

**ON LOCAL EQUIVALENCE
FOR VECTOR FIELD SYSTEMS**

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IMA Preprint Series # 543

June 1989

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Abstract. We give sufficient conditions for C^∞ vector field systems on \mathbb{R}^n with genus $g = 1$ to be diffeomorphic to a contact structure. The diffeomorphism is explicitly constructed and used to give the most general integral submanifolds for the systems. Finally the implications of these results for integrable hyperbolic partial differential equations in the plane is discussed.

The aim of this note is to give some new results on the explicit integration of certain Pfaffian systems with genus $g = 1$. The approach adopted is as follows. Given a Pfaffian system E on \mathbb{R}^n , we ask the question; when is E equivalent to a *contact structure* with respect to the diffeomorphisms of \mathbb{R}^n and how can an explicit diffeomorphism be constructed? The construction of such a diffeomorphism provides the explicit integral submanifolds for E .

The subject of the integration of Pfaffian systems has a long history dating back to 1815 when G.F. Pfaff first posed it [1]. Since that time important contributions have been made by a number of authors including von Weber, Goursat, Frobenius, Natani, Clebsch, Grassman, Darboux, and others in the 19th century; (see E. Goursat [8]). In the early 20th century, Cartan's works [2] are well known; less well known is the paper [3] of E. Vessiot, which contains a number of important results including applications to automorphic systems (see also Pommaret [12] and Kumpera [13]). More recent works are those of Gardner [4] and Bryant [5]. We also note recent applications of the theory of Pfaffian systems to control theory by Hermann [7], [11].

The motivation for solving the problem here posed comes from the theory of integrable wave equations. In his study of scalar nonlinear second order partial differential equations in two independent variables, Vessiot [6] showed that whenever the second order characteristics, \mathcal{M} , possess two or more invariants and these invariants are taken as independent coordinates, then \mathcal{M} is equivalent to a contact structure on a reduced manifold. This fact allowed Vessiot to settle the integration problem for these equations in an elegant way. These matters will be discussed briefly in section 4.

§2. A. The Theory of Pfaffian Systems. In this section I review briefly some of the known results on Pfaffian systems. Let $T\mathbb{R}^n$ and $T^*\mathbb{R}^n$ denote the tangent and cotangent bundles on \mathbb{R}^n respectively. Let $C^\infty(\mathbb{R}^n)$ denote the ring of real valued, C^∞ functions on \mathbb{R}^n . Let $\Gamma^*(T^*(\mathbb{R}^n))$ denote the $C^\infty(\mathbb{R}^n)$ -module of smooth sections of $T^*(\mathbb{R}^n)$. A *Pfaffian system on \mathbb{R}^n* is a $C^\infty(\mathbb{R}^n)$ -submodule of $\Gamma^*(T^*(\mathbb{R}^n))$. Let $\Gamma(T(\mathbb{R}^n))$ denote the $C^\infty(\mathbb{R}^n)$ -module of smooth sections of $T(\mathbb{R}^n)$. A *vector field system on \mathbb{R}^n* is a $C^\infty(\mathbb{R}^n)$ -submodule

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of $\Gamma(T(\mathbb{R}^n))$. A Pfaffian system E (vector field system \tilde{E}) is said to have dimension r at point $x \in \mathbb{R}^n$ if there exists a neighborhood $U \ni x$ and r -linearly independent 1-forms (vector fields) defined on U which generate $E(\tilde{E})$. I will assume throughout that vector fields systems and Pfaffian systems have constant dimension as x ranges over U . Such vector field systems will sometimes be referred to as *distributions*.

The first important object we may associate to E is a vector field system called the *characteristic system* of E or the *Cauchy system* of E denoted $char E$:

$$char E := \{X \in \Gamma(T\mathbb{R}^n) : X \rfloor w = 0, X \rfloor dw \in E, \text{ for all } w \in E\},$$

where \rfloor denotes the interior product. The annihilator of $char E$ is called to *Cartan system* of E denoted $C(E)$. That is

$$C(E) = (char E)^\perp.$$

The dimension of the Cartan system of E is an invariant called the *class* of E . The class of E is the minimum number of variables necessary in order to write down local generators for the system (see Gardner [4] and Cartan [2]).

The next important construction is that of the *derived map* of E which is the $C^\infty(\mathbb{R}^n)$ -linear map

$$\pi \circ d = \delta_E : E \longrightarrow \Gamma(\wedge^2 T^*(\mathbb{R}^n)) / \Gamma(T^*(\mathbb{R}^n)) \wedge E$$

obtained by exterior differentiation and then projection into the quotient structure.

Suppose δ_E has constant rank on \mathbb{R}^n . Define the *first derived system* $E^{(1)}$ of E by $E^{(1)} = \ker \delta_E$. Inductively, define the *derived map* $\delta_{E^{(i)}}$ of E

$$\delta_{E^{(i)}}; E^{(i)} \longrightarrow \Gamma(\wedge^2 T^*(\mathbb{R}^n)) / \Gamma(T^*(\mathbb{R}^n)) \wedge E^{(i)}.$$

Then the $(i + 1)$ -derived system $E^{(i+1)}$ of E is given by

$$E^{(i+1)} = \ker \delta_{E^{(i)}}.$$

Letting $g_i = \dim E^{(i)}$ we have $g_0 \geq g_i \geq \dots$

There must be an N so that $g_N = g_{N+1}$. N defined in this way is called the *derived length* of E . So if $N = 0$, then $E^{(1)} = E$ so $\delta_E = 0$. This means $d(E) \subset E \wedge \Gamma(T^*\mathbb{R}^n)$, that is, E is Frobenius-integrable. Thus δ is the zero map when E is integrable so we have the intuitive idea that δ somehow measures the amount of non-integrability of E . The collection $(E, E^{(1)}, \dots, E^{(N)})$ is called the *derived flag* of E . We have

$$E^{(N)} \subset E^{(N-1)} \subset \dots \subset E.$$

Define integers

$$p_0 = \dim E^{(N)}; \quad p_{N-i} = \dim(E^{(i)} / E^{(i+1)}), \\ 0 \leq i \leq N - 1.$$

We have that

$$\dim E = \sum_{i=0}^N p_i.$$

If p_{N+1} denotes the codimension of E in $C(E)$ then

$$\text{class } E = \sum_{i=0}^{N+1} p_i.$$

The Pfaffian system E is said to be of type (Gardner [4])

$$(p_0, p_1, \dots, p_N, p_{N+1}).$$

Finally, we mention another important invariant, namely the *rank* or *Engle rank*. E has rank ρ if ρ is the smallest integer such that

$$\alpha \wedge (d\alpha)^\rho = 0,$$

for all $\alpha \in E$.

The following theorem is proved in Bryant, Chern & Griffiths [14].

THEOREM 2.1. *In a neighborhood of \mathbb{R}^n suppose α is a 1-form of rank ρ . Then there exists a coordinate system w_1, w_2, \dots, w_n , possibly in a smaller neighborhood such that the form α becomes*

$$dw_1 + w_2 dw_3 + w_4 dw_5 + \dots + w^{2\rho} dw^{2\rho+1}.$$

Interestingly, a proof of this classical theorem is given in [3] using vector fields. The following is also classical and due to von Weber (see [8]).

THEOREM 2.2. *Let E be a Pfaffian system of type*

$$(p_0, 1, 1, 1, \dots, 1, 2)$$

and suppose

$$\dim(C(E^{(i)})/E^{(i)}) = 2, \quad 1 \leq i \leq N - 1,$$

and $E^{(N)}$ is a p_0 -dimensional completely integrable system, then there exist $N + p_0 + 2$ independent functions

$$\{x, z_1, \dots, z_{p_0}, y, y_1, \dots, y_N\}$$

such that E is locally generated by

$$\{dz_1, dz_2, \dots, dz_{p_0}, dy - y_1 dx, dy_1 - y_2 dx, \dots, dy_{N-1} - y_N dx\}.$$

For a proof of this theorem the reader is referred to [5]. Theorems 2.1 and 2.2 were the only local structure theorems known (apart from the case of Frobenius integrable systems) until Bryant gave a new structure theorem [5, Theorem 2.1].

B. Vector Field Systems. It will be convenient for what is to follow in §§3&4 to give some results and definitions for vector field systems. Just as we stated the notion of characteristic vector fields for E we can do so for the dual \tilde{E} . The characteristic system for \tilde{E} , denoted $char\tilde{E}$ is the set of all $\delta \in \tilde{E}$ such that $[\delta, \tilde{E}] \subseteq \tilde{E}$.

THEOREM 2.3. *If E and \tilde{E} are dual then $charE = char\tilde{E}$.*

Proof. Contraction of forms and fields is defined as follows:

$$(X]w)(X_1, \dots, X_{p-1}) = w(X, X_1, \dots, X_{p-1})$$

where X, X_1, \dots, X_{p-1} are vector fields and w is a p -form. We have the identity

$$d\theta(v_1, v_2) = v_1(v_2]\theta) - v_2(v_1]\theta) - \theta]([v_1, v_2]),$$

for 1-forms θ and vector fields v_1, v_2 . So if $\delta \in char\tilde{E}$ and $X \in \tilde{E}$, we have

$$d\theta(\delta, X) = \delta(X]\theta) - X(\delta]\theta) - \theta]([\delta, X])$$

and for all $\theta \in E$

$$d\theta(\delta, X) = (\delta]\theta)(X) = \theta]([\delta, X]) = 0. \quad (*)$$

So $char\tilde{E} \subseteq charE$. Assuming $\theta \in charE$ and using (*) then proves the theorem.

Once again, as for Pfaffian systems, we can define the concept of a derived system for a vector field system. The i^{th} derived system $\tilde{E}^{(i)}$ of \tilde{E} is defined recursively

$$\tilde{E}^{(i)} = \tilde{E}^{(i-1)} + [\tilde{E}^{(i-1)}, \tilde{E}^{(i-1)}], \tilde{E}^{(0)} = \tilde{E}.$$

There will always be an integer N such that

$$\tilde{E}^{(N)} = \tilde{E}^{(N+1)}$$

meaning that $\tilde{E}^{(N)}$ is Frobenius integrable.

THEOREM 2.4. *For each $i = 0, 1, 2, \dots, N$.*

$$\dim E^{(i)} + \dim \tilde{E}^{(i)} = n.$$

Proof. It will be enough to prove that $E^{(i)}$ and $\tilde{E}^{(i)}$ are dual. Let $x, y \in \tilde{E}$. Define a set of 1-forms

$$\mathfrak{F} = \{\theta \in \Gamma(T^*\mathbb{R}^n) \mid \theta]x = \theta]y = \theta][x, y] = 0\}.$$

it follows that $\mathfrak{F} \subseteq E$ and that $d\theta(x, y) = 0$. Clearly if \tilde{E} is Frobenius integrable then $\mathfrak{F} = E$. Otherwise we have the maximal subset of E such that $d\theta \in \Gamma(T^*\mathbb{R}^n) \wedge E$. That is $\theta \in \ker \delta_E = E^{(1)}$. This proves that $E^{(1)}$ is the annihilator of $\tilde{E}^{(1)}$ and the assertion for $i = 1$. To prove the result for $i > 1$, repeat the argument choosing $x, y \in E^{(i-1)}$ to show that $(\tilde{E}^{(i)})^\perp = E^{(i)}$. We need only do this a finite number of times and the theorem is proved.

Remark We now have because of theorems 2.3 and 2.4

$$\begin{aligned} p_0 &= n - \dim \tilde{E}^{(N)} \\ p_{N-i} &= \dim(\tilde{E}^{(i+1)}/\tilde{E}^{(i)}), \quad 0 \leq i \leq N-1, \\ \dim(\tilde{E}^{(i)}/\text{char} \tilde{E}^{(i)}) &= \dim(C(E^{(i)})/E^{(i)}), \quad 0 \leq i \leq N-1. \end{aligned}$$

§ 3. Vector Field Systems, Contact Structures and Integral Submanifolds. In this section we obtain some results which concern generalizations of the von Weber systems (theorem 2.2) in terms of the structure of the associated vector field system. Specifically we consider vector field systems \tilde{E} locally generated in $\mathbb{R}^{(k+1)q+1}$ by

$$\left\{ X, \frac{\partial}{\partial \xi_1}, \frac{\partial}{\partial \xi_2}, \dots, \frac{\partial}{\partial \xi_q} \right\}$$

and such that the dual Pfaffian system has type

$$(0, q, q, \dots, q+1),$$

while

$$\dim(C(E^{(i)})/E^{(i)}) = q+1, \quad 1 \leq i \leq N-1$$

and derived length $N = k$.

We will obtain sufficient conditions in terms of the structure of \tilde{E} in order that \tilde{E} be diffeomorphic to the k^{th} -order contact structure $\Omega^k(\mathbb{R}, \mathbb{R}^q)^\perp$ on $J^k(\mathbb{R}, \mathbb{R}^q)$, the k^{th} order jet bundle of maps $\mathbb{R} \rightarrow \mathbb{R}^q$. Importantly, we will show how the explicit diffeomorphism may be constructed. Notice that when $q = 1$, we obtain the von Weber system which is known to be equivalent to the k^{th} order contact structure $\Omega^k(\mathbb{R}, \mathbb{R})$ on $J^k(\mathbb{R}, \mathbb{R})$.

On $J^k(\mathbb{R}, \mathbb{R}^q)$, the k^{th} -order contact structure is locally generated by

$$\begin{aligned} & \left\{ \frac{\partial}{\partial x} + z_{11} \frac{\partial}{\partial z_{10}} + z_{12} \frac{\partial}{\partial z_{11}} + \dots + z_{1k} \frac{\partial}{\partial z_{1,k-1}} + z_{21} \frac{\partial}{\partial z_{20}} + \dots + z_{2k} \frac{\partial}{\partial z_{2,k-1}} + \dots + \right. \\ & \quad \left. + z_{q1} \frac{\partial}{\partial z_{q0}} + \dots + z_{qk} \frac{\partial}{\partial z_{q,k-1}}, \frac{\partial}{\partial z_{1k}}, \frac{\partial}{\partial z_{2k}}, \dots, \frac{\partial}{\partial z_{qk}} \right\} \\ & = \left\{ J_k^q, \frac{\partial}{\partial z_{1k}}, \dots, \frac{\partial}{\partial z_{qk}} \right\}. \end{aligned}$$

We seek a diffeomorphism $\psi : \mathbb{R}^{(k+1)q+1} \longrightarrow J^k(\mathbb{R}, \mathbb{R}^q)$ such that

$$(3.1) \quad \psi_* \tilde{E} = \Omega^k(\mathbb{R}, \mathbb{R}^q)^\perp$$

where \tilde{E} is of the stated type. Equation (3.1) means that there exists $g \in GL(q+1, \mathbb{R})$ such that

$$(3.2) \quad \psi_* \begin{pmatrix} X \\ \partial/\partial\xi_1 \\ \vdots \\ \partial/\partial\xi_q \end{pmatrix} = g \begin{pmatrix} J_k^q \\ \partial/\partial z_{1k} \\ \vdots \\ \partial/\partial z_{qk} \end{pmatrix}.$$

Suppose such a ψ exists then by the chain rule

$$(3.3) \quad \begin{aligned} \psi_*(X) &= (Xx) \frac{\partial}{\partial x} + \sum_{j=1}^q \sum_{i=0}^k (Xz_{ji}) \frac{\partial}{\partial z_{ji}}, \\ \psi_*(\partial/\partial\xi_\mu) &= \frac{\partial x}{\partial\xi_\mu} \frac{\partial}{\partial x} + \sum_{j=1}^q \sum_{i=0}^k \left(\frac{\partial z_{ji}}{\partial\xi_\mu} \right) \frac{\partial}{\partial z_{ji}}. \end{aligned}$$

If g has the form

$$g = (\gamma_{\alpha\beta})_{1 \leq \alpha, \beta \leq q+1}$$

equations (3.2), making use of (3.3) become

$$(3.4) \quad \left\{ \begin{array}{l} \left\{ \begin{array}{l} Xx = \gamma_{11}, \\ Xz_{ji} = \gamma_{11} z_{j,i+1}, \quad 1 \leq j \leq q, 0 \leq i \leq k-1, \end{array} \right. \\ \\ \left\{ \begin{array}{l} \frac{\partial x}{\partial\xi_\mu} = \gamma_{\mu+1,1}, \quad 1 \leq \mu \leq q \\ \frac{\partial z_{ji}}{\partial\xi_\mu} = \gamma_{\mu+1,1} z_{j,i+1}, \quad 1 \leq \mu \leq q, 1 \leq j \leq q, 0 \leq i \leq k-1, \end{array} \right. \\ \\ \left\{ \begin{array}{l} \frac{\partial z_{jk}}{\partial\xi_\mu} = \gamma_{\mu+1,j+1}, \quad 1 \leq \mu \leq q, 1 \leq j \leq q, \\ Xz_{jk} = \gamma_{1,j+1}, \quad 1 \leq j \leq q. \end{array} \right. \end{array} \right.$$

In general the system is overdetermined, the difference between the number of equations and unknowns being

$$kq^2 - q - 1.$$

Hence in general no diffeomorphism ψ exists. Furthermore, as the equations stand it is not clear how they can be solved in terms of the stated structure conditions. In fact, we restrict ourselves to Pfaffian systems of the stated type only to insure that the *necessary conditions* for equivalence are satisfied. In particular, see Hermann [7, p 138-141] for

a discussion of this point. However, considerable simplification is effected if we restrict ourselves to diffeomorphisms such that g lies in a subgroup of $GL(q+1, \mathbf{R})$, namely the subgroup generated by the matrices $g \in GL(q+1, \mathbf{R})$ of the form

$$g = \begin{pmatrix} \gamma_{11} & \gamma_{12} \cdots & \gamma_{1,q+1} \\ 0 & \gamma_{22} \cdots & \gamma_{2,q+1} \\ 0 & \gamma_{32} \cdots & \gamma_{3,q+1} \\ \vdots & \vdots & \cdots & \vdots \\ 0 & \gamma_{q+1,2} \cdots & \gamma_{q+1,q+1} \end{pmatrix}.$$

The equations (3.4) then became

$$(3.5) \quad \begin{aligned} (A) & \begin{cases} Xx = \gamma_{11} \\ Xz_{ji} = \gamma_{11}z_{j,i+1} \end{cases} & 1 \leq j \leq q, 0 \leq i \leq k-1 \\ (B) & \begin{cases} \frac{\partial x}{\partial \xi_\mu} = 0 \\ \frac{\partial z_{ji}}{\partial \xi_\mu} = 0 \end{cases} & 1 \leq \mu \leq q, 1 \leq j \leq q, 0 \leq i \leq k-1 \\ (C) & \frac{\partial z_{jk}}{\partial \xi_\mu} = \gamma_{\mu+1,j+1}, Xz_{jk} = \gamma_{i,j+1}, & 1 \leq \mu \leq q, 1 \leq j \leq q. \end{aligned}$$

which would seem to make matters worse since the system (3.5) is even more over determined than (3.4). In fact the later system is manageable. Indeed, it is not difficult to see that the only integrability conditions for the existence of a solution are

$$(3.6)_1 \quad 0 = \frac{\partial z_{j1}}{\partial \xi_\mu} = \frac{1}{\gamma_{11}} \left\{ \left[\frac{\partial}{\partial \xi_\mu}, X \right] z_{j0} - \frac{\partial \gamma_{11}}{\partial \xi_\mu} z_{j1} \right\}$$

$$(3.6)_2 \quad 0 = \frac{\partial z_{j2}}{\partial \xi_\mu} = \frac{1}{\gamma_{11}} \left\{ \left[\frac{\partial}{\partial \xi_\mu}, X \right] z_{j1} - \frac{\partial \gamma_{11}}{\partial \xi_\mu} z_{j2} \right\}$$

⋮

$$(3.6)_{k-1} \quad 0 = \frac{\partial z_{j,k-1}}{\partial \xi_\mu} = \frac{1}{\gamma_{11}} \left\{ \left[\frac{\partial}{\partial \xi_\mu}, X \right] z_{j,k-1} - \frac{\partial \gamma_{11}}{\partial \xi_\mu} z_{jk} \right\}$$

where

$$1 \leq j \leq q, 1 \leq \mu \leq q$$

and where by assumption on g , $\gamma_{11} \neq 0$.

Now from the first equation in 3.5(A), we have

$$\frac{\partial}{\partial \xi_\mu} Xx = \frac{\partial}{\partial \xi_\mu} Xx - X \frac{\partial x}{\partial \xi_\mu} = \left[\frac{\partial}{\partial \xi_\mu}, X \right] x = \frac{\partial \gamma_{11}}{\partial \xi_\mu}.$$

Hence, assuming, x to be an invariant of the distribution

$$(3.7) \quad \left\{ \frac{\partial}{\partial \xi_\mu}, \left[\frac{\partial}{\partial \xi_\mu}, X \right] \right\}_{\mu=1}^k$$

equation (3.6)₁ becomes

$$\left[\frac{\partial}{\partial \xi_\mu}, X \right] z_{j0} = 0, \quad 1 \leq j \leq q.$$

We have therefore shown that the compatibility condition (3.6)₁ will be satisfied if x and z_{j0} are chosen to be invariants of the subdistribution (3.7) of $\tilde{E}^{(1)}$.

Next, consider the compatibility condition (3.6)₂. Assume that x and z_{j0} are chosen as above. (3.6)₂ then becomes

$$(3.8) \quad \left[\frac{\partial}{\partial \xi_\mu}, X \right] z_{j1} = 0.$$

But, $X z_{j0} = \gamma_{11} z_{j1}$, hence (3.8) becomes

$$(3.9) \quad \left[\frac{\partial}{\partial \xi_\mu}, X \right] \left(\frac{1}{\gamma_{11}} X z_{j0} \right) = 0.$$

Now,

$$\left[\frac{\partial}{\partial \xi_\mu}, X \right] \left(\frac{1}{\gamma_{11}} \right) = -\frac{1}{(\gamma_{11})^2} \left[\frac{\partial}{\partial \xi_\mu}, X \right] \gamma_{11} = -\frac{1}{(\gamma_{11})^2} \left[\frac{\partial}{\partial \xi_\mu}, X \right] Xx.$$

But since $\left[\frac{\partial}{\partial \xi_\mu}, X \right] x = 0$, we have

$$\left[\frac{\partial}{\partial \xi_\mu}, X \right] \left(\frac{1}{\gamma_{11}} \right) = \frac{-1}{(\gamma_{11})^2} \left[\left[\frac{\partial}{\partial \xi_\mu}, X \right], X \right] x.$$

Hence

$$(3.10) \quad \begin{aligned} \left[\frac{\partial}{\partial \xi_\mu}, X \right] z_{j1} &= -\frac{(X z_{j0})}{(\gamma_{11})^2} \left[\left[\frac{\partial}{\partial \xi_\mu}, X \right], X \right] x \\ &\quad + \frac{1}{\gamma_{11}} \left[\left[\frac{\partial}{\partial \xi_\mu}, X \right], X \right] z_{j0}. \end{aligned}$$

From (3.10), it follows that compatibility conditions (3.6)₁ and (3.6)₂ will be satisfied if x and z_{j0} are chosen to be any invariants of the subdistribution

$$\left\{ \frac{\partial}{\partial \xi_\mu}, \left[\frac{\partial}{\partial \xi_\mu}, X \right], \left[\left[\frac{\partial}{\partial \xi_\mu}, X \right], X \right] \right\}_{\mu=1}^q$$

of $\tilde{E}^{(2)}$.

We may proceed to consider the compatibility condition (3.6)₃. By a similar argument and a tedious but straight-forward computation, we find that (3.6)₁, (3.6)₂ and (3.6)₃ will be satisfied if x and z_{j_0} are chosen to be any invariants of the subdistribution

$$\left\{ \frac{\partial}{\partial \xi_\mu}, \left[\frac{\partial}{\partial \xi_\mu}, X \right], \left[\left[\frac{\partial}{\partial \xi_\mu}, X \right], X \right], \left[\left[\left[\frac{\partial}{\partial \xi_\mu}, X \right], X \right], X \right] \right\}_{\mu=1}^q$$

of $\tilde{E}^{(3)}$. Proceeding in this way we finally arrive at the last compatibility condition (3.6)_{k-1} and we have the following theorem.

THEOREM 3.1. *Let $\tilde{E} : \left\{ X, \frac{\partial}{\partial \xi_1}, \frac{\partial}{\partial \xi_2} \dots \frac{\partial}{\partial \xi_q} \right\}$ be a C^∞ -vector field system on $\mathbb{R}^{(k+1)q+1}$, of type*

$$(0, q, q, q, \dots, q, q+1).$$

and derived length k and such that

$$\dim(\tilde{E}^{(i)} / \text{char } \tilde{E}^{(i)}) = q+1, \quad 1 \leq i \leq k-1.$$

Suppose the distribution $\Gamma \subseteq \tilde{E}^{(k-1)}$ generated by

$$\Gamma : \left\{ \frac{\partial}{\partial \xi_\mu}, \left[\frac{\partial}{\partial \xi_\mu}, X \right], \left[\frac{\partial}{\partial \xi_\mu}, X \right], X, \dots, \left[\left[\dots \left[\frac{\partial}{\partial \xi_\mu}, X \right], X \right], \dots, X \right] \right\}_{\mu=1}^{(k-1)\text{-times}}$$

possesses $q+1$ functionally independent invariants x and $\{z_{j_0}\}_{j=1}^q$ and define the set of functions $z_{j_1}, z_{j_2}, \dots, z_{j_k}, 1 \leq j \leq q$, by the equations

$$z_{j_1} = \frac{X z_{j_0}}{X x}, \quad z_{j_2} = \frac{X z_{j_1}}{X x}, \dots, \quad z_{j_k} = \frac{X z_{j, k-1}}{X x}.$$

Then, whenever the set of functions $\{x, z_{j_0}, \dots, z_{j_k}\}$ so defined are functionally independent, they define a diffeomorphism

$$\psi : \mathbb{R}^{(k+1)q+1} \longrightarrow J^k(\mathbb{R}, \mathbb{R}^q).$$

such that

$$\psi_* \tilde{E} = \Omega^k(\mathbb{R}, \mathbb{R}^q)^\perp.$$

The means of construction of the integral submanifolds for the Pfaffian systems of the type considered in this paper is now clear. By application of Cartan's theory or otherwise its not difficult to show that the genus of \tilde{E} is *one*. An integral submanifold is therefore a map

$$\Phi : I \longrightarrow \mathbb{R}^N, \quad I \subset \mathbb{R},$$

such that

$$\Phi^* E = 0.$$

Assuming the map ψ has been found as per theorem 3.1 we have

$$(\psi^{-1})^* E = \Omega^k(\mathbf{R}, \mathbf{R}^q).$$

Hence if $j^k f : x \rightarrow J^k(\mathbf{R}, \mathbf{R}^q)$ is the k -jet extension of any C^k map, $f : \mathbf{R} \rightarrow \mathbf{R}^q$ then

$$(j^k f)^*(\psi^{-1})^* E = (j^k f)^* \Omega^k(\mathbf{R}, \mathbf{R}^q) = 0.$$

Hence

$$(3.11) \quad \Phi = \psi^{-1} \circ j^k f : \mathbf{R} \rightarrow \mathbf{R}^N$$

is the sought after integral submanifold for E in terms of the (now known) diffeomorphism ψ .

Remark In view of the works [7,11] it may be worth mentioning the connection between the result of this paper and control theory. Note that in this section we have been studying Pfaffian systems of the form

$$dx_i - f_i(x_j, u_j) dt$$

where the x_i may be regarded as state variables and the u_j as control variables. The algorithm here developed expresses each of the variables $x_i(t), u_j(t)$ in terms of arbitrary functions and their derivatives thereby “parametrising” the state and control space. This may be compared with the so called “Problem of Monge” (see [8, p310], [11, p353]).

Example We give a simple example to illustrate the results. On \mathbf{R}^7 we have the vector field system

$$\tilde{E} : \left\{ \frac{\partial}{\partial t} + \xi_1 \left(x_1 \frac{\partial}{\partial x_1} + t x_2 \frac{\partial}{\partial x_2} \right) + \xi_2 \left(x_3 \frac{\partial}{\partial x_3} + t x_4 \frac{\partial}{\partial x_4} \right), \frac{\partial}{\partial \xi_1}, \frac{\partial}{\partial \xi_2} \right\}.$$

\tilde{E} has type $(0, 2, 2, 3)$ and satisfies $\dim(E^{(i)})/\text{char } \tilde{E}^{(i)} = 3, i = 1, 2$ and hence satisfies the necessary conditions for equivalence to the contact structure $\Omega^2(\mathbf{R}, \mathbf{R}^2)^\perp$. Calculating therefore the subdistribution Γ in theorem 3.1, that is,

$$\left\{ \frac{\partial}{\partial \xi_\mu}, \left[\frac{\partial}{\partial \xi_\mu}, X \right] \right\}_{\mu=1}^2$$

we easily find its three invariants in the form

$$\begin{aligned} x &= t, \\ z_{10} &= x_1^t / x_2, \\ z_{20} &= x_3^t / x_4. \end{aligned}$$

From the equations of the theorem 3.1, we obtain the remaining variables

$$\begin{aligned} z_{11} &= \frac{x_1^t}{x_2} \ln x_1, & z_{21} &= \frac{x_3^t}{x_4} \ln x_3, \\ z_{12} &= \frac{x_1^t}{x_2} (\ln x_1)^2 + \xi_1 \frac{x_1^t}{x_2}, & z_{22} &= \frac{x_3^t}{x_4} (\ln x_3)^2 + \xi_2 \frac{x_3^t}{x_4} \end{aligned}$$

and the map ψ is explicitly found. It's easy to check that ψ is a local diffeomorphism and that $\psi_* \tilde{E} = \Omega^2(\mathbf{R}, \mathbf{R}^2)^\perp$. Calculating ψ^{-1} , we find the integral submanifold for E is given explicitly by equation (3.11):

$$\begin{aligned} \Phi : t &\longrightarrow \left(t, \exp(f_1'/f_1), 1/f_1 \exp\left(\frac{t f_1'}{f_1}\right), \exp(f_2'/f_2), 1/f_2 \exp(t f_2'/f_2) \right. \\ &\quad \left. \frac{f_1 f_1'' - f_1'^2}{(f_1)^2}, \frac{f_2 f_2'' - f_2'^2}{(f_2)^2} \right) \\ &= (t, x_1, x_2, x_3, x_4, \xi_1, \xi_2) \end{aligned}$$

where $f_1(t), f_2(t)$ are arbitrary C^2 functions $\mathbf{R} \rightarrow \mathbf{R}$. It is straight-forward to check that $\Phi^* E = 0$.

§ 4. Relationship to Integrable Wave Equations and the Work of Vessiot. As mentioned in the introduction, Vessiot [6] used the equivalence of vector field systems to contact structures to obtain explicit general solutions for integrable nonlinear wave equations. That is, equations of the form

$$(4.1) \quad \frac{\partial^2 u}{\partial x_1 \partial x_2} = f \left(x_1, x_2, u, \frac{\partial u}{\partial x_1}, \frac{\partial u}{\partial x_2} \right)$$

when there exist at least two independent first-integrals on each characteristic. In fact, Vessiot used the algorithm worked out in this paper in the case $k = 3$ and $q = 1$ but it seems, did not write out a proof. The details of this construction are as follows. We will consider the Liouville equation as an example but the reader will find many more in Vessiot's paper [6].

Take the Liouville equation in the form

$$(4.2) \quad \frac{\partial^2 u}{\partial x_1 \partial x_2} = e^u$$

There are two sets of characteristics for (4.2) on the submanifold $\varepsilon \subset J^2(\mathbf{R}^2, \mathbf{R})$ defined by

$$z_{12} - e^z = 0.$$

These are given by the one-dimensional solutions of the vector field systems

$${}_1\tilde{\Omega}_L^2 : \left\{ \frac{\partial}{\partial x_1} + z_1 \frac{\partial}{\partial z} + z_{11} \frac{\partial}{\partial z_1} + e^z \frac{\partial}{\partial z_2} + z_2 e^z \frac{\partial}{\partial z_{22}}, \frac{\partial}{\partial z_{11}} \right\}$$

$${}_2\tilde{\Omega}_L^2 : \left\{ \frac{\partial}{\partial x_2} + z_2 \frac{\partial}{\partial z} + e^z \frac{\partial}{\partial z_1} + z_{22} \frac{\partial}{\partial z_2} + z_1 e^z \frac{\partial}{\partial z_{11}}, \frac{\partial}{\partial z_{22}} \right\}$$

where $x_1, x_2, z, z_1, z_2, z_{11}, z_{12}, z_{22}$ are local coordinates on $J^2(\mathbf{R}^2, \mathbf{R})$ such that $(j^2 u)^* z_{11} = \frac{\partial^2 u}{\partial x_1 \partial x_1}$, etc. The distributions ${}_1\tilde{\Omega}_L^2$ and ${}_2\tilde{\Omega}_L^2$ are dual to the characteristic 1-forms listed for example by Goursat [9] in his discussion of Darboux' method for the general solution of the Liouville equation. One easily finds the invariants in the form

$${}_1\tilde{\Omega}_L^2 : \left\{ \begin{array}{l} \pi_1 = x_2, \\ \pi_2 = z_{22} - \frac{(z_2)^2}{2}, \end{array} \right. \quad {}_2\tilde{\Omega}_L^2 : \left\{ \begin{array}{l} \pi_1 = x_1, \\ \pi_2 = z_{11} - \frac{(z_1)^2}{2}. \end{array} \right.$$

Now if we adapt a new coordinate system incorporating the above invariants, say

$$\varphi : \varepsilon \longrightarrow \mathbf{R}^7$$

$$\begin{aligned} \mu_1 &= x_1 & \mu_3 &= z_{22} - (z_2)^2/2 \\ \mu_2 &= x_2 & \mu_4 &= z_{11} - (z_1)^2/2 \\ \zeta_0 &= z, & \zeta_1 &= z_1, & \zeta_2 &= z_2 \end{aligned}$$

then for example $\varphi_*({}_1\tilde{\Omega}_L^2)$ is locally generated by

$$\left\{ \frac{\partial}{\partial \mu_1} + \zeta_1 \frac{\partial}{\partial \zeta_0} + (\mu_4 + \zeta_1^2/2) \frac{\partial}{\partial \zeta_1} + e^{\zeta_0} \frac{\partial}{\partial \zeta_2}, \frac{\partial}{\partial \mu_4} \right\}$$

on the reduced manifold which is locally \mathbf{R}^5 . Applying theorem 3.1 to $\varphi_*({}_1\tilde{\Omega}_L^2)$ shows that it is diffeomorphic to the third order contact structure $\Omega^3(\mathbf{R}, \mathbf{R})^\perp$. This reduction however does not by itself give the general solution of (4.2). The reduction is only a partial integration. The extra information one requires is the existence of a *nonlinear map* linking the integral submanifolds of ${}_1\tilde{\Omega}_L^2$ and ${}_2\tilde{\Omega}_L^2$. In his remarkable study [6], Vessiot discovered such a nonlinear map for each scalar integrable equation (4.1) and thereby completely cleared up the integration problem for these equations. Finally, it is worth mentioning that recently Vessiot's results were extended to coupled systems of wave equations [10]. Interestingly, in this case the characteristics are *not* in general diffeomorphic to contact structures (see also [15]).

Acknowledgements. The research reported in this paper was carried out while I was a visitor at the Institute of Mathematics and its Applications, the University of Minnesota. It is a pleasure to thank the IMA for its warm hospitality and conducive atmosphere for research. I would also like to thank Peter Olver for a useful discussion of aspects of the results here presented and for pointing out relevant parts of the literature.

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