{2,2} \oplus A_1$. Note that the change of variables

$$\bar{t} = x, \quad \bar{x} = e^t, \quad u = v(\bar{t}, \bar{x})$$

reduces (5.12) to the form (5.10).

Equation invariant under the second realization of $A_{2,2}$ is of the form (5.5). We have already obtained this equation while describing $sl(2, \mathbb{R})$ -invariant equations.

Realizations (3) and (4) give no new invariant equations as well.

New invariant equation are obtained when we consider the fifth realization from (5.10). It has the form

$$u_{tx} = x^{-1} \tilde{f}(\omega), \quad \omega = x^{-1}u.$$

If the function \tilde{f} is arbitrary, then the realization in question is the maximal invariance algebra of the above equation. Further extension of symmetry properties is only possible if $\tilde{f}(\omega) = \lambda|\omega|^{m+1}$ which gives the following invariant equation:

$$u_{tx} = \lambda|x|^{-m-2}|u|^{m+1}, \quad \lambda \neq 0, \quad m \neq 0, -1, -2.$$

Its maximal symmetry algebra is the three-dimensional Lie algebra having the basis

$$\langle \partial_t, t\partial_t + x\partial_x + u\partial_u, x\partial_x + \frac{m+1}{m}u\partial_u \rangle.$$

This algebra is isomorphic to $A_{2,2} \oplus A_1$.

We sum up the above results in the following assertion.

THEOREM 10. *The Liouville equation $u_{tx} = \lambda e^u$, $\lambda \neq 0$, has the highest symmetry among Equations (2.13). Its maximal invariance algebra is infinite-dimensional and spanned by the following infinite set of basis operators:*

$$Q = h(t)\partial_t + g(x)\partial_x - (h'(t) + g'(x))\partial_u.$$

Here h and g are arbitrary smooth functions. Next there exist exactly nine inequivalent equations of the form (2.13) maximal invariance algebras of which have dimension higher than one. These equations and their invariance algebras are given in Table 1.

7. Group Classification of Equation (2.7)

The first step of the algorithm of group classification of (2.7)

$$u_{tt} = u_{xx} + F(t, x, u, u_x), \quad F_{u_x u_x} \neq 0$$

has been partially performed in the second chapter. It follows from Theorem 1 that the invariance group of Equation (2.7) is generated by infinitesimal operator (2.4). What is more, the real constants $\lambda, \lambda_1, \lambda_2$ and real-valued functions $h = h(x), r = r(t, x), F = F(t, x, u, u_x)$ obey relation (2.5). The equivalence group of the class of Equation (2.7) is formed by transformations (3.1).

The above enumerated facts enable using the results of group classification of Equation (2.8) in order to classify invariant equations of the form (2.7). In particular, using Lemmas 1 and 2 it is straightforward to verify that the following assertions hold true.

Table 1 Invariant equation (2.13)

Number	Function f	Symmetry operators	Invariance algebra type
1	$e^t \tilde{f}(\omega), \omega = ue^{-t}, \tilde{f}_{\omega\omega} \neq 0$	$\partial_t + u\partial_u, \partial_x$	$A_{2.1}$
2	$e^{t+x} \tilde{f}(\omega), \omega = ue^{-t-x}, \tilde{f}_{\omega\omega} \neq 0$	$\partial_t + u\partial_u, \partial_x + u\partial_u$	$A_{2.1}$
3	$(t-x)^{-3} \tilde{f}(\omega), \omega = (t-x)u, \tilde{f}_{\omega\omega} \neq 0$	$-t\partial_t - x\partial_x + u\partial_u, \partial_t + \partial_x$	$A_{2.2}$
4	$x^{-1} \tilde{f}(\omega), \omega = x^{-1}u, \tilde{f}_{\omega\omega} \neq 0$	$-t\partial_t - x\partial_x - u\partial_u, \partial_t$	$A_{2.2}$
5	$(t-x)^{-2} \tilde{f}(u), \tilde{f}_{uu} \neq 0$	$\partial_t + \partial_x, t\partial_t + x\partial_x, t^2\partial_t + x^2\partial_x$	$sl(2, R)$
6	$\exp(x^{-1}u)$	$-t\partial_t + x\partial_u, \partial_t, x\partial_x + u\partial_u$	$A_{2.2} \oplus A_1$
7	$\lambda x ^{-m-2} u ^{m+1}, \lambda \neq 0, m \neq 0, -1, -2$	$\partial_t, t\partial_t - \frac{1}{m}u\partial_u, x\partial_x + \frac{m+1}{m}u\partial_u$	$A_{2.2} \oplus A_1$
8	$\tilde{f}(u), \tilde{f}_{uu} \neq 0$	$\partial_t, \partial_x, -t\partial_t - x\partial_x$	$A_{3.6}$
9	$\lambda u ^{n+1}, \lambda \neq 0, n \neq 0, -1$	$t\partial_t - \frac{1}{n}u\partial_u, x\partial_x - \frac{1}{n}u\partial_u, \partial_t, \partial_x$	$A_{2.2} \oplus A_{2.2}$

THEOREM 11. *There are at most seven inequivalent classes of nonlinear equation (2.7) invariant under one-dimensional Lie algebras.*

Below we give the full list of the invariant equations and the corresponding invariance algebras.

$$A_1^1 = \langle t\partial_t + x\partial_x \rangle : F = t^{-2}G(\xi, u, \omega), \xi = tx^{-1}, \omega = xu_x;$$

$$A_1^2 = \langle \partial_t + k\partial_x \rangle (k > 0) : F = G(\eta, u, u_x), \eta = x - kt;$$

$$A_1^3 = \langle \partial_x \rangle : F = G(t, u, u_x);$$

$$A_1^4 = \langle \partial_t \rangle : F = G(x, u, u_x);$$

$$A_1^5 = \langle \partial_t + f(x)u\partial_u \rangle (f \neq 0) :$$

$$F = -tf''u + t^2(f')^2u - 2tf'u_x + e^{tf}G(x, v, \omega),$$

$$v = e^{-tf}u, \omega = u^{-1}u_x - f'f^{-1} \ln |u|;$$

$$A_1^6 = \langle f(x)u\partial_u \rangle (f \neq 0) : F = -f^{-1}f''u \ln |u|$$

$$-2f^{-1}f'u_x \ln |u| + f^{-2}(f')^2u \ln^2 |u| + uG(t, x, \omega),$$

$$\omega = u^{-1}u_x - f'f^{-1} \ln |u|;$$

$$A_1^7 = \langle f(t, x)\partial_u \rangle (f \neq 0) : F = f^{-1}(f_{tt} - f_{xx})u + G(t, x, \omega),$$

$$\omega = u_x - f^{-1}f_x u.$$

Note that if the functions F and G are arbitrary, then the given algebras are maximal (in Lie’s sense) symmetry algebras of the respective equations.

THEOREM 12. *An equation of the form (2.7) cannot admit Lie algebra which has a subalgebra having a non-trivial Levi factor.*

With account of the above facts we conclude that nonlinear equation (2.7) admit a symmetry algebra of the dimension higher than one only if the latter is a solvable real Lie algebra. That is why, we turn to classifying Equation (2.7) whose invariance algebras are two-dimensional solvable Lie algebras.

As the calculations are similar to those performed in the third section, we present the final result only. Namely, we give the form of invariant equations and the corresponding realizations of the two-dimensional invariance algebras.

I. A_{2,1}-invariant equations

$$A_{2,1}^1 = \langle t\partial_t + x\partial_x, u\partial_u \rangle : F = x^{-2}uG(\xi, \omega),$$

$$\xi = tx^{-1}, \omega = u^{-1}xu_x;$$

$$A_{2,1}^2 = \langle t\partial_t + x\partial_x, \sigma(\xi)\partial_u \rangle (\sigma \neq 0, \xi = tx^{-1}) :$$

$$F = x^{-2}[\sigma^{-1}((1 - \xi^2)\sigma'' - 2\xi\sigma')u + G(\xi, \omega)],$$

$$\omega = \xi\sigma'u + \sigma xu_x;$$

$$A_{2,1}^3 = \langle \partial_t + k\partial_x, u\partial_u \rangle (k > 0) : F = uG(\eta, \omega),$$

$$\eta = x - kt, \omega = u^{-1}u_x;$$

$$A_{2,1}^4 = \langle \partial_t + k\partial_x, \varphi(\eta)\partial_u \rangle (k > 0, \eta = x - kt, \varphi \neq 0) :$$

$$F = (k^2 - 1)\varphi''\varphi^{-1}u + G(\eta, \omega), \omega = \varphi u_x - \varphi'u;$$

$$A_{2,1}^5 = \langle \partial_t + k\partial_x, \partial_x + u\partial_u \rangle (k > 0) :$$

$$F = e^\eta G(\omega, v), \eta = x - kt, \omega = ue^{-\eta}, v = u^{-1}u_x;$$

$$A_{2,1}^6 = \langle \partial_t, \partial_x \rangle : F = G(u, u_x);$$

$$A_{2,1}^7 = \langle \partial_x, u\partial_u \rangle : F = uG(t, \omega), \omega = u^{-1}u_x;$$

$$A_{2,1}^8 = \langle \partial_x, \varphi(t)\partial_u \rangle (\varphi \neq 0) :$$

$$F = \varphi^{-1}\varphi''u + G(t, u_x);$$

$$A_{2,1}^9 = \langle \partial_t, \partial_u \rangle : F = G(x, u_x);$$

$$A_{2,1}^{10} = \langle \partial_t, f(x)u\partial_u \rangle (f \neq 0) :$$

$$F = -u^{-1}u_x^2 + uG(x, \omega);$$

$$\omega = u^{-1}u_x - f'f^{-1}\ln|u|;$$

$$\begin{aligned}
A_{2.1}^{11} &= \langle \partial_t + f(x)u\partial_u, g(x)u\partial_u \rangle (\delta = f^{-1}f' - g^{-1}g' \neq 0) : \\
&F = -g^{-1}g''u \ln |u| - 2g^{-1}g'u_x \ln |u| \\
&\quad + g^{-2}(g')^2u \ln^2 |u| - 2f\delta tu_x + 2f\delta g'g^{-1}tu \ln |u| \\
&\quad + f^2\delta^2t^2u + f(g^{-1}g'' - f^{-1}f'')tu + uG(x, \omega), \\
&\omega = u^{-1}u_x - g'g^{-1} \ln |u| - tf\delta; \\
A_{2.1}^{12} &= \langle \partial_t + f(x)u\partial_u, e^{tf}\partial_u \rangle (f \neq 0) : \\
&F = [f^2 - tf'' + t^2(f')^2]u - 2tf'u_x + e^{tf}G(x, \omega), \\
&\omega = e^{-tf}(u_x - tf'u); \\
A_{2.1}^{13} &= \langle f(x)u\partial_u, g(x)u\partial_u \rangle (\delta = f'g - g'f \neq 0) : \\
&F = -u^{-1}u_x^2 - \delta^{-1}\delta'u_x \\
&\quad + \delta^{-1}[f''g' - g''f']u \ln |u| + uG(t, x); \\
A_{2.1}^{14} &= \langle \varphi(t)\partial_u, \psi(t)\partial_u \rangle (\varphi'\psi - \varphi\psi' \neq 0) : \\
&F = \varphi^{-1}\varphi''u + G(t, x, u_x), \varphi''\psi - \varphi\psi'' = 0.
\end{aligned}$$

II. $A_{2.2}$ -invariant equations

$$\begin{aligned}
A_{2.2}^1 &= \langle t\partial_t + x\partial_x, xu\partial_u \rangle : F = x^{-2}u \ln^2 |u| \\
&\quad - 2x^{-1}u_x \ln |u| + t^{-2}uG(\xi, \omega), \xi = tx^{-1}; \\
&\omega = xu^{-1}u_x - \ln |u|; \\
A_{2.2}^2 &= \langle t\partial_t + x\partial_x, t\varphi(\xi)\partial_u \rangle (\varphi \neq 0, \xi = tx^{-1}) : \\
&F = t^{-2}(1 - \xi^2)\varphi^{-1}\xi(2\varphi' + \xi\varphi'')u + t^{-2}G(\xi, \omega), \\
&\omega = x\varphi u_x + \xi\varphi'u; \\
A_{2.2}^3 &= \langle \partial_t + k\partial_x, \exp(k^{-1}x)u\partial_u \rangle (k > 0) : \\
&F = k^{-2}u \ln^2 |u| - 2k^{-1}u_x \ln |u| - k^{-2}u \ln |u| \\
&\quad + uG(\eta, \omega), \eta = x - kt, \omega = u^{-1}u_x - k^{-1} \ln |u|; \\
A_{2.2}^4 &= \langle \partial_t + k\partial_x, e^t\varphi(\eta)\partial_u \rangle (\eta = x - kt, k > 0, \varphi \neq 0) : \\
&F = ((k^2 - 1)\varphi''\varphi^{-1} - 2k\varphi'\varphi^{-1} + 1)u + G(\eta, \omega), \\
&\omega = \varphi u_x - \varphi'u, \varphi' = \frac{d\varphi}{d\eta};
\end{aligned}$$

$$\begin{aligned}
A_{2,2}^5 &= \langle -t\partial_t - x\partial_x, \partial_t + k\partial_x \rangle (k > 0) : \\
&F = \eta^{-2}G(u, \omega), \quad \eta = x - kt, \quad \omega = u_x\eta; \\
A_{2,2}^6 &= \langle -t\partial_t - x\partial_x + mu\partial_u, \partial_t + k\partial_x \rangle (k > 0, m \neq 0) : \\
&F = |\eta|^{-2-m}G(v, \omega), \quad \eta = x - kt, \\
&\omega = u|\eta|^m, \quad v = u_x|\eta|^{m+1}; \\
A_{2,2}^7 &= \langle \partial_x, e^x u \partial_u \rangle : \quad F = u \ln^2 |u| - u \ln |u| - 2u_x \ln |u| \\
&\quad + uG(t, \omega), \quad \omega = u^{-1}u_x - \ln |u|; \\
A_{2,2}^8 &= \langle \partial_x, e^x \varphi(t) \partial_u \rangle (\varphi \neq 0) : \\
&F = (\varphi^{-1} \varphi'' - 1)u + G(t, \omega), \quad \omega = u_x - u; \\
A_{2,2}^9 &= \langle -t\partial_t - x\partial_x, \partial_x \rangle : \quad F = t^{-2}G(u, tu_x); \\
A_{2,2}^{10} &= \langle -t\partial_t - x\partial_x + ku\partial_u, \partial_x \rangle, (k \neq 0) : \\
&F = |t|^{-2-k}G(v, \omega), \quad v = |t|^k u, \quad \omega = |t|^{k+1} u_x; \\
A_{2,2}^{11} &= \langle \partial_t, e^t \partial_u \rangle : \quad F = u + G(x, u_x); \\
A_{2,2}^{12} &= \langle -t\partial_t - x\partial_x, \partial_t \rangle : \quad F = x^{-2}G(u, \omega), \quad \omega = xu_x; \\
A_{2,2}^{13} &= \langle \partial_t + f(x)u\partial_u, e^{(1+f)t} \partial_u \rangle (f \neq 0) : \\
&F = -(tf'' - t^2(f')^2 - (1 + f^2))u - 2tf'u_x \\
&\quad + e^{tf}G(x, \omega), \quad \omega = e^{-tf}(u_x - f'(t + f^{-1})u); \\
A_{2,2}^{14} &= \langle -t\partial_t - x\partial_x, \partial_t + kx^{-1}u\partial_u \rangle (k > 0); \\
&F = -2ktx^{-3}u + k^2t^2x^{-4}u + 2ktx^{-2}u_x \\
&\quad + x^{-2} \exp(ktx^{-1})G(v, \omega), \quad v = \exp(-kx^{-1}t)u, \\
&\omega = xu^{-1}u_x + \ln |u|; \\
A_{2,2}^{15} &= \langle k(t\partial_t + x\partial_x), |x|^{k-1}u\partial_u \rangle (k \neq 0, 1) : \\
&F = -k^{-2}(1-k)x^{-2}u \ln |u| - 2k^{-1}x^{-1}u_x \ln |u| \\
&\quad + k^{-2}x^{-2}u \ln^2 |u| + x^{-2}uG(v, \omega), \\
&v = tx^{-1}, \quad \omega = xu^{-1}u_x - k^{-1} \ln |u|; \\
A_{2,2}^{16} &= \langle k(t\partial_t + x\partial_x), |t|^{k-1}\varphi(\xi)\partial_u \rangle (k \neq 0, 1, \varphi \neq 0, \\
&\xi = tx^{-1}) : F = [k^{-1}(k^{-1} - 1) + 2\xi(k^{-1} - \xi^2)\varphi^{-1}\varphi' \\
&\quad + \xi^2(1 - \xi)^2\varphi^{-1}\varphi'']t^{-2}u + t^{-2}G(\xi, \omega), \\
&\omega = x\varphi u_x + \xi\varphi'u.
\end{aligned}$$

In the above formulas G stands for an arbitrary smooth function. As customary the prime denotes derivative of a function of one variable.

7.1. Group Classification of Equation

$$u_{tt} = u_{xx} - u^{-1}u_x^2 + A(x)u_x + B(x)u \ln |u| + uD(t, x)$$

Before analyzing Equation (2.7) admitting algebras of the dimension higher than two we perform group classification of the equation

$$u_{tt} = u_{xx} - u^{-1}u_x^2 + A(x)u_x + B(x)u \ln |u| + uD(t, x). \quad (6.1)$$

Here $A(x)$, $B(x)$, $D(t, x)$ are arbitrary smooth functions. Note that class of PDEs (6.1) contains $A_{2,1}^{13}$ -invariant equation. What is more important this class contains a major part of equations of the form (2.7), maximal symmetry algebras of which have dimension three or four. We make use of this fact to simplify group classification of Equation (2.7).

LEMMA 5. *If the functions A , B and D are arbitrary, then maximal invariance algebra of PDE (6.1) is the two-dimensional Lie algebra equivalent to $A_{2,1}^{13}$ and (6.1) reduces to $A_{2,1}^{13}$ -invariant equation. Next if the maximal symmetry algebra of an equation of the form (6.1) is three-dimensional (denote it as A_3), then this equation is equivalent to one of the following ones:*

- I. $A_3 \sim A_{3,1}$, $A_3 = \langle \partial_t, f(x)u\partial_u, \varphi(x)u\partial_u \rangle$,
 $A = -\sigma^{-1}\sigma'$, $B = \sigma^{-1}\rho$, $D = 0$, $\sigma = f'\varphi - f\varphi' \neq 0$,
 $\rho = \varphi'f'' - \varphi''f'$;
- II. $A_3 \sim A_{3,1}$, $A_3 = \langle f(x)u\partial_u, \varphi(x)u\partial_u, \partial_t + \psi(x)u\partial_u \rangle$,
 $A = -\sigma^{-1}\sigma'$, $B = \sigma^{-1}\rho$,
 $D = t\sigma^{-1}[\sigma'\psi' - \psi\rho - \sigma\psi'']$,
 $\sigma = f'\varphi - \varphi'f \neq 0$, $\rho = f''\varphi' - \varphi''f'$,
 $f'\psi - f\psi' \neq 0$, $\varphi'\psi - \varphi\psi' \neq 0$;
- III. $D = x^{-2}G(\xi)$, $\xi = tx^{-1}$, $G \neq 0$:
 - 1) $A_3 \sim A_{3,2}$, $A_3 = \langle t\partial_t + x\partial_x, u\partial_u, |x|^{1-n}u\partial_u \rangle$,
 $A = nx^{-1}$ ($n \neq 1$), $B = 0$;
 - 2) $A_3 \sim A_{3,3}$, $A_3 = \langle t\partial_t + x\partial_x, u\partial_u, u \ln |x|\partial_u \rangle$, $A = x^{-1}$, $B = 0$;
 - 3) $A_3 \sim A_{3,4}$, $A_3 = \langle t\partial_t + x\partial_x, \sqrt{|x}|u\partial_u, \sqrt{|x}| \ln |x|u\partial_u \rangle$,
 $A = 0$, $B = \frac{1}{4}x^{-2}$;
 - 4) $A_3 \sim A_{3,9}$, $A_3 = \langle t\partial_t + x\partial_x, \sqrt{|x}| \cos(\frac{1}{2}\beta \ln |x|)u\partial_u,$
 $\sqrt{|x}| \sin(\frac{1}{2}\beta \ln |x|)u\partial_u \rangle$, $A = 0$, $B = mx^{-2}$, $m > \frac{1}{4}$, $\beta = \sqrt{4m-1}$;
 - 5) $A_3 \sim A_{3,7}$, $A_3 = \langle t\partial_t + x\partial_x, (\sqrt{|x|})^{1+\beta}u\partial_u, (\sqrt{|x|})^{1-\beta}u\partial_u \rangle$,
 $A = 0$, $B = mx^{-2}$, $m < \frac{1}{4}$, $m \neq 0$, $\beta = \sqrt{1-4m}$;
 - 6) $A_3 \sim A_{3,8}$, $A_3 = \langle t\partial_t + x\partial_x, \cos(\sqrt{m} \ln |x|)u\partial_u,$
 $\sin(\sqrt{m} \ln |x|)u\partial_u \rangle$, $A = x^{-1}$, $B = mx^{-2}$, $m > 0$;
 - 7) $A_3 \sim A_{3,6}$, $A_3 = \langle t\partial_t + x\partial_x, |x|^{\sqrt{|m|}}u\partial_u, |x|^{-\sqrt{|m|}}u\partial_u \rangle$,
 $A = x^{-1}$, $B = mx^{-2}$, $m < 0$;
 - 8) $A_3 \sim A_{3,4}$, $A_3 = \langle t\partial_t + x\partial_x, (\sqrt{|x|})^{1-n}u\partial_u, (\sqrt{|x|})^{1-n}$
 $\times \ln |x|u\partial_u \rangle$, $A = nx^{-1}$ ($n \neq 0, 1$), $B = \frac{1}{4}(n-1)^2x^{-2}$;

- 9) $A_3 \sim A_{3.9}, A_3 = \langle t\partial_t + x\partial_x, (\sqrt{|x|})^{1-n} \cos(\frac{1}{2}\beta \ln|x|)u\partial_u, (\sqrt{|x|})^{1-n} \sin(\frac{1}{2}\beta \ln|x|)u\partial_u \rangle, A = nx^{-1} (n \neq 0, 1), B = mx^{-2} (m > \frac{1}{4}(n-1)^2), \beta = \sqrt{4m - (n-1)^2};$
- 10) $A_3 \sim A_{3.7}, A_3 = \langle t\partial_t + x\partial_x, (\sqrt{|x|})^{1-\beta-n}u\partial_u, (\sqrt{|x|})^{1-n+\beta} \times u\partial_u \rangle, A = nx^{-1} (n \neq 0, 1), B = mx^{-2} (m < \frac{1}{4}(n-1)^2, m \neq 0), \beta = \sqrt{(n-1)^2 - 4m}.$

IV. $D = G(t),$

- 1) $A_3 \sim A_{3.3}, A_3 = \langle \partial_x, u\partial_u, xu\partial_u \rangle, A = B = 0;$
- 2) $A_3 \sim A_{3.2}, A_3 = \langle \partial_x, u\partial_u, e^x u\partial_u \rangle, A = -1, B = 0;$
- 3) $A_3 \sim A_{3.8}, A_3 = \langle \partial_x, \cos(x)u\partial_u, \sin(x)u\partial_u \rangle, A = 0, B = 1;$
- 4) $A_3 \sim A_{3.6}, A_3 = \langle \partial_x, e^x u\partial_u, e^{-x} u\partial_u \rangle, A = 0, B = -1;$
- 5) $A_3 \sim A_{3.4}, A_3 = \langle \partial_x, \exp(\frac{1}{2}x) u\partial_u, \exp(\frac{1}{2}x) xu\partial_u \rangle, A = -1, B = \frac{1}{4};$
- 6) $A_3 \sim A_{3.7}, A_3 = \langle \partial_x, \exp(\frac{1}{2}(1 + \beta)x) u\partial_u, \exp(\frac{1}{2}(1 - \beta)x) u\partial_u \rangle, A = -1, B = m (m < \frac{1}{4}), m \neq 0, \beta = \sqrt{1 - 4m};$
- 7) $A_3 \sim A_{3.9}, A_3 = \langle \partial_x, \exp(\frac{1}{2}x) \cos(\frac{1}{2}\beta x)u\partial_u, \exp(\frac{1}{2}x) \sin(\frac{1}{2}\beta x)u\partial_u \rangle, A = -1, B = m (m > \frac{1}{4}), \beta = \sqrt{4m - 1};$

V. $D = G(\eta), \eta = x - kt, k > 0,$

- 1) $A_3 \sim A_{3.3}, A_3 = \langle \partial_t + k\partial_x, u\partial_u, xu\partial_u \rangle, A = B = 0;$
- 2) $A_3 \sim A_{3.2}, A_3 = \langle \partial_t + k\partial_x, u\partial_u, e^x u\partial_u \rangle, A = -1, B = 0;$
- 3) $A_3 \sim A_{3.8}, A_3 = \langle \partial_t + k\partial_x, \cos(x)u\partial_u, \sin(x)u\partial_u \rangle, A = 0, B = 1;$
- 4) $A_3 \sim A_{3.6}, A_3 = \langle \partial_t + k\partial_x, e^x u\partial_u, e^{-x} u\partial_u \rangle, A = n, B = -1;$
- 5) $A_3 \sim A_{3.4}, A_3 = \langle \partial_t + k\partial_x, \exp(\frac{1}{2}x) u\partial_u, \exp(\frac{1}{2}x) xu\partial_u \rangle, A = -1, B = \frac{1}{4};$
- 6) $A_3 \sim A_{3.7}, A_3 = \langle \partial_t + k\partial_x, \exp(\frac{1}{2}(1 + \beta)x) u\partial_u, \exp(\frac{1}{2}(1 - \beta)x) u\partial_u \rangle, A = -1, B = m (m < \frac{1}{4}), m \neq 0, \beta = \sqrt{1 - 4m};$
- 7) $A_3 \sim A_{3.9}, A_3 = \langle \partial_t + k\partial_x, \exp(\frac{1}{2}x) \cos(\frac{1}{2}\beta x)u\partial_u, \exp(\frac{1}{2}x) \sin(\frac{1}{2}\beta x)u\partial_u \rangle, A = -1, B = m (m > \frac{1}{4}), \beta = \sqrt{4m - 1}.$

Proof. Inserting the expression

$$F = -u^{-1}u_x^2 + A(x)u_x + B(x)u \ln|u| + uD(t, x)$$

into classifying Equation (2.5) we get the system of determining equations for the functions $h(x), r(t, x)$ and constants $\lambda, \lambda_1, \lambda_2$:

$$\begin{aligned} r &= 0, \quad (\lambda x + \lambda_2)A' + \lambda A = 0, \\ (\lambda x + \lambda_2)B' + 2\lambda B &= 0, \quad h'' + Ah' + Bh = -(\lambda t + \lambda_1)D_t - (\lambda x + \lambda_2)D_x - 2\lambda D. \end{aligned} \tag{6.2}$$

First, we consider the case of arbitrary functions A, B, D . The left-hand side of the fourth equation from Equation (6.2) depends on x only. Since D is arbitrary, the relation $D_t \neq 0$ holds. Hence it immediately follows that the constants $\lambda, \lambda_1, \lambda_2$ must be equal to zero. As a consequence, the fourth equation becomes the linear ordinary differential equation for the function $h = h(x)$

$$h'' + Ah' + Bh = 0. \quad (6.3)$$

The general solution of the above equation reads as

$$h = C_1 f(x) + C_2 \varphi(x), \quad C_1, C_2 \in \mathbb{R},$$

($f(x), \varphi(x)$) being the fundamental system of solutions of the equation

$$y'' + Ay' + By = 0, \quad y = y(x). \quad (6.4)$$

Inserting the expression for h into (6.3) yields

$$A = -\sigma^{-1}\sigma', \quad B = \sigma^{-1}(\varphi' f'' - f' \varphi''),$$

where $\sigma = \varphi f' - \varphi' f \neq 0$, which proves the first part of the lemma.

Suppose now that $D = 0$. If at least one of the functions A or B is arbitrary, then $\lambda = \lambda_2 = 0$ and the function h is a solution of (6.3). This completes the proof of the case I of the second part of the lemma statement.

Provided functions A and B are not arbitrary, it follows from the second and third equations of (6.2) that one of the following relations

- 1) $A = B = 0$;
- 2) $A = n, B = m, m, n \in \mathbb{R}, |n| + |m| \neq 0$;
- 3) $A = nx^{-1}, B = mx^{-2}, m, n \in \mathbb{R}, |n| + |m| \neq 0$ (6.5)

holds. With these conditions the maximal invariance algebra of (6.1) has the dimension higher than three. Consequently, without any loss of generality we can suggest that $D \neq 0$. Integrating the equation

$$(\lambda t + \lambda_1)D_t + (\lambda x + \lambda_2)D_x + 2\lambda D = H(x),$$

under $D \neq 0$ yields the following (inequivalent) expressions for the function $D(t, x)$:

$$\begin{aligned} D &= x^{-2}G(\xi) + x^{-2} \int x H(x) dx, \quad \xi = tx^{-1}; \\ D &= G(\eta) + k^{-1} \int H(x) dx, \quad \eta = x - kt, \quad k > 0; \\ D &= G(t) + \int H(x) dx, \\ D &= tH(x) + \tilde{H}(x). \end{aligned} \quad (6.6)$$

The change of variables

$$\bar{t} = t, \quad \bar{x} = x, \quad u = \theta(x)v(\bar{t}, \bar{x}), \quad \theta \neq 0, \quad (6.7)$$

where θ is a solution of equation

$$\theta^{-1}\theta'' - \theta^{-2}(\theta')^2 + A\theta^{-1}\theta' + B \ln |\theta| + \Lambda(x) = 0,$$

preserves the form of Equation (6.1). We use this fact to simplify the form of the function D

$$\begin{aligned}
 D &= x^{-2}G(\xi), \quad \xi = tx^{-1}; \\
 D &= G(\eta), \quad \eta = x - kt, \quad k > 0; \\
 D &= G(t), \\
 D &= tH(x).
 \end{aligned}
 \tag{6.8}$$

If the function D is given by the one of the first three expressions, then $H(x) \equiv 0$ and h satisfies (6.3).

Given the condition $D = tH(x)$, we have

$$h'' + Ah' + Bh = -\lambda_1 H, \quad (\lambda x + \lambda_2)H' + 3\lambda H = 0.$$

So that the maximal invariance algebra of the corresponding Equation (6.1) is three-dimensional iff $\lambda = \lambda_2 = 0$, which yields the case II of the second part of the lemma statement.

Turn now to the case when $D = x^{-2}G(\xi)$, $\xi = tx^{-1}$. Then the (non-vanishing identically) function G obeys the equation

$$(\lambda_2 \xi - \lambda_1)G' + 2\lambda_2 G = 0. \tag{6.9}$$

If G is an arbitrary function, then $\lambda_1 = \lambda_2 = 0$. In addition we have $\lambda \neq 0$ (otherwise the maximal invariance algebra is two-dimensional). Hence we get

$$xA' + a = 0, \quad xB' + 2B = 0.$$

Consequently, the functions A and B are given by either first or third formula from (6.5). Analyzing these expressions yields 10 cases of the case III of the second part of the lemma statement.

If the function G is not arbitrary, then integrating Equation (6.9) we get

$$\begin{aligned}
 G &= p, \quad p \in \mathbb{R}, \quad p \neq 0; \\
 G &= p(\xi - q)^{-2}, \quad p \neq 0, \quad q \geq 0.
 \end{aligned}$$

Given the condition $G = p$ the parameter λ_2 vanishes. Hence in view of the requirement for the maximal algebra to be three-dimensional it follows that λ vanishes as well. This yields the case when A and B in (6.1) are arbitrary functions (the case I of the second part of lemma statement). If $G = p(\xi - q)^{-2}$, $p \neq 0$, then $\lambda_1 = \lambda_2 q$. Hence we conclude that the maximal invariance algebra of the corresponding Equation (6.1) is three-dimensional iff the functions A, B are given by formulas 3) from (6.5) (which implies that $\lambda_1 = \lambda_2 = 0$) and we get the case III of the second part of the lemma statement. Next if A, B are given by formulas 2 from (6.5) (which implies that $\lambda = 0$, $D = p(t - qx)^{-2}$), then we arrive at the case IV (under $q = 0$) or the case V (under $q > 0$) of the second part of the lemma statement.

Turn now to the case $D = G(\eta)$, $\eta = x - kt$, $k > 0$. If these relations hold, then

$$\lambda(\eta G' + 2G) + (\lambda_2 - k\lambda_1)G' = 0.$$

Hence it follows that if G is an arbitrary function of η , then $\lambda = 0$, $\lambda_2 = k\lambda_1$. That is why, the maximal invariance algebra of (6.1) is three-dimensional iff either $A = B = 0$ or the functions A, B are given by formulas 2 from (6.5). Thus we have derived all equations listed in the case V of the second part of the lemma statement.

The cases when either $G = p$ ($p \neq 0$) or $G = p\eta^{-2}$ ($p \neq 0$) yield no new invariant equation (6.5).

Consider now the last possible case $D = G(t)$. If this is the case, then the equation

$$(\lambda t + \lambda_1)G' + 2\lambda G = 0$$

holds. Hence it follows that if G is an arbitrary function, then $\lambda = \lambda_1 = 0$. Consequently, the maximal invariance algebra of Equation (6.1) is three-dimensional iff A, B are given by the formula 2 from (6.5). This yields all invariant equations from the case IV of the second part of the assertion of the lemma. If either of the relations $G = p$ ($p \neq 0$) or $G = pt^{-2}$ ($p \neq 0$) hold, then no new invariant equations admitting three-dimensional maximal invariance algebras can be obtained.

To complete the proof of the lemma we need to establish non-equivalence of the obtained invariant equations. To this end it suffices to prove that there are no transformations from the group \mathcal{E} , reducing their invariance algebras one into another.

As we have already mentioned in Section 3 there exist nine non-isomorphic three-dimensional solvable Lie algebras $A_{3,i} = \langle e_1, e_2, e_3 \rangle$ ($i = 1, 2, \dots, 9$). We analyze in some detail the case of the algebra $A_{3,3}$. The list of invariant equations and algebras contains three algebras which are isomorphic to $A_{3,3}$, namely,

$$L_1 = \langle t\partial_t + x\partial_x, u\partial_u, u \ln |x|\partial_u \rangle;$$

$$L_2 = \langle \partial_x, u\partial_u, xu\partial_u \rangle;$$

$$L_2 = \langle \partial_t + k\partial_x, u\partial_u, xu\partial_u \rangle \quad (k > 0).$$

Denote the basis elements of the algebra L_2 as e_1, e_2, e_3 . Suppose that there is a transformation φ from the group \mathcal{E} transforming L_2 into L_3 . In other words we suppose that there exist constants $\alpha_i, \beta_i, \delta_i \in \mathbb{R}$ ($i = 1, 2, 3$) such that the relations

$$\varphi(e_1) = \sum_{i=1}^3 \alpha_i \tilde{e}_i, \quad \varphi(e_2) = \sum_{i=1}^3 \beta_i \tilde{e}_i, \quad \varphi(e_3) = \sum_{i=1}^3 \delta_i \tilde{e}_i$$

and

$$\Delta = \begin{vmatrix} \alpha_1 & \alpha_2 & \alpha_3 \\ \beta_1 & \beta_2 & \beta_3 \\ \delta_1 & \delta_2 & \delta_3 \end{vmatrix} \neq 0$$

hold. In the above formulas $\tilde{e}_1 = \partial_t + k\partial_x$, $\tilde{e}_2 = v\partial_v$, $\tilde{e}_3 = \bar{x}v\partial_v$. Equating the coefficients of linearly independent operators $\partial_{\bar{t}}, \partial_{\bar{x}}, \partial_v$ yields that $\alpha_1 = \beta_1 = \delta_1 = 0$. Hence we get the contradictory equation $\Delta = 0$. This means that realizations L_2 and L_3 are non-isomorphic. Analogously, we prove that L_1 and L_2 (as well as L_1 and L_3) are non-isomorphic.

The remaining algebras are considered in a similar way. The Lemma is proved.

In what follows we will use the results on classification of abstract four-dimensional solvable real Lie algebras $A_4 = \langle e_1, e_2, e_3, e_4 \rangle$ [42, 43]. There are 10 decomposable

$$4A_1 = 3A_1 \oplus A_1 = A_{3,1} \oplus A_1, \quad A_{2,2} \oplus 2A_1 = A_{2,2} \oplus A_{2,1} = A_{3,2} \oplus A_1,$$

$$2A_{2,2} = A_{2,2} \oplus A_{2,2}, \quad A_{3,i} \oplus A_1 \quad (i = 3, 4, \dots, 9);$$

and 10 non-decomposable four-dimensional solvable real Lie algebras (note that we give below non-zero commutation relations only).

- $A_{4.1} : [e_2, e_4] = e_1, [e_3, e_4] = e_2;$
- $A_{4.2} : [e_1, e_4] = qe_1, [e_2, e_4] = e_2,$
 $[e_3, e_4] = e_2 + e_3, q \neq 0;$
- $A_{4.3} : [e_1, e_4] = e_1, [e_3, e_4] = e_2;$
- $A_{4.4} : [e_1, e_4] = e_1, [e_2, e_4] = e_1 + e_2,$
 $[e_3, e_4] = e_2 + e_3;$
- $A_{4.5} : [e_1, e_4] = e_1, [e_2, e_4] = qe_2,$
 $[e_3, e_4] = pe_3, -1 \leq p \leq q \leq 1, p \cdot q \neq 0;$
- $A_{4.6} : [e_1, e_4] = qe_1, [e_2, e_4] = pe_2 - e_3,$
 $[e_3, e_4] = e_2 + pe_3, q \neq 0, p \geq 0;$
- $A_{4.7} : [e_2, e_3] = e_1, [e_1, e_4] = 2e_1,$
 $[e_2, e_4] = e_2, [e_3, e_4] = e_2 + e_3;$
- $A_{4.8} : [e_2, e_3] = e_1, [e_1, e_4] = (1 + q)e_1,$
 $[e_2, e_4] = e_2, [e_3, e_4] = qe_3, |q| \leq 1;$
- $A_{4.9} : [e_2, e_3] = e_1, [e_1, e_4] = 2qe_1,$
 $[e_2, e_4] = qe_2 - e_3, [e_3, e_4] = e_2 + qe_3, q \geq 0;$
- $A_{4.10} : [e_1, e_3] = e_1, [e_2, e_3] = e_2,$
 $[e_1, e_4] = -e_2, [e_2, e_4] = e_1.$

THEOREM 13. Equation $u_{tt} = u_{xx} - u^{-1}u_x^2$ has the widest symmetry group amongst equations of the form (6.1). Its maximal invariance algebra is the five-dimensional Lie algebra

$$A_5^1 = \langle \partial_t, \partial_x, t\partial_t + x\partial_x, xu\partial_u, u\partial_u \rangle.$$

There are no equations of the form (6.1) which are inequivalent to the above equation and admit invariance algebra of the dimension higher than four. Below we give all inequivalent equation (6.1) admitting four-dimensional together with their symmetry algebras.

I. $D = 0,$

- 1) $A_4 \sim A_{3.6} \oplus A_1, A_4 = \langle \partial_t, \partial_x, u \operatorname{ch}(\beta x)\partial_u, u \sinh(\beta x)\partial_u \rangle,$
 $A = 0, B = -\beta^2, \beta \neq 0;$
- 2) $A_4 \sim A_{3.8} \oplus A_1, A_4 = \langle \partial_t, \partial_x, u \cos(\beta x)\partial_u, u \sin(\beta x)\partial_u \rangle,$
 $A = 0, B = \beta^2, \beta \neq 0;$
- 3) $A_4 \sim A_{2.1} \oplus A_{2.2}, A_4 = \langle \partial_t, \partial_x, u\partial_u, e^{-x}u\partial_u \rangle, A = 1, B = 0;$
- 4) $A_4 \sim A_{3.4} \oplus A_1, A_4 = \langle \partial_t, \partial_x, e^{-x}u\partial_u, xe^{-x}u\partial_u \rangle, A = 2, B = 1;$

- 5) $A_4 \sim A_{3,9} \oplus A_1$, $A_4 = \langle \partial_t, \partial_x, ue^{-x} \cos(\beta x) \partial_u, ue^{-x} \sin(\beta x) \partial_u \rangle$,
 $A = 2$, $B = m$, $m > 1$, $\beta = \sqrt{m-1}$;
- 6) $A_4 \sim A_{3,7} \oplus A_1$, $A_4 = \langle \partial_t, \partial_x, ue^{-x} \operatorname{ch}(\beta x) \partial_u, ue^{-x} \sinh(\beta x) \partial_u \rangle$,
 $A = 2$, $B = m$, $m > 1$, $m \neq 0$, $\beta = \sqrt{1-m}$;
- 7) $A_4 \sim A_{4,2}$, $A_4 = \langle \partial_t, t \partial_t + x \partial_x, \sqrt{|x|} u \partial_u, u \sqrt{|x|} \ln |x| \partial_u \rangle$,
 $A = 0$, $B = \frac{1}{4} x^{-2}$;
- 8) $A_4 \sim A_{4,5}$, $A_4 = \langle \partial_t, t \partial_t + x \partial_x, |x|^{\frac{1}{2}+\beta} u \partial_u, |x|^{\frac{1}{2}-\beta} u \partial_u \rangle$, $A = 0$,
 $B = mx^{-2}$, $m < \frac{1}{4}$, $m \neq 0$, $\beta = \sqrt{\frac{1}{4} - m}$;
- 9) $A_4 \sim A_{4,6}$, $A_4 = \langle \partial_t, t \partial_t + x \partial_x, \sqrt{|x|} \cos(\beta \ln |x|) u \partial_u, \sqrt{|x|} \sin(\beta \ln |x|) u \partial_u \rangle$, $A = 0$, $B = mx^{-2}$, $m > \frac{1}{4}$, $\beta = \sqrt{m - \frac{1}{4}}$;
- 10) $A_4 \sim A_{4,3}$, $A_4 = \langle \partial_t, t \partial_t + x \partial_x, u \ln |x| \partial_u, u \partial_u \rangle$, $A = x^{-1}$, $B = 0$;
- 11) $A_4 \sim A_{3,7} \oplus A_1$, $A_4 = \langle \partial_t, t \partial_t + x \partial_x, |x|^{1-n} u \partial_u, u \partial_u \rangle$,
 $A = nx^{-1}$, $B = 0$, $n \neq 0, 1$;
- 12) $A_4 \sim A_{4,5}$, $A_4 = \langle \partial_t, t \partial_t + x \partial_x, |x|^{\frac{1}{2}(1-n)} u \partial_u, |x|^{\frac{1}{2}(1-n)} u \ln |x| \partial_u \rangle$,
 $A = nx^{-1}$, $B = \frac{1}{4}(n-1)^2 x^{-2}$, $n \neq 0, 1$;
- 13) $A_4 \sim A_{4,5}$, $A_4 = \langle \partial_t, t \partial_t + x \partial_x, |x|^{\frac{1}{2}(1-n+\beta)} u \partial_u, |x|^{\frac{1}{2}(1-n-\beta)} u \partial_u \rangle$,
 $A = nx^{-1}$, $B = mx^{-2}$, $m < \frac{1}{4}(n-1)^2$, $m \neq 0$, $n \neq 0$, $\beta = \sqrt{(n-1)^2 - 4m}$;
- 14) $A_4 \sim A_{4,6}$, $A_4 = \langle \partial_t, t \partial_t + x \partial_x, |x|^{\frac{1}{2}(1-n)} \cos(\beta \ln |x|) u \partial_u, |x|^{\frac{1}{2}(1-n)} \sin(\beta \ln |x|) u \partial_u \rangle$,
 $A = nx^{-1}$, $B = mx^{-2}$, $m \neq 0$, $n \neq 0$, $m > \frac{1}{4}(n-1)^2$, $\beta = \sqrt{m - \frac{1}{4}(n-1)^2}$;

II. $D = ktx^{-3}$, $k > 0$,

- 1) $A_4 \sim A_{4,1}$, $A_4 = \langle \partial_t - \frac{1}{2} kx^{-1} u \partial_u, t \partial_t + x \partial_x, xu \partial_u, u \partial_u \rangle$, $A = B = 0$;
- 2) $A_4 \sim A_{4,2}$, $A_4 = \langle \partial_t - \frac{4}{9} kx^{-1} u \partial_u, t \partial_t + x \partial_x, \sqrt{|x|} u \partial_u, \sqrt{|x|} \ln |x| u \partial_u \rangle$,
 $A = 0$, $B = \frac{1}{4} x^{-2}$;
- 3) $A_4 \sim A_{4,5}$, $A_4 = \langle \partial_t - \frac{k}{m+2} x^{-1} u \partial_u, t \partial_t + x \partial_x, |x|^{\frac{1}{2}+\beta} u \partial_u, |x|^{\frac{1}{2}-\beta} u \partial_u \rangle$,
 $A = 0$, $B = mx^{-2}$, $m \neq 0$, -2 , $m < \frac{1}{4}$, $\beta = \sqrt{\frac{1}{4} - m}$;
- 4) $A_4 \sim A_{4,2}$, $A_4 = \langle \partial_t + \frac{1}{9} kx^{-1} (1 + 3 \ln |x|) u \partial_u, t \partial_t + x \partial_x, x^2 u \partial_u, x^{-1} u \partial_u \rangle$,
 $A = 0$, $B = -2x^{-2}$;
- 5) $A_4 \sim A_{4,6}$, $A_4 = \langle \partial_t - \frac{k}{m+2} x^{-1} u \partial_u, t \partial_t + x \partial_x, \sqrt{|x|} u \cos(\beta \ln |x|) \partial_u, \sqrt{|x|} u \sin(\beta \ln |x|) \partial_u \rangle$,
 $A = 0$, $B = mx^{-2}$, $m > \frac{1}{4}$, $\beta = \sqrt{m - \frac{1}{4}}$;
- 6) $A_4 \sim A_{4,3}$, $A_4 = \langle \partial_t - kx^{-1} u \partial_u, t \partial_t + x \partial_x, u \partial_u, u \ln |x| \partial_u \rangle$,
 $A = x^{-1}$, $B = 0$;
- 7) $A_4 \sim A_{3,4} \oplus A_1$, $A_4 = \langle \partial_t + kx^{-1} (1 + \ln |x|) u \partial_u, t \partial_t + x \partial_x, u \partial_u, x^{-1} u \partial_u \rangle$,
 $A = 2x^{-1}$, $B = 0$;
- 8) $A_4 \sim A_{3,7} \oplus A_1$, $A_4 = \langle \partial_t + \frac{k}{n-2} x^{-1} u \partial_u, t \partial_t + x \partial_x, u \partial_u, |x|^{1-n} u \partial_u \rangle$,
 $A = nx^{-1}$, $B = 0$, $n \neq 0, 1, 2$;
- 9) $A_4 \sim A_{4,4}$, $A_4 = \langle \partial_t - \frac{1}{2} kx^{-1} \ln^2 |x| u \partial_u, t \partial_t + x \partial_x, x^{-1} u \partial_u, x^{-1} \ln |x| u \partial_u \rangle$,
 $A = 3x^{-1}$, $B = x^{-2}$;
- 10) $A_4 \sim A_{4,2}$, $A_4 = \langle \partial_t - \frac{4k}{(m-3)^2} x^{-1} u \partial_u, t \partial_t + x \partial_x, |x|^{\frac{1}{2}(1-n)} u \partial_u, |x|^{\frac{1}{2}(1-n)} \ln |x| u \partial_u \rangle$,
 $A = nx^{-1}$, $B = \frac{1}{4}(n-1)^2 x^{-2}$, $n \neq 0, 3$;

- 11) $A_4 \sim A_{4.5}$, $A_4 = \langle t\partial_t + x\partial_x, \partial_t - \frac{k}{2-n+m}x^{-1}u\partial_u, |x|^{\frac{1}{2}(1-n+\beta)}u\partial_u, |x|^{\frac{1}{2}(1-n-\beta)}u\partial_u \rangle$, $A = nx^{-1}$, $B = mx^{-2}$, $n \neq 0, 2$, $m \neq n - 2$, $m < \frac{1}{4}(n - 1)^2$, $\beta = \sqrt{(n - 1)^2 - 4m}$;
- 12) $A_4 \sim A_{4.2}$, $A_4 = \langle t\partial_t + x\partial_x, \partial_t + \frac{k}{3-n}x^{-1} \ln |x|u\partial_u, x^{-1}u\partial_u, |x|^{2-n}u\partial_u \rangle$, $A = nx^{-1}$, $B = (n - 2)x^{-2}$, $n \neq 0, 2, 3$;
- 13) $A_4 \sim A_{4.6}$, $A_4 = \langle t\partial_t + x\partial_x, \partial_t - \frac{k}{2-n+m}x^{-1}u\partial_u, |x|^{\frac{1}{2}(1-n)}u \cos(\beta \ln |x|)\partial_u, |x|^{\frac{1}{2}(1-n)}u \sin(\beta \ln |x|)\partial_u \rangle$, $A = nx^{-1}$, $B = mx^{-2}$, $n \neq 0$, $m \neq 0$, $m > \frac{1}{4}(n - 1)^2$, $\beta = \sqrt{m - \frac{1}{4}(n - 1)^2}$;

III. $D = kt$, $k > 0$,

- 1) $A_4 \sim A_{4.1}$, $A_4 = \langle \partial_x, \partial_t - \frac{1}{2}kx^2u\partial_u, xu\partial_u, u\partial_u \rangle$, $A = B = 0$;
- 2) $A_4 \sim A_{4.3}$, $A_4 = \langle \partial_x, \partial_t - kxu\partial_u, e^{-x}u\partial_u, u\partial_u \rangle$, $A = 1$, $B = 0$;
- 3) $A_4 \sim A_{3.8} \oplus A_1$, $A_4 = \langle \partial_x, \partial_t - k\beta^{-2}u\partial_u, u \cos(\beta x)\partial_u, u \sin(\beta x)\partial_u \rangle$, $A = 0$, $B = \beta^2$, $\beta \neq 0$;
- 4) $A_4 \sim A_{3.6} \oplus A_1$, $A_4 = \langle \partial_x, \partial_t + k\beta^{-2}u\partial_u, u \cosh(\beta x)\partial_u, u \sinh(\beta x)\partial_u \rangle$, $A = 0$, $B = -\beta^2$, $\beta \neq 0$;
- 5) $A_4 \sim A_{3.4} \oplus A_1$, $A_4 = \langle \partial_x, \partial_t - 4ku\partial_u, \exp(-\frac{1}{2}x)u\partial_u, x \exp(-\frac{1}{2}x)u\partial_u \rangle$, $A = 1$, $B = \frac{1}{4}$;
- 6) $A_4 \sim A_{3.7} \oplus A_1$, $A_4 = \langle \partial_x, \partial_t - km^{-1}u\partial_u, \exp(-\frac{1}{2}(1-\beta)x)u\partial_u, \exp(-\frac{1}{2}(1+\beta)x)u\partial_u \rangle$, $A = 1$, $B = m$, $m < \frac{1}{4}$, $m \neq 0$, $\beta = \sqrt{1 - 4m}$;
- 7) $A_4 \sim A_{3.9} \oplus A_1$, $A_4 = \langle \partial_x, \partial_t - km^{-1}u\partial_u, \exp(-\frac{1}{2}x) \cos(\beta x)u\partial_u, \exp(-\frac{1}{2}x) \sin(\beta x)u\partial_u \rangle$, $A = 1$, $B = m$, $m > \frac{1}{4}$, $\beta = \sqrt{m - \frac{1}{4}}$;

IV. $D = kt^{-2}$, $k \neq 0$,

$A_4 \sim A_{4.8}$ ($q = -1$), $A_4 = \langle \partial_x, t\partial_t + x\partial_x, xu\partial_u, u\partial_u \rangle$, $A = B = 0$;

V. $D = m(x - kt)^{-2}$, $k > 0$, $m \neq 0$,

$A_4 \sim A_{4.8}$ ($q = -1$), $A_4 = \langle \partial_t + k\partial_x, t\partial_t + x\partial_x, xu\partial_u, u\partial_u \rangle$, $A = B = 0$.

Proof. According to Lemma 5 to get the list of inequivalent equations of the form (6.1) we need to analyze the cases when either $D = 0$ or D does not vanish identically and is obtained through Equation (6.8).

If $D = 0$, then the function h satisfies Equation (6.3) and the functions A, B are given by one of the formulas (6.5). It follows from (6.2) that five is the highest possible dimension of an invariance algebra admitted by (6.1). Equation admitting this algebra is equivalent to the following one:

$$u_{tt} = u_{xx} - u^{-1}u_x^2. \tag{6.10}$$

This proves the first part of the assertion of theorem.

The remaining expressions for the functions A, B from (6.5) yield the 14 sub-cases of the case I of the second assertion of the theorem.

The fourth expression for the function D given in (6.8) within the equivalence relation (6.7) boils down to either $D = ktx^{-3}$, $k > 0$ or $D = kt$, $k > 0$. The function h satisfies one of the equations

$$h'' + Ah' + Bh = -\lambda_1 H,$$

where either $H = kx^{-3}$ or $H = k$. Analysis of the corresponding expressions for the functions A, B (6.5) yields expressions listed in the cases II and III.

Next the function D given by the third formula from (6.8) simplifies to $D = kt^{-2}$ ($k \neq 0$) whence we get the results listed in the case IV of the second assertion of the theorem. Similarly, the second expression for the function D gives rise to the formulas of the case V.

The first expression for D from (6.8) gives no new invariant equations.

What is left is to prove that the so obtained invariant equations are inequivalent. We skip the proof of this fact.

Theorem is proved.

7.2. Nonlinear Equations (2.7) Invariant Under Three-dimensional Lie Algebras

Class of PDEs (2.7) does not contain an equation whose invariance algebra is isomorphic to a Lie algebra with a non-trivial Levi ideal (see, Theorem 12). That is why, to complete the second step of our classification algorithm it suffices to consider three-dimensional solvable real Lie algebras only. We begin by considering two decomposable three-dimensional solvable Lie algebras.

Note that while classifying invariant equation (2.7) we skip equations belonging to the class (6.1) which has been analyzed in the previous subsection.

7.2.1. Invariance Under Decomposable Lie Algebras

As $A_{3,1} = 3A_1 = A_{2,1} \oplus A_1$, $A_{3,2} = A_{2,2} \oplus A_1$, to construct all realizations of $A_{3,1}$ it suffices to compute all possible extensions of the (already known) realizations of the algebras $A_{2,1} = \langle e_1, e_2 \rangle$ and $A_{2,2} = \langle e_1, e_2 \rangle$. To this end we need to supplement the latter by a basis operator e_3 of the form (2.4) in order to satisfy the commutation relations

$$[e_1, e_3] = [e_2, e_3] = 0. \quad (6.11)$$

What is more, to simplify the form of e_3 we can use transformations from \mathcal{E} which do not alter the remaining basis operators of the corresponding two-dimensional Lie algebras.

We do not present full calculation details. Instead, we give several examples illustrating the main steps which we need to perform in order to extend $A_{2,1}$ to a realization of $A_{3,1}$.

Consider the realization $A_{2,1}^1$. Upon checking commutation relations (6.11), where e_3 is of form (2.4), we get

$$\lambda_1 = \lambda_2 = r(t, x) = 0, \quad h = k = \text{const.}$$

Consequently e_3 is the linear combination of e_1, e_2 , namely, $e_3 = \lambda e_1 + ke_2$ which is impossible by the assumption that the algebra under study is three-dimensional. Hence we conclude that the above realization of $A_{2,1}^1$ cannot be extended to a realization of the algebra $A_{3,1}$.

Turn now to the realization $A_{2,1}^2$. Checking commutation relations (6.11), where e_3 is of form (2.4), yields the following realization of $A_{3,1}$:

$$\langle t\partial_t + x\partial_x, \sigma(\xi)\partial_u, \gamma(\xi)\partial_u \rangle, \quad \xi = tx^{-1},$$

where $\gamma'\sigma - \gamma\sigma' \neq 0$. However, the corresponding invariant equation (2.7) is linear.

Finally, consider the realization $A_{2,1}^3$. Inserting its basis operators and the operator e_3 of the form (2.4) into (6.11) and solving the obtained equations gives the following realization of $A_{3,1}$:

$$\langle \partial_t, \partial_x, u\partial_u \rangle.$$

Inserting the obtained coefficients for e_3 into the classifying Equation (2.5) we get invariant equation

$$u_{tt} = u_{xx} + uG(\omega), \quad \omega = u^{-1}u_x,$$

where in order to ensure nonlinearity we need to have $G_{\omega\omega} \neq 0$.

A similar analysis of the realizations $A_{2,1}^i$ ($i = 4, 5, \dots, 12, 14$) yields three new invariant equations. For two of so obtained $A_{3,1}$ -invariant equations the corresponding three-dimensional algebras are maximal. The other two may admit four-dimensional invariance algebras provided arbitrary elements are properly specified.

Handling in a similar way extensions of $A_{2,2}$ up to realizations of $A_{3,2}$ results in 10 inequivalent nonlinear equations the maximal invariance algebras of which are realizations of the three-dimensional algebra $A_{3,2}$ and four inequivalent equation (2.7) admitting four-dimensional symmetry algebras.

We perform analysis of the equations admitting four-dimensional algebras in the next sub-section. Here we present the complete list of nonlinear equation (2.7) the maximal symmetry algebras of which are realizations of three-dimensional Lie algebras $A_{3,1}$ and $A_{3,2}$.

$A_{3,1}$ -invariant equations

$$\begin{aligned} A_{3,1}^1 &= \langle \partial_t, \partial_x, u\partial_u \rangle : \\ &F = uG(\omega), \quad \omega = u^{-1}u_x; \\ A_{3,1}^2 &= \langle \partial_x, \varphi(t)\partial_u, \psi(t)\partial_u \rangle : \\ &\sigma = \psi'\varphi - \psi\varphi' \neq 0, \quad \sigma' = 0 : \\ &F = \varphi^{-1}\varphi''u + G(t, u_x). \end{aligned}$$

$A_{3,2}$ -invariant equations

$$\begin{aligned} A_{3,2}^1 &= \langle \partial_t, \partial_x, e^x u\partial_u \rangle : \\ &F = -u^{-1}u_x^2 - u \ln |u| + uG(\omega), \\ &\omega = u^{-1}u_x - \ln |u| : \\ A_{3,2}^2 &= \langle -t\partial_t - x\partial_x, \partial_t + k\partial_x, u\partial_u \rangle \ (k \geq 0) : \\ &F = u\eta^{-2}G(\omega), \quad \eta = x - kt, \\ &\omega = \eta u^{-1}u_x; \end{aligned}$$

$$A_{3.2}^3 = \langle -t\partial_t - x\partial_x + mu\partial_u, \partial_t + k\partial_x, |\eta|^{-m}\partial_u \rangle$$

$$(\eta = x - kt, k = m = 0 \text{ or } k > 0, m \in \mathbb{R}) :$$

$$F = m(k^2 - 1)(m + 1)\eta^{-2}u + |\eta|^{-2-m}G(\omega),$$

$$\omega = |\eta|^m(mu + \eta u_x);$$

$$A_{3.2}^4 = \langle \partial_x, e^x u \partial_u, \partial_t + mu \partial_u \rangle (m > 0) :$$

$$F = -u^{-1}u_x^2 - u_x + uG(\omega),$$

$$\omega = u^{-1}u_x - \ln|u| + mt;$$

$$A_{3.2}^5 = \langle -t\partial_t - x\partial_x, \partial_x, u\partial_u \rangle :$$

$$F = ut^{-2}G(\omega), \quad \omega = tu^{-1}u_x;$$

$$A_{3.2}^6 = \langle -t\partial_t - x\partial_x, \partial_t + kx^{-1}u\partial_u, u\partial_u \rangle (k > 0) :$$

$$F = 2ktx^{-2}u_x - 2ktx^{-3}u + k^2t^2x^{-4}u + x^{-2}uG(\omega),$$

$$\omega = xu^{-1}u_x + ktx^{-1};$$

$$A_{3.2}^7 = \langle -t\partial_t - x\partial_x, \partial_t + kx^{-1}u\partial_u, \exp(ktx^{-1})\partial_u \rangle (k > 0) :$$

$$F = 2ktx^{-2}u_x + (k^2t^2x^{-4} - 2ktx^{-3} + k^2x^{-2})u +$$

$$+ x^{-2}\exp(ktx^{-1})G(\omega), \quad \omega = \exp(-ktx^{-1})(xu_x + ktx^{-1}u);$$

$$A_{3.2}^8 = \langle \frac{1}{2k}(\partial_t + k\partial_x), e^{x+kt}\partial_u, e^\eta\partial_u \rangle (k > 0, \eta = x - kt) :$$

$$F = (k^2 - 1)u + G(\eta, \omega), \quad \omega = u_x - u;$$

$$A_{3.2}^9 = \langle \partial_t + f(x)u\partial_u, e^{(1+f(x))t}\partial_u, f(x)e^{f(x)t}\partial_u \rangle :$$

$$F = -(tf'' - t^2(f')^2 - (1 + f)^2)u - 2tf'u_x + e^{tf}G(x, \omega),$$

$$\omega = e^{-tf}(u_x - f'(t + f^{-1})u), \quad f'' + 2f^2 + f = 0, \quad f \neq 0;$$

$$A_{3.2}^{10} = \langle k(t\partial_t + x\partial_x), |t|^{k-1}|\xi|^{\frac{k-1}{2k}}\partial_u, |\xi|^{\frac{k-1}{2k}}\partial_u \rangle (k \neq 0; 1) :$$

$$F = \left[\frac{1-k}{k}\xi^2 + \frac{1-k^2}{4k^2}(1-\xi^2) \right] t^{-2}u + t^{-2}G(\xi, \omega),$$

$$\omega = |\xi|^{\frac{k-1}{2k}} \left[xu_x + \frac{k-1}{2k}u \right], \quad \xi = tx^{-1}.$$

7.2.2. Invariance Under Non-decomposable Three-dimensional Solvable Lie Algebras

There exist seven non-decomposable three-dimensional solvable Lie algebras over the field of real numbers. All these algebras contain a subalgebra which is the two-dimensional Abelian ideal. Consequently, we can use the results of classification of $A_{2,1}$ -invariant equations in order to describe equations admitting non-decomposable three-dimensional solvable real Lie algebras. We remind that equations of the form (6.1) have already been analyzed and therefore are not considered in the sequel.

As an example we compute extension of the realization $A_{2,1}^{10}$ to all possible realizations of non-decomposable three-dimensional solvable real Lie algebras. The remaining realizations are handled in a similar way.

It is straightforward to verify that transformations

$$\bar{t} = \gamma t + \gamma_1, \quad \bar{x} = \epsilon \gamma x + \gamma_2, \quad v = \rho(x)u + \theta(x), \tag{6.12}$$

where $\gamma, \gamma_1, \gamma_2 \in \mathbb{R}, \gamma \neq 0, \epsilon = \pm 1, \rho \neq 0$, are equivalence transformations for the realization $A_{2,1}^{10} = \langle \partial_t, f(x)u\partial_u \rangle (f \neq 0)$. That is why, we can utilize the above transformation to simplify the form of operator e_3 . As a result, we get three inequivalent expressions for e_3

- $e_3 = t\partial_t + x\partial_x + r(t, x)\partial_u (r_t \neq 0 \text{ or } r = 0)$;
- $e_3 = \partial_x + r(t, x)\partial_u (r_t \neq 0 \text{ or } r = 0)$;
- $e_3 = r(t, x)\partial_u (r_t \neq 0 \text{ or } r = 1)$.

Let $e_1 = \partial_t, e_2 = f(x)u\partial_u$ and $e_3 = t\partial_t + x\partial_x + r(t, x)\partial_u$, then

$$[e_1, e_3] = \partial_t + r_t\partial_u, \quad [e_2, e_3] = -xf'u\partial_u - rf\partial_u.$$

Analyzing commutation relations for the algebras $A_{3,i} (i = 3, 4, \dots, 9)$ we obtain that the necessary conditions for $A_{2,1}^{10}$ to admit extension to a realization of $A_{3,5}$ are $r = 0, xf' = -f$, of $A_{3,6}$ as $r = 0, xf' = f$, and of $A_{3,7}$ as $r = 0, xf' = -qf (0 < |q| < 1)$. So $A_{2,1}^{10}$ gives rise to the following realizations:

- $A_{3,5} : e_1 = \partial_t, e_2 = x^{-1}u\partial_u, e_3 = t\partial_t + x\partial_x$;
- $A_{3,6} : e_1 = \partial_t, e_2 = xu\partial_u, e_3 = t\partial_t + x\partial_x$;
- $A_{3,7} : e_1 = \partial_t, e_2 = |x|^{-q}u\partial_u, e_3 = t\partial_t + x\partial_x (0 < |q| < 1)$.

If $e_3 = \partial_x + r(t, x)\partial_u$, then

$$[e_1, e_3] = r_t\partial_u, \quad [e_2, e_3] = -f'u\partial_u - rf\partial_u.$$

Analyzing the commutation relations for $A_{3,i} (i = 3, 4, \dots, 9)$ we come to conclusion that the realization $A_{2,1}^{10}$ cannot be extended to a realization of the above three-dimensional Lie algebras.

The same conclusion holds true when $e_3 = r(t, x)\partial_u (r_t \neq 0 \text{ or } r = 1)$.

Let $e_1 = f(x)u\partial_u, e_2 = \partial_t$. If $e_3 = t\partial_t + x\partial_x + r(t, x)\partial_u (r_t \neq 0 \text{ or } r = 0)$, then it follows from commutation relations

$$[e_1, e_3] = -(rf + xf'u)\partial_u, \quad [e_2, e_3] = \partial_t + r_t\partial_u$$

that the only possible extension of the realization $A_{2,1}^{10}$ is the realization of $A_{3,5}$:

$$\langle x^{-1}u\partial_u, \partial_t, t\partial_t + x\partial_x \rangle.$$

This realization coincides within notation with the already obtained one.

Next if $e_3 = \partial_x + r(t, x)\partial_u$ ($r_t \neq 0$ or $r = 0$), then

$$[e_1, e_3] = -(f'u + rf)\partial_u, \quad [e_2, e_3] = r_t\partial_u.$$

Analyzing the commutation relations for $A_{3,i}$ ($i = 3, 4, \dots, 9$) we come to conclusion that the realization $A_{2,1}^{10}$ cannot be extended to a realization of the above three-dimensional Lie algebras.

The same conclusion holds true for the case $e_3 = r(t, x)\partial_u$ ($r_t \neq 0$ or $r = 1$).

Summing up the above considerations we see that the realization $A_{2,1}^{10}$ can be extended to the following realizations of non-decomposable three-dimensional solvable real Lie algebras:

$$L^1 \sim A_{3,5}, \quad L^1 = \langle \partial_t, x^{-1}u\partial_u, t\partial_t + x\partial_x \rangle;$$

$$L^2 \sim A_{3,6}, \quad L^2 = \langle \partial_t, xu\partial_u, t\partial_t + x\partial_x \rangle;$$

$$L^3 \sim A_{3,7}, \quad L^3 = \langle \partial_t, |x|^{-q}u\partial_u, t\partial_t + x\partial_x \rangle \quad (0 < |q| < 1).$$

Solving the corresponding classifying equations yields the following invariant equations:

$$L^1 : u_{tt} = u_{xx} - u^{-1}u_x^2 - 2x^{-2}u \ln |u| + x^{-2}uG(\omega), \quad \omega = xu^{-1}u_x + \ln |u|;$$

$$L^2 : u_{tt} = u_{xx} - u^{-1}u_x^2 + x^{-2}uG(\omega),$$

$$\omega = xu^{-1}u_x - \ln |u|;$$

$$L^3 : u_{tt} = u_{xx} - u^{-1}u_x^2 - q(q+1)x^{-2}u \ln |u| + ux^{-2}G(\omega),$$

$$\omega = xu^{-1}u_x + q \ln |u| \quad (0 < |q| < 1).$$

Note that the algebras L^1, L^2, L^3 are maximal (in Lie's sense) invariance algebras of the corresponding equations.

While classifying nonlinear equations invariant under non-decomposable three-dimensional solvable Lie algebras we have discovered equations whose maximal invariance algebras are four-dimensional. For example, after extending the realization $A_{2,1}^9$ up to a realization of the algebra $A_{3,3}$ we got the following realization of the latter:

$$\langle \partial_u, \partial_t, \partial_x + t\partial_u \rangle.$$

The corresponding invariant equation (2.7) is $u_{tt} = u_{xx} + G(u_x)$. However, the maximal invariance algebra of this equation is the four-dimensional Lie algebra $\langle \partial_t, t\partial_u, \partial_u, \partial_x \rangle$, which is a realization of $A_{3,3} \oplus A_1$. Note also that we have obtained the above invariant equation when classifying $A_{3,1}$ -invariant equations.

By the above reason, we give below only those nonlinear invariant equations the maximal symmetry algebras of which are three-dimensional non-decomposable solvable real Lie algebras.

A_{3,3}-invariant equations

$$\begin{aligned}
 A_{3,3}^1 &= \langle u\partial_u, \partial_t + k\partial_x, m\partial_t + k^{-1}xu\partial_u \rangle \ (k > 0, m \neq 0) : \\
 &F = -u^{-1}u_x^2 + uG(\omega), \ \omega = x - kt + mk^2u^{-1}u_x; \\
 A_{3,3}^2 &= \langle u\partial_u, \partial_x, m\partial_t + xu\partial_u \rangle \ (m > 0) : \\
 &F = -u^{-1}u_x^2 + uG(\omega), \ \omega = t - mu^{-1}u_x; \\
 A_{3,3}^3 &= \langle |t|^{\frac{1}{2}}\partial_u, -|t|^{\frac{1}{2}}\ln|t|\partial_u, t\partial_t + x\partial_x + \frac{1}{2}u\partial_u \rangle : \\
 &F = -\frac{1}{4}t^{-2}u + u_x^3G(\xi, \omega), \ \xi = tx^{-1}, \ \omega = xu_x^2; \\
 A_{3,3}^4 &= \langle \partial_u, -t\partial_u, \partial_t + k\partial_x \rangle \ (k \geq 0); \\
 &F = G(\eta, u_x), \ \eta = x - kt.
 \end{aligned}$$

A_{3,4}-invariant equations

$$\begin{aligned}
 A_{3,4}^1 &= \langle |\eta|^{m-1}\partial_u, \partial_t + k\partial_x, t\partial_t + x\partial_x + (mu + t|\eta|^{m-1})\partial_u \rangle \\
 &(\eta = x - kt, \ k > 0, m \neq 1) : \\
 &F = (k^2 - 1)(m - 1)(m - 2)\eta^{-2}u - 2k(m - 1)\eta^{m-2}\ln|\eta| \\
 &+ |\eta|^{m-2}G(\omega), \ \omega = [\eta u_x - (m - 1)u]|\eta|^{-m}; \\
 A_{3,4}^2 &= \langle \partial_u, -t\partial_u, \partial_t + k\partial_x + u\partial_u \rangle \ (k \geq 0) : \\
 &F = e^tG(\eta, \omega), \ \eta = x - kt, \ \omega = e^{-t}u_x; \\
 A_{3,4}^3 &= \langle |t|^{\frac{1}{2}}\partial_u, -|t|^{\frac{1}{2}}\ln|t|\partial_u, t\partial_t + x\partial_x + \frac{3}{2}u\partial_u \rangle : \\
 &F = -\frac{1}{4}t^{-2}u + u_x^{-1}G(\xi, \omega), \ \xi = tx^{-1}, \ \omega = x^{-1}u_x^2; \\
 A_{3,4}^4 &= \langle kx^{-1}u\partial_u, \partial_t - kx^{-1}\ln|x|u\partial_u, t\partial_t + x\partial_x \rangle \ (k > 0) : \\
 &F = -3ktx^{-3}u - 2x^{-2}u\ln|u| - u^{-1}u_x^2 + x^{-2}uG(\omega), \\
 &\omega = xu^{-1}u_x + \ln|u| + ktx^{-1}; \\
 A_{3,4}^5 &= \langle \exp(ktx^{-1})\partial_u, \partial_t + kx^{-1}u\partial_u, t\partial_t + x\partial_x + (u + t\exp(ktx^{-1}))\partial_u \rangle \ (k > 0) : \\
 &F = k^2x^{-4}u(t^2 + x^2) + 2x^{-1}(ktx^{-1} + 1)u_x \\
 &+ 2k\exp(ktx^{-1})x^{-1}\ln|x| + x^{-1}\exp(ktx^{-1})G(\omega), \\
 &\omega = \exp(-ktx^{-1})(u_x + ktx^{-2}u).
 \end{aligned}$$

A_{3.5}-invariant equations

$$A_{3,5}^1 = \langle |\eta|^{m-1} \partial_u, \partial_t + k \partial_x, t \partial_t + x \partial_x + mu \partial_u \rangle \quad (k > 0, m \neq 1)$$

$$F = (k^2 - 1)(m - 1)(m - 2)\eta^{-2}u + |\eta|^{m-2}G(\omega),$$

$$\omega = |\eta|^{-m}[\eta u_x - (m - 1)u], \eta = x - kt;$$

$$A_{3,5}^2 = \langle \partial_t, \partial_x, t \partial_t + x \partial_x \rangle :$$

$$F = u_x^2 G(u);$$

$$A_{3,5}^3 = \langle \partial_t, \partial_x, t \partial_t + x \partial_x + mu \partial_u \rangle \quad (m \neq 0) :$$

$$F = |u|^{1-\frac{2}{m}} G(\omega), \quad \omega = |u_x|^m |u|^{1-m};$$

$$A_{3,5}^4 = \langle \partial_t, \partial_x, t \partial_t + x \partial_x + \partial_u \rangle :$$

$$F = e^{-2u} G(\omega), \quad \omega = e^u u_x;$$

$$A_{3,5}^5 = \langle \partial_t, x^{-1} u \partial_u, t \partial_t + x \partial_x \rangle :$$

$$F = -u^{-1} u_x^2 - 2x^{-2} u \ln |u| + x^{-2} u G(\omega),$$

$$\omega = x u^{-1} u_x + \ln |u|;$$

$$A_{3,5}^6 = \langle \partial_t + kx^{-1} u \partial_u, \exp(ktx^{-1}) \partial_u, t \partial_t + x \partial_x + u \partial_u \rangle \quad (k > 0) :$$

$$F = kx^{-4} u [kt^2 - 2tx + kx^2] + 2ktx^{-2} u_x + x^{-1} \exp(ktx^{-1}) G(\omega),$$

$$\omega = \exp(-ktx^{-1})(u_x + ktx^{-2}u);$$

$$A_{3,5}^7 = \langle \varphi(t) \partial_u, \psi(t) \partial_u, \partial_x + u \partial_u \rangle \quad (\varphi' \psi - \varphi \psi' \neq 0) :$$

$$F = \varphi^{-1} \varphi'' u + u_x G(t, \omega),$$

$$\omega = e^{-x} u_x, \quad \varphi'' \psi - \varphi \psi'' = 0.$$

A_{3.6}-invariant equations

$$A_{3,6}^1 = \langle \partial_t + k \partial_x, |\eta|^{m+1} \partial_u, t \partial_t + x \partial_x + mu \partial_u \rangle \quad (k > 0, m \neq -1) :$$

$$F = m(k^2 - 1)(m + 1)\eta^{-2}u + |\eta|^{m-2}G(\omega),$$

$$\omega = |\eta|^{1-m}[u_x - \eta^{-1}(m + 1)u], \eta = x - kt;$$

$$A_{3,6}^2 = \langle \partial_t + mx^{-1}u\partial_u, xu\partial_u, t\partial_t + x\partial_x \rangle \ (m \geq 0) :$$

$$F = -u^{-1}u_x^2 - 2mtx^{-3}u + x^{-2}uG(\omega),$$

$$\omega = xu^{-1}u_x - \ln |u| + 2mtx^{-1};$$

$$A_{3,6}^3 = \langle \partial_t + kx^{-1}u\partial_u, \exp(ktx^{-1})\partial_u, t\partial_t + x\partial_x - u\partial_u \rangle \ (k > 0) :$$

$$F = x^{-4}[k^2x^2 - 2ktx + k^2t^2]u + 2ktx^{-2}u_x + x^{-3} \exp(ktx^{-1})G(\omega),$$

$$\omega = \exp(-ktx^{-1})(x^2u_x + ktu);$$

$$A_{3,6}^4 = \langle e^{-t}\partial_u, e^t\partial_u, \partial_t + k\partial_x \rangle \ (k \geq 0) :$$

$$F = u + G(\eta, u_x), \ \eta = x - kt;$$

$$A_{3,6}^5 = \langle |t|^{-\frac{1}{2}}\partial_u, |t|^{\frac{3}{2}}\partial_u, t\partial_t + x\partial_x + \frac{1}{2}u\partial_u \rangle :$$

$$F = \frac{3}{4}t^{-2}u + |t|^{-\frac{3}{2}}G(\xi, \omega), \ \xi = tx^{-1}, \ \omega = x^{-1}u_x^2.$$

A_{3,7}-invariant equations

$$A_{3,7}^1 = \langle \partial_t + k\partial_x, |\eta|^{m-q}\partial_u, t\partial_t + x\partial_x + mu\partial_u \rangle$$

$$(k > 0, m \neq q, 0 < |q| < 1) :$$

$$F = (k^2 - 1)(m - q)(m - q - 1)\eta^{-2}u + |\eta|^{m-2}G(\omega),$$

$$\omega = |\eta|^{1-m}[u_x - (m - q)\eta^{-1}u], \ \eta = x - kt;$$

$$A_{3,7}^2 = \langle \partial_t + kx^{-1}u\partial_u, \exp(ktx^{-1})\partial_u, t\partial_t + x\partial_x + qu\partial_u \rangle$$

$$(k > 0, 0 < |q| < 1) :$$

$$F = [k^2x^{-2} + k^2x^{-4}t^2 - 2ktx^{-3}]u + 2ktx^{-2}u_x + |x|^{q-2} \exp(ktx^{-1})G(\omega),$$

$$\omega = |x|^{1-q} \exp(-ktx^{-1})(u_x + ktx^{-2}u);$$

$$A_{3,7}^3 = \langle |t|^{\frac{1}{2}q}\partial_u, |t|^{1-\frac{1}{2}q}\partial_u, t\partial_t + x\partial_x + (1 + \frac{1}{2}q)u\partial_u \rangle \ (q \neq 0, \pm 1) :$$

$$F = \frac{1}{4}q(q - 2)t^{-2}u + |t|^{\frac{1}{2}(q-2)}G(\xi, \omega),$$

$$\xi = tx^{-1}, \ \omega = |t|^{-\frac{1}{2}q}u_x;$$

$$A_{3,7}^4 = \langle \exp\left(\frac{1}{2}(q-1)t\right) \partial_u, \exp\left(\frac{1}{2}(1-q)t\right) \partial_u, \partial_t + k\partial_x + \frac{1}{2}(1+q)u\partial_u \rangle$$

$$(q \neq 0, \pm 1; k \geq 0) :$$

$$F = \frac{1}{4}(q-1)^2 u + \exp\left(\frac{1}{2}(1+q)t\right) G(\eta, \omega),$$

$$\eta = x - kt, \quad \omega = \exp\left(-\frac{1}{2}(1+q)t\right) u_x;$$

$$A_{3,7}^5 = \langle \partial_t + kx^{-1}u\partial_u, |x|^{-q}u\partial_u, t\partial_t + x\partial_x \rangle (k \geq 0, q \neq 0, \pm 1) :$$

$$F = -u^{-1}u_x^2 - q(q+1)x^{-2}u \ln |u| + k(q-1)(q+2)tx^{-3}u$$

$$+ ux^{-2}G(\omega), \quad \omega = xu^{-1}u_x + q \ln |u| + k(1-q)tx^{-1}.$$

A_{3,8}-invariant equations

$$A_{3,8}^1 = \langle \cos t \partial_u, -\sin t \partial_u, \partial_t + k\partial_x \rangle (k \geq 0) :$$

$$F = -u + G(\eta, u_x), \quad \eta = x - kt;$$

$$A_{3,8}^2 = \langle |t|^{\frac{1}{2}} \cos(\ln |t|) \partial_u, -|t|^{\frac{1}{2}} \sin(\ln |t|) \partial_u, t\partial_t + x\partial_x + \frac{1}{2}u\partial_u \rangle :$$

$$F = -\frac{5}{4}t^{-2}u + |t|^{-\frac{3}{2}}G(\xi, \omega),$$

$$\xi = tx^{-1}, \quad \omega = |t|^{\frac{1}{2}}u_x.$$

A_{3,9}-invariant equations

$$A_{3,9}^1 = \langle \sin t \partial_u, \cos t \partial_u, \partial_t + k\partial_x + qu\partial_u \rangle (k \geq 0, q > 0) :$$

$$F = -u + e^{qt}G(\eta, \omega), \quad \eta = x - kt, \quad \omega = e^{-qt}u_x;$$

$$A_{3,9}^2 = \langle |t|^{\frac{1}{2}} \sin(\ln |t|) \partial_u, |t|^{\frac{1}{2}} \cos(\ln |t|) \partial_u, t\partial_t + x\partial_x + \left(\frac{1}{2} + q\right)u\partial_u \rangle$$

$$(q \neq 0) : F = -\frac{5}{4}t^{-2}u + |t|^{q-\frac{3}{2}}G(\xi, \omega),$$

$$\xi = tx^{-1}, \quad \omega = |t|^{\frac{1}{2}-q}u_x.$$

7.3. Complete Group Classification of Equation (2.7)

The aim of this subsection is to finalize group classification of (2.7). The majority of invariant equations obtained in the previous sub-section contain arbitrary functions of one variable. That is why we can utilize the standard Lie–Ovsyannikov approach in order to complete their group classification.

7.3.1. Equations Containing Arbitrary Functions of One Variable

Note that we do not consider the equations belonging to the already studied class (6.1).

As our computations show, new results could be obtained for the equations

$$u_{tt} = u_{xx} + uG(\omega), \quad \omega = u^{-1}u_x, \tag{6.13}$$

$$u_{tt} = u_{xx} + G(u_x) \tag{6.14}$$

only. Below we give (without proof) the assertions describing their group properties.

ASSERTION 7. Equation (6.13) admits wider symmetry group iff it is equivalent to the following PDE:

$$u_{tt} = u_{xx} + mu^{-1}u_x^2 \quad (m \neq 0, -1). \tag{6.15}$$

The maximal invariance algebra of (6.15) is the four-dimensional Lie algebra

$$A_4 \sim A_{3.5} \oplus A_1, \quad A_4 = \langle \partial_t, \partial_x, t\partial_t + x\partial_x, u\partial_u \rangle.$$

ASSERTION 8. Equation (6.14) admits wider symmetry group iff it is equivalent to one of the following PDEs:

$$u_{tt} = u_{xx} + e^{u_x}; \tag{6.16}$$

$$u_{tt} = u_{xx} + m \ln |u_x|, \quad m > 0; \tag{6.17}$$

$$u_{tt} = u_{xx} + |u_x|^k, \quad k \neq 0, 1. \tag{6.18}$$

The maximal invariance algebras of the above equations are five-dimensional solvable Lie algebras given below.

$$A_5^2 = \langle \partial_t, \partial_x, \partial_u, t\partial_u, t\partial_t + x\partial_x + (u - x)\partial_u \rangle;$$

$$A_5^3 = \langle \partial_t, \partial_x, \partial_u, t\partial_u, t\partial_t + x\partial_x + (2u + \frac{1}{2}mt^2)\partial_u \rangle;$$

$$A_5^4 = \langle \partial_t, \partial_x, \partial_u, t\partial_u, t\partial_t + x\partial_x + \frac{k-2}{k-1}u\partial_u \rangle.$$

Analyzing the remaining equations containing arbitrary functions of one variable we come to conclusion that for any them to admit wider invariance group one of the following conditions should hold true:

- equation in question is equivalent to PDE of the form (6.1), or
- equation in question is equivalent to PDE of the form (6.15).

Skipping the proof, we present two typical examples. We begin with the equation

$$u_{tt} = u_{xx} + u + G(u_x). \tag{6.19}$$

This equation is invariant under the four-dimensional algebra $\langle \partial_t, \partial_x, e^t \partial_u, e^{-t} \partial_u \rangle$ isomorphic to $A_{3,6} \oplus A_1$. Inserting $F = u + G(u_x)$ into classifying Equation (2.5) yields the system of two equations for G

$$h'G' = -h'' - 2\lambda, \quad [(h - \lambda)u_x + r_x]G' - (h - 2\lambda)G = r_{tt} - r_{xx} - 2h'u_x - r.$$

As $G'' \neq 0$, the first equation implies that $\lambda = h' = 0$. The second equation takes the form

$$(hu_x + r_x)G' - hG = r_{tt} - r_{xx} - r.$$

Upon differentiating the above equation twice with respect to u_x we get $(hu_x + r_x)G'' = 0$. As $G'' \neq 0$, the relations $h = r_x = 0$ hold. Hence we conclude that the class of PDEs (6.19) does not contain equations admitting five-dimensional symmetry algebras.

The system of determining equations for the symmetry group of $A_{3,2}^2$ -invariant equation

$$u_{tt} = u_{xx} + u\eta^{-2}G(\omega), \quad \eta = x - kt, \quad \omega = \eta u^{-1}u_x, \quad k \geq 0, \quad (6.20)$$

is

$$\begin{aligned} (\eta^{-2}r\omega - \eta^{-1}r_x)G_\omega - \eta^{-2}rG &= r_{tt} - r_{xx}, \\ [(\lambda_2 - k\lambda_1)\eta^{-3}\omega + \eta^{-1}h']G_\omega - 2(\lambda_2 - k\lambda_1)\eta^{-3}G &= -2h'\eta^{-1}\omega - h''. \end{aligned}$$

Differentiating the first equation with respect to ω yields

$$(\eta^{-2}r\omega - \eta^{-1}r_x)G_{\omega\omega} = 0,$$

whence in view of inequality $G_{\omega\omega} \neq 0$ we get $r = 0$. Next differentiating the second equation twice by ω we get

$$[(\lambda_2 - k\lambda_1)\eta^{-3}\omega + \eta^{-1}h']G_{\omega\omega\omega} = 0,$$

whence it follows that $G_{\omega\omega\omega} = 0$. Indeed, if this relation does not hold, we have $\lambda_2 = k\lambda_1$, $h' = 0$ and operator (2.4) takes the form

$$\lambda(t\partial_t + x\partial_x) + \lambda_1(\partial_t + k\partial_x) + C_1u\partial_u, \quad \lambda, \lambda_1, C_1 \in \mathbb{R}, \quad k \geq 0.$$

As the above operator contains at most three arbitrary constants it cannot generate a four-parameter Lie transformation group.

By the above argument we can restrict our considerations to the following class of functions G :

$$G = A\omega^2 + B\omega + C, \quad A \neq 0, -1, \quad B, C \in \mathbb{R}. \quad (6.21)$$

Without any loss of generality we can suppose that $A \neq -1$ in (6.21) (since otherwise (6.20) belongs to the class of PDEs (6.1)). Inserting (6.21) into the second equation from (6.21) yields

$$2(A+1)\eta^2h' = B(\lambda_2 - k\lambda_1), \quad \eta^2Bh' + \eta^3h'' = 2C(\lambda_2 - k\lambda_1). \quad (6.22)$$

If $k > 0$, then this system splits into the three equations (note that $h = h(x)$)

$$h' = 0, \quad B(\lambda_2 - k\lambda_1) = C(\lambda_2 - k\lambda_1) = 0.$$

Provided $|B| + |C| \neq 0$ there is no way to extend symmetry of Equation (6.20). If, on the contrary, $B = C = 0$, then $F = Au^{-1}u_x^2 (A \neq 0, -1)$ and we obtain the equation equivalent to (6.15). Under $k = 0$ system (6.22) takes the form

$$2(A + 1)x^2h' = \lambda_2 B, \quad x^2 Bh' + h''x^3 = 2\lambda_2 C.$$

Hence

$$h = -\frac{1}{2}\lambda_2(A + 1)^{-1}Bx^{-1} + C_1, \quad C_1 \in \mathbb{R}, \quad C = \frac{B^2 - 2B}{4(A + 1)}.$$

In this case Equation (6.20) does admit additional symmetry operator

$$\partial_x - \frac{B}{2(A + 1)}x^{-1}u\partial_u$$

but the change of variables

$$\bar{t} = t, \quad \bar{x} = x, \quad u = |x|^\nu v, \quad \nu = -\frac{B}{2(A + 1)},$$

reduces it to the form (6.15).

So Equation (6.20) admits wider symmetry group iff it is either belongs to the class of (6.1) or is equivalent to (6.15).

To finalize the procedure of group classification of Equation (2.7) we need to consider invariant equations obtained in the previous section that contain arbitrary functions of two variables.

7.3.2. Classification of Equations with Arbitrary Functions of Two Variables

In the case of equations with arbitrary functions of two variables the standard Lie–Ovsyannikov method is inefficient and we apply our classification algorithm. To this end, we compute extensions of three-dimensional solvable Lie algebras to all possible realizations of four-dimensional solvable Lie algebras. The subsequent step will be to check which of the obtained realizations are symmetry algebras of nonlinear equations of the form (2.7). In what follows we utilize the results of the paper [44], where all inequivalent (within the action of inner automorphism group) four-dimensional solvable abstract Lie algebras are given.

We give full computation details for the case of $A_{3,6}$ -invariant equations. As shown in the previous sub-section, there are two inequivalent $A_{3,6}$ -invariant equations, namely,

$$\begin{aligned} A_{3,6}^4 &= \langle e^{-t}\partial_u, e^t\partial_u, \partial_t + k\partial_x \rangle \\ &(k \geq 0) : F = u + G(\eta, u_x), \quad \eta = x - kt; \\ A_{3,6}^5 &= \langle |t|^{-\frac{1}{2}}\partial_u, |t|^{\frac{3}{2}}\partial_u, t\partial_t + x\partial_x + \frac{1}{2}u\partial_u \rangle : \\ &F = \frac{3}{4}t^{-2}u + |t|^{-\frac{3}{2}}G(\xi, \omega), \quad \xi = tx^{-1}, \quad \omega = x^{-1}u_x^2. \end{aligned}$$

According to [44] the algebra $A_{3,6}$ is the subalgebra of the following four-dimensional solvable Lie algebras: $2A_{2,2}$, $A_{3,6} \oplus A_1$; $A_{4,2}(q = -1)$; $A_{4,8}(q = -\frac{1}{2})$.

Algebra $2A_{2,2}$. The algebra $2A_{2,2} = \langle e_1, e_2, e_3, e_4 \rangle$ is defined by the following commutation relations (note that we give non-zero relations only):

$$[e_1, e_2] = [e_1, e_4] = [e_2, e_3] = [e_2, e_4] = 0, \quad [e_1, e_2] = e_2, [e_3, e_4] = e_4.$$

It contains a subalgebra $A_{3,6} = \langle e_1 - e_3, e_2, e_4 \rangle$. That is why, we can choose the basis operators of the realization of $A_{3,6}$ as $e_1 - e_3, e_2, e_4$. Next, we take an arbitrary operator of the form (2.4) as $e_1 + e_3$ and require for the commutation relations

$$[e_1 - e_3, e_1 + e_3] = 0, \quad [e_1 + e_3, e_2] = e_2, \quad [e_1 + e_3, e_4] = e_4 \quad (6.23)$$

to hold.

Realization $A_{3,6}^4$. In this case

$$\begin{aligned} e_1 - e_3 &= -\partial_t - k\partial_x, \quad e_2 = e^{-t}\partial_u, \quad e_4 = e^t\partial_u, \\ e_1 + e_3 &= (\lambda t + \lambda_1)\partial_t + (\lambda x + \lambda_2)\partial_x + (hu + r)\partial_u. \end{aligned}$$

It follows from (6.23) that

$$\lambda = \lambda_1 = 0, \quad r = \gamma = \gamma(\eta), \quad h = -1.$$

Using the change of variables

$$\bar{t} = t, \quad \bar{x} = x, \quad v = u + \Lambda(\eta),$$

where $\Lambda = \Lambda(\eta)$ is a solutions of equation $\lambda_2\Lambda' + \Lambda = -\gamma$, we simplify the operator $e_1 + e_3$ to become

$$e_1 + e_3 = \alpha\partial_x - u\partial_u, \quad \alpha \in \mathbb{R}.$$

Requiring invariance under the above operator yields that $\alpha \neq 0$ (otherwise G would be linear in u_x). With this result in hand we rewrite the invariant equation to become

$$G = \exp(-\alpha^{-1}\eta)H(\omega), \quad \omega = \exp(\alpha^{-1}\eta)u_x.$$

Thus we arrive at the following realization of the algebra $2A_{2,2}$:

$$\langle e^{-t}\partial_u, e^t\partial_u, \partial_t + k\partial_x, \alpha\partial_x - u\partial_u \rangle \quad (k \geq 0, \alpha \neq 0).$$

This algebra is admitted by the equation

$$u_{tt} = u_{xx} + u + \exp(-\alpha^{-1}\eta)G(\omega), \quad \eta = x - kt, \quad \omega = \exp(\alpha^{-1}\eta)u_x.$$

If the function G ($G_{\omega\omega} \neq 0$) is arbitrary, then the obtained realization is the maximal symmetry algebra of the equation under study. What is more, there is no such G that the above equation admits a wider invariance algebra.

Realization $A_{3,6}^5$. In this case

$$e_1 - e_3 = -t\partial_t - x\partial_x - \frac{1}{2}u\partial_u, \quad e_2 = |t|^{-\frac{1}{2}}\partial_u,$$

$$e_1 + e_3 = (\lambda_t + \lambda_1)\partial_t + (\lambda_x + \lambda_2)\partial_x + (hu + r)\partial_u, \quad e_4 = |t|^{\frac{3}{2}}\partial_u.$$

It follows from commutation relations (6.23) that $\lambda_1 = \lambda_2 = \lambda = 0$, $h = -1$, $r = |t|^{\frac{1}{2}}\gamma(\xi)$, $\xi = tx^{-1}$.

Making the change of variables

$$\bar{t} = t, \quad \bar{x} = x, \quad v = u - |t|^{\frac{1}{2}}\gamma(\xi)$$

we get $r = 0$ in $e_1 + e_3$. Consequently, without loss of generality we can choose $e_1 + e_3 = -u\partial_u$. Requiring for the $A_{3,6}^5$ -invariant equation to admit the operator $e_1 + e_3$ yields the constraint $2\omega G_\omega = -G$, whence $G = |\omega|^{-\frac{1}{2}}H(\xi)$. Consequently, the function F is linear in u_x . This means that $A_{3,6}^5$ does not admit extension to a realization $2A_{2,2}$ that can be a symmetry algebra of an essentially nonlinear equation of the form (2.7).

Algebra $A_{3,6} \oplus A_1$. What we need to do here is to supplement the set of operators e_1, e_2, e_3 forming the basis of $A_{3,6}$ by the operator e_4 of the form (2.4) and verify the commutation relations

$$[e_1, e_4] = [e_2, e_4] = [e_3, e_4] = 0. \tag{6.24}$$

Realization $A_{3,6}^4$. It follows from (6.24) that $h = \lambda = \lambda_1 = 0$, $r = \gamma(\eta)$, $\eta = x - kt$ in the operator e_4 so that

$$e_4 = \alpha\partial_x + \gamma(\eta)\partial_u, \quad \alpha \in \mathbb{R}.$$

If $\alpha \neq 0$, then e_4 is equivalent to ∂_x . Hence we get two possible realizations of the algebra $A_{3,6} \oplus A_1$:

$$\langle e^{-t}\partial_u, e^t\partial_u, \partial_t, \partial_x \rangle;$$

$$\langle e^{-t}\partial_u, e^t\partial_u, \partial_t + k\partial_x, \gamma(\eta)\partial_u \rangle.$$

Analyzing the above realizations we come to conclusion that the second one cannot be invariance algebra of a nonlinear equation of the form (2.7). The first realization is the maximal invariance algebra of the (6.19), if G is an arbitrary function.

Realization $A_{3,6}^3$. It follows from (6.24) that $\lambda_1 = \lambda_2 = h = \lambda = 0$ and

$$r = |t|^{\frac{1}{2}}\gamma(\xi), \quad \xi = tx^{-1},$$

so that the operator e_4 necessarily takes the form $e_4 = |t|^{\frac{1}{2}}\gamma(\xi)\partial_u$. As the straightforward verification shows the so obtained realization cannot be invariance algebra of a nonlinear equation of the form (2.7).

Algebra $A_{4,2}$ ($q = -1$). We need to supplement the set of operators e_1, e_2, e_4 forming the basis of $A_{3,6}$ by the operator e_3 of the form (2.4) so that the following commutation relation hold

$$[e_1, e_3] = [e_2, e_3] = 0, \quad [e_3, e_4] = e_2 + e_3. \tag{6.25}$$

Realization $A_{3,6}^4$. In this case

$$e_1 = e^{-t} \partial_u, \quad e_2 = e^t \partial_u, \quad e_4 = -\partial_t - k \partial_x$$

and it follows from (6.25) that the coefficients of e_3 satisfy equations $h = \lambda = \lambda_1 = \lambda_2 = 0$, the function r being a solution of the equation

$$r_t + kr_x = r + e^t.$$

Further analysis shows that this realization cannot be invariance algebra of nonlinear equation of the form (2.7).

Realization $A_{3,6}^5$. In this case

$$e_1 = |t|^{-\frac{1}{2}} \partial_u, \quad e_2 = |t|^{\frac{3}{2}} \partial_u, \quad e_4 = -t \partial_t - x \partial_x - \frac{1}{2} u \partial_u.$$

It follows from Equation (6.25) that the coefficients of the operator e_3 satisfy equations $\lambda = \lambda_1 = \lambda_2 = h = 0$ and the function r is a solution of the equation

$$tr_t + xr_x = \frac{3}{2}r + |t|^{\frac{3}{2}}.$$

Further analysis shows that the so obtained realization cannot be invariance algebra of a nonlinear equation of the form (2.7).

Algebra $A_{4,8}$ ($q = -\frac{1}{2}$). We need to supplement the set of operators e_1, e_3, e_4 forming the basis of $A_{3,6}$ by the e_2 of the form (2.4) in order to satisfy the commutation relations

$$[e_1, e_2] = 0, \quad [e_2, e_3] = e_1, \quad [e_2, e_4] = e_2. \quad (6.26)$$

Realization $A_{3,6}^4$. In this case

$$e_1 = e^{-t} \partial_u, \quad e_3 = e^t \partial_u, \quad e_4 = \frac{1}{2} \partial_t + \frac{1}{2} k \partial_x$$

and the second commutation relation yields the false equality $1 = 0$.

Realization $A_{3,6}^5$. In this case

$$e_1 = |t|^{-\frac{1}{2}} \partial_u, \quad e_2 = |t|^{\frac{3}{2}} \partial_u, \quad e_4 = -t \partial_t - x \partial_x - \frac{1}{2} u \partial_u$$

and again the second commutation relation from (6.26) cannot be satisfied.

Consequently, there is no extension of the realization of $A_{3,6}$ to a realization of the algebra $A_{3,8}$ ($q = -\frac{1}{2}$).

The remaining equations containing arbitrary functions of two variables are handled in a similar way. The results can be summarized as follows

- if the functions contained in the equations under study are arbitrary, then the corresponding realizations are their maximal invariance algebras, and
- except for Equation (6.14), all the equations in question do not allow for extension of their symmetry.

Below we give the complete list of invariant equations obtained through group analysis of equations with arbitrary functions of two variables.

7.3.3. Equations Invariant Under Four-dimensional Solvable Lie Algebras

$A_{2,2} \oplus 2A_1$ -invariant equations

- 1) $\langle \partial_x, \partial_t + u\partial_u, e^t\partial_u, e^{-t}\partial_u \rangle : F = u + e^t G(\omega), \omega = u^{-t}u_x;$
- 2) $\langle \frac{1}{2k}(\partial_t + k\partial_x), e^{x+kt}\partial_u, e^\eta\partial_u, \partial_x + u\partial_u \rangle (k > 0, \eta = x - kt) :$
 $F = (k^2 - 1)u + e^\eta G(\omega), \omega = e^{-\eta}(u_x - u).$

$2A_{2,2}$ -invariant equations

- 1) $\langle \partial_t + \epsilon u\partial_u, \partial_x, e^{x+kt}\partial_u, e^{x-kt}\partial_u \rangle (\epsilon = 0, 1; k > 0) :$
 $F = (k^2 - 1)u + e^{\epsilon t} G(\omega), \omega = e^{-\epsilon t}(u_x - u);$
- 2) $\langle \alpha\partial_x - u\partial_u, \partial_t + k\partial_x, e^{-t}\partial_u, e^t\partial_u \rangle (k \geq 0, \alpha > 0) :$
 $F = u + \exp(-\alpha^{-1}\eta)G(\omega), \eta = x - kt, \omega = \exp(\alpha^{-1}\eta)u_x.$

$A_{3,3} \oplus A_1$ -invariant equations

- 1) $\langle \partial_t, \partial_x, \partial_u, t\partial_u \rangle : F = G(u_x).$

$A_{3,4} \oplus A_1$ -invariant equations

- 1) $\langle \partial_u, \partial_x, t\partial_t + x\partial_x + (u + x)\partial_u, t\partial_u \rangle :$
 $F = t^{-1}G(\omega), \omega = u_x - \ln|t|;$
- 2) $\langle \partial_t + u\partial_u, \partial_x, t\partial_u, \partial_u \rangle : F = e^t G(\omega), \omega = e^{-t}u_x;$
- 3) $\langle x^{-1}\partial_u, \partial_x - x^{-1}(u + \ln|x|)\partial_u, t\partial_t + x\partial_x, tx^{-1}\partial_u \rangle :$
 $F = 2x^{-1}u_x + x^{-2} + t^{-1}x^{-1}G(\omega), \omega = xu_x + u - \ln|tx^{-1}|.$

$A_{3,5} \oplus A_1$ -invariant equations

- 1) $\langle \partial_x, \partial_u, t\partial_t + x\partial_x + u\partial_u, t\partial_u \rangle : F = t^{-1}G(u_x);$
- 2) $\langle x^{-1}\partial_u, \partial_x - x^{-1}u\partial_u, t\partial_t + x\partial_x, tx^{-1}\partial_u \rangle :$
 $F = -2x^{-2}u + 2t^{-1}(u_x + x^{-1}u) \ln|t(u_x + x^{-1}u)|$
 $+ t^{-1}(u_x + x^{-1}u)G(\omega), \omega = xu_x + u.$

$A_{3,6} \oplus A_1$ -invariant equations

- 1) $\langle \partial_x, t\partial_u, t\partial_t + x\partial_x, \partial_u \rangle : F = t^{-2}G(\omega), \omega = t^{-1}u_x;$
- 2) $\langle \partial_t, \partial_x, e^t\partial_u, e^{-t}\partial_u \rangle : F = u + G(u_x).$

 $A_{3,7} \oplus A_1$ -invariant equations

- 1) $\langle \exp\left(-\frac{1}{2}(1-q)t\right)\partial_u, \exp\left(\frac{1}{2}(1-q)t\right)\partial_u, \partial_t + \frac{1}{2}(1+q)u\partial_u, \partial_x \rangle$
 $(q \neq 0, \pm 1) : F = \frac{1}{4}(1-q)^2u + \exp\left(\frac{1}{2}(1+q)t\right)G(\omega),$
 $\omega = \exp\left(-\frac{1}{2}(1+q)t\right)u_x;$
- 2) $\langle \partial_x, |t|^{\frac{1}{2}(1-q)}\partial_u, |t|^{\frac{1}{2}(1+q)}\partial_u, t\partial_t + x\partial_x + \frac{1}{2}(1+q)u\partial_u \rangle$
 $(q \neq 0, \pm 1) : F = \frac{1}{4}(q^2 - 1)t^{-2}u + |t|^{\frac{1}{2}(q-3)}G(\omega),$
 $\omega = |t|^{\frac{1}{2}(1-q)}u_x;$
- 3) $\langle |t|^{-\frac{1}{q}}|\xi|^{\frac{q+1}{2q}}\partial_u, \partial_x - \frac{1+q}{2q}x^{-1}u\partial_u, -q(t\partial_t + x\partial_x), |\xi|^{\frac{1+q}{2q}}\partial_u \rangle$
 $(q \neq 0, \pm 1) : F = \left[\frac{1-q^2}{4q^2}(t^{-2} + x^{-2}) \right]u + \frac{1+q}{q}x^{-1}u_x$
 $+ t^{-2}|\xi|^{\frac{1+q}{2q}}G(\omega), \xi = tx^{-1}, \omega = |\xi|^{\frac{q-1}{2q}} \left[xu_x + \frac{q+1}{2q}u \right].$

 $A_{3,8} \oplus A_1$ -invariant equations

- 1) $\langle \sin t\partial_u, \cos t\partial_u, \partial_t, \partial_x \rangle : F = -u + G(u_x).$

 $A_{3,9} \oplus A_1$ -invariant equations

- 1) $\langle \sin t\partial_u, \cos t\partial_u, \partial_t + qu\partial_u, \partial_x \rangle (q > 0) :$
 $F = -u + e^{qt}G(\omega), \omega = e^{-qt}u_x.$

 $A_{4,1}$ -invariant equations

- 1) $\langle \partial_u, -t\partial_u, \partial_x, \partial_t - tx\partial_u \rangle : F = G(\omega), \omega = u_x + \frac{1}{2}t^2;$
- 2) $\langle \partial_u, -t\partial_u, \alpha\partial_x + \frac{1}{2}t^2\partial_u, \partial_t + kx\partial_x \rangle (k \geq 0, \alpha > 0) :$
 $F = \alpha^{-1}(x - kt) + G(u_x).$

$A_{4,2}$ -invariant equations

$$1) \langle |t|^{1-\frac{1}{2}q} \partial_u, |t|^{\frac{1}{2}q} \partial_u, \partial_x, t \partial_t + x \partial_x + \left[\left(1 + \frac{1}{2}q\right) u + x |t|^{\frac{1}{2}q} \right] \partial_u \rangle$$

$$(q \neq 0, 1) : F = \frac{1}{4}q(q-2)t^{-2}u + |t|^{\frac{1}{2}(q-3)}G(\omega),$$

$$\omega = |t|^{\frac{1}{2}(1-q)}u_x - 2|t|^{\frac{1}{2}};$$

$$2) \langle \partial_x, \sqrt{|t|} \partial_u, \sqrt{|t|} \ln |t| \partial_u, t \partial_t + x \partial_x + \left(q + \frac{1}{2}\right) u \partial_u \rangle$$

$$(q \neq 0) : F = -\frac{1}{4}t^{-2}u + |t|^{q-\frac{3}{2}}G(\omega), \quad \omega = |t|^{\frac{1}{2}-q}u_x.$$

$A_{4,3}$ -invariant equations

$$1) \langle \partial_x, |t|^{\frac{1}{2}} \partial_u, -|t|^{\frac{1}{2}} \ln |t| \partial_u, t \partial_t + x \partial_x + \frac{1}{2}u \partial_u \rangle :$$

$$F = -\frac{1}{4}t^{-2}u + |t|^{-\frac{3}{2}}G(\omega), \quad \omega = |t|^{\frac{1}{2}}u_x;$$

$$2) \langle \partial_x, t \partial_u, \partial_u, t \partial_t + x \partial_x \rangle : F = t^{-2}G(\omega), \quad \omega = tu_x;$$

$$3) \langle e^{kt} \partial_u, \partial_t + ku \partial_u, \beta \partial_x + te^{kt} \partial_u, e^{-kt} \partial_u \rangle (k \neq 0, \beta > 0) :$$

$$F = k^2u + 2k\beta^{-1}xe^{kt} + e^{kt}G(\omega), \quad \omega = e^{-kt}u_x;$$

$$4) \langle e^{x+kt} \partial_u, e^\eta \partial_u, \alpha(\partial_x + u \partial_u) + 2kte^\eta \partial_u, -\frac{1}{2k}(\partial_t + k \partial_x) \rangle$$

$$(\alpha \neq 0, k > 0) : F = (k^2 - 1)u - 4k^2\alpha^{-1}\eta e^\eta + e^\eta G(\omega),$$

$$\omega = e^{-\eta}(u_x - u), \eta = x - kt.$$

$A_{4,4}$ -invariant equations

$$1) \langle |t|^{\frac{1}{2}} \partial_u, -|t|^{\frac{1}{2}} \ln |t| \partial_u, \partial_x, t \partial_t + x \partial_x + \left[\frac{3}{2}u - x |t|^{\frac{1}{2}} \ln |t| \right] \partial_u \rangle :$$

$$F = \frac{1}{4}t^{-2}u + |t|^{-\frac{1}{2}}G(\omega), \quad \omega = |t|^{-\frac{1}{2}}u_x + \frac{1}{2} \ln^2 |t|.$$

$A_{4,5}$ -invariant equations

$$1) \langle \partial_x, |t|^{m-\alpha} \partial_u, |t|^{1-m+\alpha} \partial_u, t \partial_t + x \partial_x + mu \partial_u \rangle$$

$$(m \neq \frac{1}{2}(1 + \alpha), \frac{1}{2} + \alpha; \alpha \neq 0) :$$

$$F = (m - \alpha)(m - \alpha - 1)t^{-2}u + |t|^{m-2}G(\omega), \quad \omega = |t|^{1-m}u_x.$$

A_{4.6}-invariant equations

$$1) \langle \partial_x, |t|^{\frac{1}{2}} \sin(q^{-1} \ln |t|) \partial_u, |t|^{\frac{1}{2}} \cos(q^{-1} \ln |t|) \partial_u, qt \partial_t + qx \partial_x \\ \left(\frac{1}{2}q + p \right) u \partial_u \rangle (q \neq 0, p \geq 0) :$$

$$F = - \left(\frac{1}{4} + q^{-2} \right) t^{-2} u + |t|^{q^{-1}(p - \frac{3}{2}q)} G(\omega), \quad \omega = |t|^{q^{-1}(\frac{1}{2}q - p)} u_x.$$

A_{4.7}-invariant equations

$$1) \langle \partial_u, -t \partial_u, \partial_t + k \partial_x, t \partial_t + x \partial_x + \left(2u - \frac{1}{2}t^2 \right) \partial_u \rangle (k \geq 0) :$$

$$F = -\ln |\eta| + G(\omega), \quad \omega = \eta^{-1} u_x, \quad \eta = x - kt.$$

A_{4.8}-invariant equations

$$1) \langle \partial_t + \epsilon u \partial_u, \partial_x, e^x \partial_u, t e^x \partial_u \rangle (\epsilon = 0; 1) :$$

$$F = -u + e^{\epsilon t} G(\omega), \quad \omega = e^{-\epsilon t} (u_x - u);$$

$$2) \langle |x|^{m-q} \partial_u, \partial_t, t |x|^{m-q} \partial_u, t \partial_t + x \partial_x + mu \partial_u \rangle (q \neq 0, m \in \mathbb{R}) :$$

$$F = -(m - q)(m - q - 1)x^{-2}u + |x|^{m-2}G(\omega),$$

$$\omega = |x|^{1-m}[u_x - (m - q)x^{-1}u];$$

$$3) \langle \partial_t + k \partial_x, \partial_u, t \partial_u, t \partial_t + x \partial_x + qu \partial_u \rangle (k > 0, q \in \mathbb{R}) :$$

$$F = |\eta|^{q-2}G(\omega), \quad \omega = |\eta|^{1-q}u_x, \quad \eta = x - kt;$$

$$4) \langle x^{-1} \partial_u, \partial_t + \partial_x - x^{-1}u \partial_u, tx^{-1} \partial_u, t \partial_t + x \partial_x \rangle :$$

$$F = 2x^{-1}u_x + x^{-1}(t - x)^{-1}G(\omega), \quad \omega = xu_x + u;$$

$$5) \langle \partial_u, -t \partial_u, \partial_t + k \partial_x + u \partial_u, \alpha \partial_x + u \partial_u \rangle (\alpha \neq 0, k \geq 0) :$$

$$F = \exp(\alpha^{-1}\eta + t)G(\omega), \quad \omega = \exp(-\alpha^{-1}\eta - t)u_x, \quad \eta = x - kt.$$

A_{4.10}-invariant equations

$$1) \langle \sin t \partial_u, \cos t \partial_u, \partial_x + u \partial_u, \partial_t + k \partial_x \rangle (k \geq 0) :$$

$$F = -u + e^\eta G(\omega), \quad \omega = e^{-\eta} u_x, \quad \eta = x - kt.$$

In the above formulas $G = G(\omega)$ is an arbitrary function satisfying the condition $F_{u_x u_x} \neq 0$.

8. Symmetry Reduction and Solutions of Nonlinear Wave Equations

Among the various applications of Lie symmetry groups the most prominent and remarkable one is a possibility to construct exact solutions of nonlinear PDEs. The basic idea is reducing multi-dimensional differential equations to ordinary differential equations via special ansatzes (invariant solutions). A regular (but not the only!) way to derive those ansatzes is to utilize symmetry group admitted by the equation under study (for more details see, e.g., [14, 15]). Though the obtained

ordinary differential equations are, as a rule, nonlinear, they possess in many cases a residual symmetry allowing for constructing their general or particular solutions. Inserting the latter into the corresponding ansatz yields the exact solution of initial nonlinear multi-dimensional PDE. This method is often referred to in the literature as symmetry reduction of PDEs.

The majority of papers on exact solutions of Equation (1.1) deal with the PDEs of the form

$$u_{tx} = f(u), \tag{7.1}$$

where f is a smooth function. The particular cases of the above equation are

- the Liouville equation ($f(u) = e^u$),
- the Bonnet (or sin-Gordon) equation ($f(u) = \sin u$),
- the sinh-Gordon equation ($f(u) = \sinh u$),
- the Tzitzeica equation ($f(u) = e^u + e^{-2u}$).

It is a common knowledge that the Liouville equation is integrated in a closed form. The sin/sinh-Gordon equations are integrable by the inverse scattering method. Some exact solutions of the Tzitzeica equation were found in [4, 5].

A number of explicit solutions of Equation (7.1) different from those mentioned above have been constructed in [45] (see, also [40] and the references therein). The broad classes of exact solutions of Equation (7.1) under $f(u) = \lambda u^k$ are obtained in [46]. The case $f = a \sin(\lambda u) + b \sin(\frac{1}{2}\lambda u)$ has been studied in [5, 12, 47].

In what follows we concentrate on constructing solutions of nonlinear wave equation (2.7) having the richest symmetry properties. To this end we apply the symmetry reduction method.

To perform reduction of PDEs (2.7) to ordinary differential equations we need to describe all inequivalent one-dimensional subalgebras of the symmetry algebras of the equations under study. What is more, basis operators of these algebras

$$\tau(t, x, u)\partial_t + \xi(t, x, u)\partial_x + \eta(t, x, u)\partial_u,$$

have to obey the following restriction [14]:

$$|\tau| + |\xi| \neq 0 \tag{7.2}$$

in some open domain Ω of the space $V = \mathbb{R}^2 \times \mathbb{R}^1$ of independent $\mathbb{R}^2 = \langle t, x \rangle$ and dependent $\mathbb{R}^1 = \langle u \rangle$ variables.

As we proved in the previous sections, equations

$$u_{tt} = u_{xx} - u^{-1}u_x^2; \tag{7.3}$$

$$u_{tt} = u_{xx} + e^{ux}; \tag{7.4}$$

$$u_{tt} = u_{xx} + m \ln |u_x| \ (m > 0); \tag{7.5}$$

$$u_{tt} = u_{xx} + |u_x|^k \ (k \neq 0, 1). \tag{7.6}$$

enjoy the highest symmetry properties amongst PDEs of the form (2.7).

We consider in some detail the main steps of the symmetry reduction algorithm for Equation (7.3). To classify one-dimensional subalgebras we utilize the method

suggested in [44] and the lists of one-dimensional subalgebras of four-dimensional subalgebras obtained in [44].

Equation (7.3) admits the algebra

$$A_5^1 = (A_{3,3} \oplus A_1) \ni \langle e_5 \rangle,$$

where $A_{3,3} = \langle e_1, e_2, e_3 \rangle = \langle u\partial_u, \partial_x, xu\partial_u \rangle$, $A_1 = \langle e_4 \rangle = \langle \partial_t \rangle$, $e_5 = t\partial_t + x\partial_x$.

In what follows we need the commutation relations of the basis operators of the algebra $A_{3,3} \oplus A_1$ with the operator e_5 :

$$[e_1, e_5] = 0, \quad [e_2, e_5] = e_2, \quad [e_3, e_5] = -e_3, \quad [e_4, e_5] = e_4.$$

According to [44] the one-dimensional subalgebras of $A_{3,3} \oplus A_1$ defined within the action of inner automorphism group of this algebra are $\langle e_1 \rangle$, $\langle e_1 + \alpha e_4 \rangle$, $\langle e_4 \rangle$, $\langle e_2 \rangle$, $\langle e_2 + \alpha e_4 \rangle$, $\langle e_3 \rangle$, $\langle e_3 + \alpha e_4 \rangle$, $\langle e_2 + \beta e_3 \rangle$, $\langle e_2 + \beta e_3 + \alpha e_4 \rangle$ ($\alpha, \beta \neq 0$). The above subalgebras can be further simplified by using transformation group generated by the operator e_5 . For example, using the Campbell–Hausdorff formula we transform $e_1 + \alpha e_4$ as follows:

$$\exp(\theta e_5)(e_1 + \alpha e_4) \exp(-\theta e_5) = e_1 + \alpha e^\theta e_4.$$

Consequently, choosing $\theta = -\ln|\alpha|$ we simplify $e_1 + \alpha e_4$ to become $e_1 \pm e_4$. Similarly, we prove that we can put $\alpha = \pm 1$ in $e_3 + \alpha e_4$ and $\beta = \pm 1$ in $e_2 + \beta e_3$, $e_3 + \beta e_3 + \alpha e_4$.

To complete classification of one-dimensional subalgebras we have to describe all inequivalent subalgebras with non-zero projection on the basis operator e_5 , i.e., subalgebras of the form

$$\Lambda = e_5 + \alpha_1 e_1 + \alpha_2 e_2 + \alpha_3 e_3 + \alpha_4 e_4, \quad \alpha_1, \alpha_2, \alpha_3, \alpha_4 \in \mathbb{R}. \tag{7.7}$$

Utilizing the automorphism $\exp(\theta e_4)$ with a properly chosen θ we have $\alpha_4 = 0$ in (7.7). Next applying the transformation $\exp(\theta_1 e_2 + \theta_2 e_3)$ to operator (7.7) reduces it to one of the following operators $e_5, e_5 + \alpha e_1$ ($\alpha \neq 0$).

So the list of the one-dimensional subalgebras of the five-dimensional algebra A_5^1 determined within the action of inner automorphism group is exhausted by the following algebras: $\langle e_1 \rangle$, $\langle e_1 \pm e_4 \rangle$, $\langle e_4 \rangle$, $\langle e_2 \rangle$, $\langle e_2 + \alpha e_4 \rangle$, $\langle e_3 \rangle$, $\langle e_3 \pm e_4 \rangle$, $\langle e_2 \pm e_3 \rangle$, $\langle e_2 \pm e_3 + \alpha e_4 \rangle$, $\langle e_5 \rangle$, $\langle e_5 + \alpha e_1 \rangle$ ($\alpha \neq 0$). By direct verification we prove that the basis operators of the algebras $\langle e_1 \rangle$, $\langle e_3 \rangle$ do not satisfy condition (7.2).

Finally, we make use of the fact that the discrete groups of transformations

$$\begin{aligned} \bar{t} &= -t, \quad \bar{x} = x, \quad v = u; \\ \bar{t} &= t, \quad \bar{x} = -x, \quad v = u; \\ \bar{t} &= -t, \quad \bar{x} = -x, \quad v = u, \end{aligned}$$

also belong to the equivalence group of (7.3). Using the above transformations enables to further simplify the optimal system of inequivalent subalgebras

$$\begin{aligned} &\langle e_1 + e_4 \rangle, \langle e_4 \rangle, \langle e_2 \rangle, \langle e_2 + \alpha e_4 \rangle, \langle e_3 + e_4 \rangle, \\ &\langle e_2 \pm e_3 \rangle, \langle e_2 \pm e_3 + \alpha e_4 \rangle, \langle e_5 \rangle, \langle e_5 + \alpha e_1 \rangle \quad (\alpha > 0). \end{aligned} \tag{7.8}$$

The second step of the method of symmetry reduction is constructing the complete set of invariants $f(t, x, u)$ for each inequivalent one-dimensional subalgebra. As a

typical example, we consider the case of the subalgebra $\langle e_1 + e_4 \rangle$. To construct its invariants we need to integrate the first-order PDE

$$(e_1 + e_4) \circ F(t, x, u) = 0$$

or

$$uF_u + F_t = 0.$$

The complete set of first integrals of the above equation is $\omega_1 = x, \omega_2 = e^{-t}u$. Hence we get the general form of invariant solution (ansatz) $\omega_2 = \varphi(\omega_1)$. Solving this equation with respect to u we finally have

$$u = e^t \varphi(x). \tag{7.9}$$

Inserting (7.9) into (7.3) yields ordinary differential equation for unknown function φ

$$\varphi'' - \varphi^{-1}(\varphi')^2 - \varphi = 0,$$

which is easily integrated

$$\varphi = \exp \left[\frac{1}{2}x^2 + C_1x + C_2 \right], \quad C_1, C_2 \in \mathbb{R}.$$

Inserting the so obtained expression for ϕ into ansatz (7.9) yields the explicit form of invariant solution of equation (7.3)

$$u = \exp \left[t + \frac{1}{2}x^2 + C_1x + C_2 \right], \quad C_1, C_2 \in \mathbb{R}.$$

The full list of so obtained exact solutions of (7.3) is given in the Appendix. In a similar way we perform symmetry reduction and construct exact solutions of Equations (7.4)–(7.6). These solutions are also listed in Appendix.

9. Concluding Remarks

Let us briefly summarize the results obtained in this paper.

We prove that the problem of group classification of the general quasi-linear hyperbolic type Equation (1.1) reduces to classifying equations of more specific forms

- I. $u_{tt} = u_{xx} + F(t, x, u, u_x), \quad F_{u_x u_x} \neq 0;$
- II. $u_{tt} = u_{xx} + g(t, x, u)u_x + f(t, x, u), \quad g_u \neq 0;$
- III. $u_{tx} = g(t, x)u_x + f(t, x, u), \quad g_x \neq 0, \quad f_{uu} \neq 0;$
- IV. $u_{tx} = f(t, x, u), \quad f_{uu} \neq 0.$

The cases of PDEs that are essentially nonlinear in u_x (the class of PDEs I) and either linear in u_x or do not depend on u_x (the classes II–IV) need to be considered separately.

If we denote as \mathcal{DE} the set of PDEs II–III, then the results of application of our algorithm for group classification of equations I–IV can be summarized as follows.

- We perform complete group classification of the class \mathcal{DE} . We prove that the Liouville equation has the highest symmetry properties among equations from \mathcal{DE} . Next we prove that the only equation belonging to this class and admitting the four-dimensional invariance algebra is the nonlinear d'Alembert equations. It is established that there are 12 inequivalent equations from \mathcal{DE} invariant under three-dimensional Lie algebras. We give the lists of all inequivalent equations from \mathcal{DE} that admit one- and two-dimensional symmetry algebras.
- We have studied the structure of invariance algebras admitted by nonlinear equations from the class I. It is proved, in particular, that their invariance algebras are necessarily solvable.
- We perform complete group classification of nonlinear equations from the class of PDEs I. We prove that the highest symmetry algebras admitted by those equations are five-dimensional and construct all inequivalent classes of equations invariant with respect to five-dimensional Lie algebras. We also construct all inequivalent equations of the form I admitting one-, two-, three- and four-dimensional Lie algebras.

The results of group classification of the class of nonlinear wave equation (1.1) are utilized for constructing their explicit solutions. Namely, we perform symmetry reduction of all Equation (1.1) admitting five-dimensional symmetry algebras to ordinary differential equations and constructed multi-parameter families of their exact solutions.

Appendix. Exact solutions of Equations (7.3)–(7.6)

Below we present the lists of invariant solutions of nonlinear wave equations (7.3)–(7.6), which are obtained using the symmetry reduction approach.

Exact solutions of (7.3):

$$\langle e_1 + e_4 \rangle : u = \exp \left[t + \frac{1}{2}x^2 + C_1x + C_2 \right],$$

$$\langle e_4 \rangle : u = \exp[C_1x + C_2],$$

$$\langle e_2 \rangle : u = C_1t + C_2,$$

$$\langle e_2 + \alpha e_4 \rangle, (\alpha > 0) : u = C_2|t - \alpha x + C_1|^{1-\alpha^2},$$

$$\langle e_3 + e_4 \rangle : u = \exp \left[tx + \frac{1}{12}x^4 + C_1x + C_2 \right],$$

$$\langle e_2 - e_3 \rangle : u = C_1 \exp \left(-\frac{1}{2}x^2 \right) \cos(t + C_2),$$

$$\langle e_2 + e_3 \rangle : u = C_1 \exp\left(\frac{1}{2}x^2\right) \cosh(t + C_2),$$

$$\langle e_2 + e_3 + e_4 \rangle : u = \exp\left[C_1 + \frac{1}{2}x^2 \pm (t - x)\right],$$

$$\langle e_2 - e_3 + \alpha e_4 \rangle (\alpha \neq 1, \alpha > 0) : u = C_2 \exp\left(-\frac{1}{2}x^2\right) \left[\cos\left(C_1 + \frac{t - \alpha x}{\alpha^2 - 1}\right)\right]^{1 - \alpha^2},$$

$$\langle e_2 + e_3 + \alpha e_4 \rangle (\alpha \neq 1, \alpha > 0) : u = C_2 \exp\left(-\frac{1}{2}x^2\right) \left[\cosh\left(C_1 + \frac{t - \alpha x}{1 - \alpha^2}\right)\right]^{1 - \alpha^2},$$

$$\langle e_5 \rangle : u = \exp\left\{\int^\xi \left(C_1(\eta^2 - 1) + \frac{1}{4}(\eta^2 - 1) \ln \left|\frac{1 + \eta}{1 - \eta}\right| - \frac{1}{2}\eta\right)^{-1} d\eta + C_2\right\}, \quad \xi = tx^{-1},$$

$$\langle e_5 + e_1 \rangle : u = t \exp\left\{\int^\xi \left[C_1\eta^2 + \frac{1}{2}\eta \ln \left|\frac{1 + \eta}{1 - \eta}\right| - \eta\right]^{-1} d\eta + C_2\right\},$$

$$\langle e_5 + \alpha e_1 \rangle : u = |t|^\alpha \exp\left[\int^\xi g(\eta) d\eta + C_1\right], \quad \text{where } \xi = tx^{-1} \text{ and } g(\xi) \text{ is a solution of}$$

the Riccati equation, $\xi^2(\xi^2 - 1)g' = \xi^2g^2 - 2\xi(\xi^2 - \alpha)g + \alpha(\alpha - 1)$.

Here $e_1 = u\partial_u$, $e_2 = \partial_x$, $e_3 = xu\partial_u$, $e_4 = \partial_t$, $e_5 = t\partial_t + x\partial_x$ and C_1, C_2 are arbitrary real constants.

Exact solutions of (7.4):

$$\langle e_4 \rangle : u = C_1t + C_2,$$

$$\langle e_2 \rangle : u = (x + C_1)[1 - \ln|x + C_1|] + C_2,$$

$$\langle e_2 + \alpha e_4 \rangle (\alpha > 0, \alpha \neq 1) : u = (x - \alpha t) \ln|1 - \alpha^2| - (x - \alpha t + C_1)[\ln|x - \alpha t + C_1| - 1] + C_2,$$

$$\langle e_3 + \alpha e_4 \rangle (\alpha > 0) : u = \alpha^{-1}tx + \alpha^2 \exp(\alpha^{-1}t) + C_1t + C_2,$$

$$\langle e_2 + \epsilon e_3 \rangle (\epsilon = \pm 1) : u = \frac{1}{2}\epsilon t^2 + \varphi(x), \quad \varphi' = y(x),$$

$$y - \ln|\epsilon - e^y| = \epsilon x + C_1,$$

$$\langle e_2 + \beta e_3 + e_4 \rangle (\beta \neq 0) : u = \frac{1}{2}\beta t^2 + (x - t) \ln|\beta| + C_1,$$

$$\langle e_2 + \beta e_3 + \alpha e_4 \rangle (\beta \neq 0, \alpha > 0, \alpha \neq 1) : u = \frac{1}{2}\beta t^2 + \varphi(\eta), \quad \eta = x - \alpha t,$$

$$\varphi = y'(\eta), \quad y - \ln|\beta - e^y| = \beta(1 - \alpha^2)^{-1}\eta + C_1,$$

$$\langle e_5 \rangle : u = -x \ln|x| + C_1t + x,$$

$$\langle e_5 + \alpha e_3 \rangle (\alpha > 0) : u = (\alpha t - x) \ln|x| + x\varphi(\xi), \quad \xi = tx^{-1},$$

$$(\xi^2 - 1)\varphi'' + \exp[-\xi\varphi' + \varphi + \alpha\xi - 1] = 1 + \alpha\xi.$$

Here $e_1 = \partial_u$, $e_2 = \partial_t$, $e_3 = t\partial_u$, $e_4 = \partial_x$, $e_5 = t\partial_t + x\partial_x + (u-x)\partial_u$ and C_1, C_2 are arbitrary real constants.

Exact solutions of (7.5):

$$\langle e_1 + e_4 \rangle : u = x + C_1 t + C_2,$$

$$\langle e_4 \rangle : u = C_1 t + C_2,$$

$$\langle e_2 \rangle : u = \varphi(x), \quad \varphi' = f(x), \quad \int \frac{df}{\ln|f|} = -mx + C_1,$$

$$\langle e_2 + e_4 \rangle : u = x - t + C_1,$$

$$\langle e_2 + \alpha e_4 \rangle (\alpha > 0, \alpha \neq 1) : u = \varphi(\eta), \quad \eta = x - \alpha t, \quad \varphi' = f(\eta),$$

$$\int \frac{df}{\ln|f|} = -\frac{m}{1-\alpha^2}(x - \alpha t) + C_1,$$

$$\langle e_3 + \alpha e_4 \rangle (\alpha > 0) : u = \alpha^{-1}tx + \frac{1}{2}mt^2(1 + \ln \alpha) + \frac{1}{2}mt^2 \left(\ln|t| - \frac{1}{2} \right) + C_1 t + C_2, \quad \alpha > 0,$$

$$\langle e_5 \rangle : u = \frac{1}{2}m^2 \ln|t| + t^2\varphi, \quad \text{where } \varphi = \varphi(\xi) \quad (\xi = tx^{-1}) \text{ satisfies equation}$$

$$\xi^2(\xi^2 - 1)\varphi'' + 2\xi(\xi^2 - 2)\varphi' - 2\varphi + m \ln|\xi^2\varphi'| - \frac{3}{2}m = 0.$$

Here $e_1 = \partial_u$, $e_2 = \partial_t$, $e_3 = t\partial_u$, $e_4 = \partial_x$, $e_5 = t\partial_t + x\partial_x + (2u + \frac{1}{2}mt^2)\partial_u$ and C_1, C_2 are arbitrary real constants.

Exact solutions of (7.6):

$$\langle e_1 + e_4 \rangle : u = \frac{1}{2}t^2 + x + C_1 t + C_2,$$

$$\langle e_4 \rangle : u = C_1 t + C_2,$$

$$\langle e_2 \rangle : u = (2-k)^{-1}(C_1 + (k-1)x)^{\frac{2-k}{1-k}} + C_2 \text{ and } k \neq 2, \\ u = C_2 \ln|x - C_1| \text{ and } k = 2;$$

$$\langle e_2 + \alpha e_4 \rangle (\alpha > 0, \alpha \neq 1) : u = (1 - \alpha^2) \ln|C_1 - x + \alpha t| + C_2 \text{ and } k = 2,$$

$$u = \frac{1-k}{2-k} \left(\frac{1-k}{\alpha^2 - 1} \right)^{\frac{1}{1-k}} |x - \alpha t + C_1|^{\frac{2-k}{1-k}} + C_2 \text{ and } k \neq 0, 1, 2,$$

$$\langle e_3 + e_4 \rangle : u = tx + t(\ln|t| - 1) + C_1 t + C_2 \text{ and } k = -1,$$

$$u = tx - \ln|t| + C_1 t + C_2 \text{ and } k = -2;$$

$$u = tx + (k^2 + 3k + 2)^{-1}|t|^{k+2} + C_1 t + C_2 \text{ and } k \neq 0, 1, -1, -2,$$

$(e_2 + \epsilon e_3)$ ($\epsilon = \pm 1$) : $u = \frac{1}{2}\epsilon t^2 \int^x f(\eta)d\eta + C_2$, where f is defined by

$$\int \frac{df}{\epsilon - |f|^k} = \eta + C_1,$$

$(e_2 + \epsilon e_3 + \alpha e_4)$ ($\epsilon = \pm 1, \alpha > 0$) : $u = \frac{1}{2}\epsilon t^2 \int^\eta f(z)dz + C_2$,

$$\eta = x - \alpha t, \text{ where } f \text{ is defined by } \int \frac{df}{\epsilon - |f|^k} = (1 - \alpha^2)^{-1}z + C_1,$$

(e_5) : $u = |t|^{\frac{k-2}{k-1}}\varphi(\xi), \xi = tx^{-1}$, where φ is defined by

$$\xi^2(\xi^2 - 1)\varphi'' + 2\xi\left(\xi^2 - \frac{k-2}{k-1}\right)\varphi' + (-1)^k \xi^{2k}|\varphi'|^k + \frac{k-2}{(k-1)^2}\varphi = 0$$

and $k \neq 0, 1, 2$,

$$u = \int^\xi \left[C_1(1 - \eta^2) + \frac{1}{4}(1 - \eta^2) \ln \left| \frac{1 + \eta}{1 - \eta} \right| - \frac{1}{2}\eta \right]^{-1} d\eta + C_2,$$

$\xi = tx^{-1}$ and $k = 2$,

$(e_5 + \alpha e_1)$ ($\alpha > 0$) : $u = \alpha \ln |t| + \int^\xi f(\eta)d\eta + C_1$,

$\xi = tx^{-1}$, $f = f(\eta)$ is a solution of Riccati equation

$$\eta^2(\eta^2 - 1)f' + \eta^4 f^2 + 2\eta^3 f + \alpha = 0 \text{ and } k = 2.$$

Here $e_1 = \partial_u, e_2 = \partial_t, e_3 = t\partial_u, e_4 = \partial_x, e_5 = t\partial_t + x\partial_x + \frac{k-2}{k-1}u\partial_u, k \neq 0, 1$ and C_1, C_2 are arbitrary real constants.

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