Solvable symmetry structures in differential form applications

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Abstract

We investigate symmetry te
hniques for expressing various exterior differential forms in terms of simplified coordinate systems. In particular, we give extensions of the Lie symmetry approa
h to integrating Frobenius integrable distributions based on a solvable stru
ture of symmetries and show how a solvable structure of symmetries may be used to find local coordinates for the Pfaffian problem and Darboux's theorem.

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Key words: Frobenius integrable, PfaÆan equations, Darboux's theorem.

Introduction $\mathbf 1$

In this paper we present several methods based on symmetry te
hniques for expressing various differential forms in simplified coordinate systems. We use work by Lie $[10]$ and Cartan $[3]$ to explore how symmetries may be used to integrate Frobenius integrable distributions. In re
ent times, Barasab-Horwath $[1]$, Duzhin and Lychagin $[6]$, Hartl and Athorne $[9]$, and Sherring and Prince [13] have extended Lie and Cartan's work from the perspective of constructing first integrals of a completely integrable distribution by quadratures. Our work uses such results to examine conditions under which a given differential form can be expressed in a simpler coordinate system.

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The plan of this paper is to first review Lie's solvable symmetry structure approach to integrating Frobenius integrable vector field distributions. For a given Frobenius integrable distribution an exterior produ
t of one-forms is used to generate invariants of the distribution. We then apply the approa
h to PfaÆan and Darboux systems, and in both ases present an example.

It is assumed throughout this paper that our expressions apply locally on some n -dimensional, open, simply connected differentiable submanifold, U . of \mathbb{R}^n , with coordinates $x^1, \ldots, x^n \in C^{\infty}(U)$. One further assumption that we make on U is that it be convex. This allows us to use the converse of the Poincare Lemma on the whole of U, i.e. if $\Delta \in \Lambda^2(U)$ is closed $(a\Delta = 0)$, $\lim_{\Omega} s_i = a \cup \text{ for some } \cup \in \Lambda \quad (U)$ [12, 15].

Consider the differentiable manifold U of dimension n. TU is the tangent bundle of vector fields with $T_p(U)$, $p \in U$ as its fibres. Let $\mathfrak{X}(U)$ be be the module of all smooth vector fields over $C^+(U)$. The state cotangent bundle with $T_p(U)$, $p \in U$ as its fibres. The set of exterior differential m-forms is a $\mathbf r$ section of the bundle of all homogeneous differential forms, $\Lambda(U)$. For any $\Omega \in \Lambda^{\infty}(U)$ define its kernel by ker(Ω) := { $\Lambda \in \mathcal{X}(U)$: $\Lambda \mathbf{I} \Omega = 0$ }.

For the remainder of this paper we will also assume all vector field distributions non-singular in the sense that their dimension is constant on U , and also that all one-forms have constant rank on U.

2Ideals, Cauchy characteristics and symme-

FOILOWING BIVANT *et al.* [2], for any $\alpha^*, \ldots, \alpha^r \in \Lambda(U)$ up to some $p \in \mathbb{N}$, we write $I := \langle \alpha^*, \ldots, \alpha^r \rangle$ to mean that I is the (nomogeneous) algebraic ideal generated by the elements $\alpha^*,\ldots,\alpha^r.$ An ideal I is a $\emph{aifferential}$ ideal if the exterior derivative of every member of I is also in I . A vector field Y is called a *Cauchy characteristic* vector field of an ideal I if $Y \cup I \subset I$. Define $A(I)$ to be the set of all Cauchy characteristic vector fields of I . It is not hard to show that $A(I)$ is Frobenius integrable.

A vector field $X \in \mathfrak{X}(U)$ is said to be a *symmetry* (or *isovector*) of an ideal, I, if $\mathcal{L}_X I \subset I$. It is easy to see that in order to show that X is a symmetry of I , it is enough to show that the Lie derivative with respect to X of merely the generators of I, is also in I. A vector field $X \in \mathfrak{X}(U)$ is a symmetry of a vector field distribution $D \subset \mathfrak{X}(U)$ if $\mathcal{L}_X D \subset D$. Once again, it is enough to look at simply the generators of D when determining whether a vector field is a symmetry of the distribution.

We now present some results connecting symmetries, ideals, and Cauchy characteristic spaces.

Proposition 2.1. Let I be an ideal. Suppose $A(I)$ is not zero-dimensional. If a vector field X is a symmetry of I then X is a symmetry of $A(I)$.

Proof. Let X be a symmetry of the ideal I. Let $Y \in A(I)$ and $\beta \in I$. Then, from rearranging the identity $\mathcal{L}_X(Y \cup \beta) = [X, Y] \cup \beta + Y \cup (\mathcal{L}_X \beta)$, we obtain

$$
[X,Y]\Box \beta = \mathcal{L}_X(Y \Box \beta) - Y \Box (\mathcal{L}_X \beta).
$$

Now the first term on the right hand side is in I since $Y \perp \beta \in I$ and X is a symmetry of I. The second term is also in I since $\mathcal{L}_X\beta \in I$ and $Y \in A(I)$. Hence $[X, Y] \cup \beta \in I$. Therefore $[X, Y] \in A(I)$. \Box

Proposition 2.2. Let $\{\alpha_1, \ldots, \alpha_r\}$ be some finite set of unearly independent one-forms in $\Lambda^*(U)$, and define $I := \langle \alpha^*, \ldots, \alpha^r, a\alpha^*, \ldots, a\alpha^r \rangle$. With the awat space of the Pfaffian system generated by α ,..., α defined by $D :=$ $\{X \in \mathfrak{X}(U) : X \rvert \alpha^i = 0, \text{ for all } 1 \leq i \leq p\},\$ then a vector field, $Y \in D$ is a Cauchy characteristic of I if and only if $[X, Y] \in D$ for all $X \in D$.

Proof. Let Y be a Cauchy characteristic vector field of I, i.e. $Y \Box \alpha^i = 0$ and $\chi_1 a \alpha^* \in I$ for all $1 \leq i \leq p$. This implies that for all i ,

$$
\mathcal{L}_Y \alpha^i = Y \Box d \alpha^i \in I.
$$

Hence Y is a symmetry of I. Let $X \in D$, where D is defined in the theorem. Using the property

$$
\mathcal{L}_Y(X \mathsf{J} \alpha^i) = [Y, X] \mathsf{J} \alpha^i + X \mathsf{J} \left(\mathcal{L}_Y \alpha^i \right),
$$

we know that the term on the left is zero and the second term on the right is also zero. Hence for all i, $[X, Y] \Box \alpha^i = 0$, so that $[X, Y] \in D$.

Conversely, let $Y \in D$ and $[X, Y] \in D$ for all $X \in D$. We therefore have that for all i ,

$$
Y \mathbf{1} \alpha^i = 0 = [X, Y] \mathbf{1} \alpha^i.
$$

Now on
e again using the property

$$
\mathcal{L}_Y(X \mathsf{J} \alpha^i) = [Y, X] \mathsf{J} \alpha^i + X \mathsf{J} \left(\mathcal{L}_Y \alpha^i \right),
$$

we have that for each i ,

$$
X \Box \left(\mathcal{L}_Y \alpha^i \right) = 0. \tag{1}
$$

Since (1) must note for all $A \in D$, we must have that $Ly \alpha^2 \in I$. Since $Y \mathbf{I} \alpha^i = 0,$

$$
\mathcal{L}_Y \alpha^i = Y \lrcorner \, d\alpha^i \in I,
$$

so Y is a Cauchy characteristic vector field of I .

At this point we will introduce the idea of a *trivial symmetry*. Given a $differential$ ideal I, we call all Cauchy characteristics of I trivial symmetries of I . The reason for this is contained in the next proposition:

Proposition 2.3. Let I be a differential ideal, and let Y be a Cauchy characteristic vector field of I . Then Y is a symmetry of I .

 \Box

Proof. let $\beta \in I$ and $Y \in A(I)$.

$$
\mathcal{L}_Y \beta = d(Y \mathbf{1} \beta) + Y \mathbf{1} d\beta.
$$

The first term on the right is in I because $Y \cup \beta \in I$, and consequently $d(Y \cup \beta) \in I$, since I is a differential ideal. The second term is obviously in I_{\cdot} \Box

Similarly, given a vector field distribution D , a *trivial symmetry* of D is a symmetry of D that is also in D.

A fundamental distin
tion between trivial and non-trivial symmetries is as follows: Given a trivial symmetry, multiplying it by any nononstant function will yield a trivial symmetry, however doing the same to a nontrivial symmetry will in general not produ
e a non-trivial symmetry.

For a differential ideal generated by a Pfaffian system we have the following extension of Proposition 2.3:

Proposition 2.4. Let I be a differential ideal generated by some finite coliection of unearly independent one-forms $\alpha^*,\ldots,\alpha^r \in \Lambda^*(U)$. A vector field $X \in \mathfrak{X}(U)$ is a symmetry of I in the annihilating space $D := \{X \in \mathfrak{X}(U) :$ $X \perp \alpha^{i} = 0$, for all $1 \leq i \leq p$ if and only if X is a trivial symmetry (Cauchy characteristic vector field) of I .

Proof. With X as a symmetry of I, if $X \perp \alpha^i = 0$ for all $1 \leq i \leq p$, then for ea
h i

$$
I \ni \mathcal{L}_X \alpha^i = X \lrcorner \, d\alpha^i.
$$

 \Box

The onverse is also obvious using Proposition 2.3.

Definition 2.5. A differential *p*-form said to be *decomposable* (or *simple*) if it can be written as the wedge product of p one-forms.

Decomposability is a local property, and a p -form is decomposable if and only if the dimension of the kernel is of dimension $n - p$.

Consider the following two simple propositions, the first of which is proved in Sherring and Prince [13]:

Proposition 2.6. A vector field $X \in \mathfrak{X}(U)$ is a symmetry of a decomposable m -form $\Omega \in \Lambda^m(U)$ if and only if Λ is a symmetry of $\ker(\Omega)$.

Proposition 2.7. Let $\Omega \in \Lambda^m(U)$ and $I := \langle \Omega, d\Omega \rangle$. If $d\Omega = 0$ mod Ω , then ker() = A(I).

Proof. First suppose ker() is not zero-dimensional, so that there exists a non-zero version and who were the communications with the contract who were contracted to the contract of the mod Ω , $W \square a \Omega \equiv W \square (\alpha \wedge \Omega) = (W \square \alpha) \wedge \Omega$ for some $\alpha \in \Lambda^{\perp}(U)$. Therefore $W \in A(I)$.

Now suppose $A(I)$ is not zero-dimensional. This means there exists a that is a contract the subset of \mathcal{S} is a contract of the form of the form in the form of \mathcal{S} some smooth $f \in C^{\infty}(U)$. Hence from the first part, $\Lambda \in \text{Ker}(\Omega)$.

 \blacksquare) is zero-dimensional, then \blacksquare This means Y 2= I , and hen
e Y 2= A(I). Therefore A(I) is zero-dimensional.

Finally, if A(I) is zero-dimensional, then Z 6= 0 for all Z ² X(U). \Box Hen
e ker() is zero-dimensional.

Using the above two results, we obtain the following extension to differential ideals thus giving us a ondition under whi
h the onverse of Proposition 2.1 holds true:

Proposition 2.8. Let 1 be a afferential ideal generated by some $\Omega \in \Lambda$ (U) with definition of the d zero-dimensional. Then X is a symmetry of I if and only if X is a symmetry of $A(I)$.

Proof. From Proposition 2.7, d = 0 mod implies that ker() = A(I). Hen
e the result follows from Proposition 2.6. \Box

 $\textit{Remark.}$ If $\Omega \in \Lambda$ (*U)* with $m = n$, and *I* is the differential ideal generated by (note d = 0), then any non-zero ve
tor eld in X(U) is a symmetry of I . Moreover, A(h i) is zero-dimensional, and therefore any non-zero ve
tor eld in X(U) is also a symmetry of a zero-dimensional A(h i).

3 The Frobenius theorem and integration via symmetry

First, we present a basic result:

Lemma 3.1. [5] Let $\Omega \in \Lambda$ (U) for some $m \leq n - 1$. Then ker(Ω) can be at most m-dimensional. Moreover, ker() is pre
isely m-dimensional if and only if the contract of the

Lemma 3.1 has the following corollary:

Corollary 3.2. Let $D := sp\{Y_1, \ldots, Y_m\}$ be some m-dimensional distribution in $\mathfrak{X}(U)$, where $m \leq n-1$. If $\Omega := Y_1 \cup \ldots \cup Y_m \cup (ax^2 \wedge \cdots \wedge ax^2) \in$ Λ^{--} (U), then Ω is decomposable and equal to the wedge product of some $n - m$ unearly independent generators of $D^-.$

Proof. With D and dened as in the orollary, let X ² X(U) be any non-zero vez eld in D. Then from the definition of the definition of the definition of the definition of the d ker() is at least m-dimensional. But from Lemma 3.1, sin
e is an (nm) form, its kernel an not be greater than m-dimensional, and therefore is de
omposable.

Now we can write $\Omega = \theta^* \wedge \cdots \wedge \theta^*$ for some imearly independent $\sigma^2, \ldots, \sigma^{n-1} \in \Lambda^2(U)$. Since for each $1 \leq i \leq m$, $Y_i \mathbf{I} \setminus \mathbf{I} \subset \mathbf{I}$, we then have that for each $1 \leq j \leq (n-m)$, $Y_i \perp b^j = 0$. Hence b^2, \ldots, b^m m generate \Box $D^-.$

Theorem 3.3. (Frobenius) Let D be an m-dimensional distribution generated by the vector fields $Y_1, \ldots, Y_m \in \mathfrak{X}(U)$, where $m \leq n-1$. Define D^- to be the submodule of all one-forms that annihilate D . Let Ω := Y_1 , Y_2 , Y_m , \exists $\{ax^m \wedge \cdots \wedge ax^m\} \in \Lambda^m$ (U). Then D has m-aimensional inte $gral submanifolds on U if and only if either of the following two equivalent$ onditions are true:

- 1. For all $X, Y \in D$, $[X, Y] \in D$,
- 2. For all $\theta \in D^{\perp}$, $d\theta \wedge \Omega = 0$. , . . , , . .

We say that a distribution D is Frobenius integrable (or generates a foli*ation* of U) if the first condition in the Frobenius theorem holds. The Frobenius theorem means that D generates an m -dimensional foliation of U whose leaves are described by some set of $n-m$ functions $\gamma^1 = c_1, \ldots, \gamma^{n-m} = c_{n-m}$ of rank $n-m$, where $\gamma^*, \ldots, \gamma^*$ are consequently and c_1, \ldots, c_{n-m} are some appropriate onstant fun
tions.

Using Corollary 3.2, we have the following corollary to the Frobenius theorem:

Corollary 3.4. Let D be an m-dimensional distribution generated by the vector jieus $Y_1,\ldots,Y_m\in\mathfrak{X}(U)$, where $m\leq n-1$. Let $\mathfrak{U}:=Y_1\mathfrak{z}\ldots\mathfrak{z}Y_m\mathfrak{z}(ax)$ $\cdots \wedge ax \in \Lambda$ (U). For all $b \in D$, ab $\wedge \Omega = 0$ (i.e. D is Frobenius integrable) if and only if and only if the contract of the contract of the contract of the contract of the con

Proof. With Ω defined as in the corollary, Corollary 3.2 implies $\Omega = \theta^*$ /\ $\cdots \wedge \sigma$ are not some imearly independent $\sigma^1, \ldots, \sigma^n$ are $\Lambda^1(U)$ that generate D^+ . Now for each $1 \leq i \leq (n-m)$, the Frobenius condition $u\sigma/\Delta u = 0$ is equivalent to the condition that $a\theta^+ = 0$ mod θ^- ,..., θ^+ Hence

$$
d\Omega = d\left(\theta^1 \wedge \cdots \wedge \theta^{n-m}\right),
$$

= 0 \mod \Omega.

To prove the onverse, suppose d = 0 mod . Now for all i,

 \mathcal{L}

$$
d\theta^i \wedge \Omega = d\left(\theta^i \wedge \Omega\right) + \theta^i \wedge d\Omega. \tag{2}
$$

Since $\theta^* \wedge \Omega = 0$, and Ω is closed modulo itself, we find from (2) that $d\theta^* \wedge \Omega =$ $\overline{0}$. \Box

From Sherring and Prince [13] we have the following definition:

Definition 3.5. A differential m -form $\alpha \in \Lambda$ (U) is *Probenius integrable* if ker() is Frobenius integrable and of dimension n m.

From this definition we have the following lemma:

Lemma 3.0. A differential m-form $\Omega \in \Lambda^{\infty}(U)$ is Provenius integrable if and only if the distribution of the distribution of the distribution of the second second second second second

Proof. First suppose $\Omega \in \Lambda$ (U) is Frobenius integrable. By definition, ker() is maximal dimension, and hen
e is de
omposable. We an write $\Omega = \sigma^* \wedge \ldots \sigma^*$ for some $\sigma^*, \ldots \sigma^*$ if $\Lambda^*(U)$. Since ker(Ω) is frobenius integrable, it follows that for each $1 \leq i \leq n-m$, $a\sigma =$ mod $\sigma^2, \ldots, \sigma^{n-m}$.

composes it is a distinct that the definition of the second contract α is form α is a rank-rank. Further, the state α is α integrable, the α is from α integrable. from Proposition 2.7. \Box

Theorem 3.7. Let $\Omega \in \Lambda^m(U)$ for some $m > 1$ be decomposable, and let $\Lambda \in \mathfrak{X}(U)$ with the property $\Lambda \rightrightarrows \Omega \neq 0$. Then there exists $\theta \in \Lambda^*(U)$ such \cdots \cdots \cdots \cdots

Proof. Let $\Omega \in \Lambda^m(U)$ be decomposable, and let $\Lambda \in \mathcal{X}(U)$ with $\Lambda \downarrow \Omega \neq 0$. τ . τ , τ , τ , τ , τ , τ , τ are a basis for the vector τ . The vector τ fields X, Y_{m+1}, \ldots, Y_n are linearly independent. We can extend these vector neids to a basis by including some $r_2, \ldots, r_m \in \mathcal{X}$. Let $\{ \theta^*, \ldots, \theta^m \}$ be a dual basis of one forms for $\{A, Y_2, \ldots, Y_m\}$. Then $\Omega = f \sigma^2 \wedge \cdots \wedge \sigma^m$, and \Box moreover, $\Lambda \exists M \equiv U U^T \wedge \cdots \wedge U^T$. Hence the result follows.

By an obvious iteration, we have the following corollary to Theorem 3.7:

Corollary 3.8. Let $\Omega \in \Lambda^{\infty}(U)$ be decomposable. Let $\Lambda_1, \ldots, \Lambda_p \in \mathfrak{X}(U)$ up to some $p < m$ such that Λ_1 , \ldots , Λ_p , $\iota \neq 0$. Then there exist $\sigma^2, \ldots, \sigma^r \in$ A (U) such that

$$
\Omega = \theta^p \wedge \cdots \wedge \theta^1 \wedge (X_1 \sqcup \cdots \sqcup X_p \sqcup \Omega),
$$

\n
$$
X_p \sqcup \Omega = \theta^{p-1} \wedge \cdots \wedge \theta^1 \wedge (X_1 \sqcup \cdots \sqcup X_p \sqcup \Omega),
$$

\n
$$
\vdots
$$

\n
$$
X_2 \sqcup \cdots \sqcup X_p \sqcup \Omega = \theta^1 \wedge (X_1 \sqcup \cdots \sqcup X_p \sqcup \Omega).
$$

Proposition 3.9. Let $\Omega \in \Lambda^m(U)$ be decomposable, and let $\Lambda \in \mathfrak{X}(U)$ such that X 6= 0. Then ker(X) = ker() spfXg.

Proof. It is lear that ker(X) ker(). Sin
e X ² ker(X), we therefore have the form of ρ as the form ρ as a spirit of ρ assumption to the contract of the solution of the solution of ρ Lemma 3.1 implies ker() has maximal dimension n m. Sin
e X 2= ker(), it follows that ker() spfXg has dimension n m + 1. Hen
e Lemma 3.1 \Box implies X is de
omposable.

We have the following corollary to Proposition 3.9, which can also be found in Sherring and Prince [13]:

Corollary 3.10. Let $\Omega \in \Lambda^m(U)$ for some $m > 1$ be aecomposable, and let a that is defined in the contract of the contract of the contract of the α

Before we present the next result, we require the following central definition:

Definition 3.11. Let D be a distribution in $\mathfrak{X}(U)$. Then a set of p linearly independent vector fields, $X_1, \ldots, X_p \in \mathfrak{X}(U)$, form a solvable symmetry structure for D if

$$
\mathcal{L}_{X_1} \left(sp\{X_2, \ldots, X_p\} \oplus D \right) \subset sp\{X_2, \ldots, X_p\} \oplus D,
$$

$$
\vdots
$$

$$
\mathcal{L}_{X_{p-1}} \left(sp\{X_p\} \oplus D \right) \subset sp\{X_p\} \oplus D,
$$

$$
\mathcal{L}_{X_p} D \subset D.
$$

Theorem 3.12. Let $\Omega \in \Lambda^m(U)$ be *Frobenius integrable. Further, let* $\Lambda \in$ X(U) su
h that A(h i) spfXg is Frobenius integrable and X 6= 0. Then X is Frobenius integrable.

Proof. This theorem is obvious from Definition 3.5, Propositions 2.7 and 3.7, and Corollary 3.10. \Box

We have the following corollary to Theorem 3.12:

Corollary 3.13. Let $\Omega \in \Lambda^m(U)$ be *Frobenius integrable, and suppose there* exist X1; : : : ; Xp ² X(U) up to some ^p < ^m su
h that X1 : : : Xp 6= 0. \mathcal{I} is a frobenius integrable distribution, and for all late \mathcal{I} and \mathcal{I} and \mathcal{I} is and \mathcal{I} α , β , β , γ , γ is also found integrable, then α is also found integrating integrating γ : : ; , X : : : : Y : Y : Y : Y integrable. Moreover, functional contracts integrable. The set of Y a solvable symmetry structure for self-rily if and only if

$$
\mathcal{L}_{X_p}\Omega = \lambda_p \Omega,
$$

\n
$$
\mathcal{L}_{X_{p-1}}(X_p \mathbf{1}\Omega) = \lambda_{p-1}(X_p \mathbf{1}\Omega),
$$

\n
$$
\vdots
$$

\n
$$
\mathcal{L}_{X_1}(X_2 \mathbf{1} \dots \mathbf{1} X_p \mathbf{1}\Omega) = \lambda_1(X_2 \mathbf{1} \dots \mathbf{1} X_p \mathbf{1}\Omega),
$$
\n(3)

for some $\lambda_1, \ldots, \lambda_p \in \cup \cup$.

Corollary 3.13 provides a direct connection between a solvable symmetry structure for the equation $\mathcal{L}^{(t)}$ and $\mathcal{L}^{(t)}$ and $\mathcal{L}^{(t)}$ and $\mathcal{L}^{(t)}$ and $\mathcal{L}^{(t)}$ and $\mathcal{L}^{(t)}$ frequently referred to as a solvable symmetry stru
ture for).

The papers by Sherring and Prince [13] and Basarab-Horwath [1] extend Lie's approa
h to integrating a Frobenius integrable distribution via a solvable stru
ture of symmetries. In those papers, a Frobenius integrable distribution is given first. The one-form annihilating space is then generated and all generators wedged to give a de
omposable form with a Frobenius integrable kernel. The result is reprodu
ed below:

Theorem 3.14. [13] Let $D := sp\{Y_1, \ldots, Y_q\} \subset \mathfrak{X}(U)$ be a q-dimensional Frobenius integrable vector field distribution. Define $\Omega := Y_1 \cup \ldots \cup Y_d \cup (dx^1 \wedge$ $\cdots \wedge dx^n) \in \Lambda^{n-q}(U)$, and suppose there exists a solvable structure of linearly independent symmetries $X_1, \ldots, X_{n-q} \in \mathfrak{X}(U)$ such that X_{n-q} is a non-trivial symmetry of D, and that for all $1 \leq i < n-q$, X_i is a non-trivial symmetry of $D \oplus sp\{X_{i+1},\ldots,X_{n-q}\}$. For all $1 \leq i \leq n-q$, define ω^i by

$$
\omega^{i} := \frac{X_{1} \mathbf{1} \dots \mathbf{1} X_{i-1} \mathbf{1} X_{i+1} \mathbf{1} \dots \mathbf{1} X_{n-q} \mathbf{1} \Omega}{X_{i} \mathbf{1} X_{1} \mathbf{1} \dots \mathbf{1} X_{i-1} \mathbf{1} X_{i+1} \mathbf{1} \dots \mathbf{1} X_{n-q} \mathbf{1} \Omega}
$$

Then $\{\omega^1,\ldots,\omega^{n-q}\}\$ is dual to $\{X_1,\ldots,X_{n-q}\}\$, and for all ω^i up to $i=n-q$,

$$
\omega^1 = d\gamma^1,
$$

\n
$$
\omega^2 = d\gamma^2 - X_1(\gamma^2)d\gamma^1,
$$

\n
$$
\omega^3 = d\gamma^3 - X_2(\gamma^3)(d\gamma^2 - X_1(\gamma^2)d\gamma^1) - X_1(\gamma^3)d\gamma^1,
$$

\n
$$
\vdots
$$

\n
$$
\omega^{n-q} = d\gamma^{n-q} \mod d\gamma^1, \dots, d\gamma^{n-q-1},
$$

for some functionally independent $\gamma^1, \ldots, \gamma^{n-q} \in C^{\infty}(U)$. Moreover, on U, the submanifolds described by D generate a q-dimensional foliation of U whose leaves have $\gamma^1, \ldots, \gamma^{n-q}$ constant.

In our work, we will start with a decomposable m -form Ω with a Frobenius integrable kernel. This is achieved by also demanding that $d\Omega = 0 \text{ mod } \Omega$. Hence by Proposition 2.7, the Cauchy characteristic space of the differential ideal generated by Ω is Frobenius integrable and equal to ker(Ω). Using these facts, we show below in Theorem 3.15 how a solvable structure of symmetries for Ω (as in Corollary 3.13) can assist in generating a simplified expression for Ω . Theorem 3.15 is the key result of this paper.

Theorem 3.15. Let $\Omega \in \Lambda^m(U)$ be Frobenius integrable. Suppose there exists a solvable structure of linearly independent symmetries $X_1, \ldots, X_m \in \mathfrak{X}(U)$ such that X_m is a non-trivial symmetry of $A(\langle\Omega\rangle)$, and that for all $1 \leq i < m$, X_i is a non-trivial symmetry of $A(\langle\Omega\rangle)\oplus sp\{X_{i+1},\ldots,X_m\}$. For all $1\leq i\leq n$ m, define ω^i by

$$
\omega^{i} := \frac{X_{1}\mathbf{1} \dots \mathbf{1} X_{i-1}\mathbf{1} X_{i+1}\mathbf{1} \dots \mathbf{1} X_m \mathbf{1} \Omega}{X_{i}\mathbf{1} X_{1}\mathbf{1} \dots \mathbf{1} X_{i-1}\mathbf{1} X_{i+1}\mathbf{1} \dots \mathbf{1} X_m \mathbf{1} \Omega}.
$$
(4)

Then $\{\omega^1,\ldots,\omega^m\}$ is dual to $\{X_1,\ldots,X_m\}$, and for all ω^i up to $i=m$,

$$
\omega^{1} = d\gamma^{1},
$$

\n
$$
\omega^{2} = d\gamma^{2} - X_{1}(\gamma^{2})d\gamma^{1},
$$

\n
$$
\omega^{3} = d\gamma^{3} - X_{2}(\gamma^{3})(d\gamma^{2} - X_{1}(\gamma^{2})d\gamma^{1}) - X_{1}(\gamma^{3})d\gamma^{1},
$$

\n
$$
\vdots
$$

\n
$$
\omega^{m} = d\gamma^{m} \mod d\gamma^{1}, \dots, d\gamma^{m-1},
$$

\n(5)

for some functionally independent $\gamma^1, \ldots, \gamma^m \in C^{\infty}(U)$. Finally, define $\gamma^0 :=$ $\Omega(X_1,\ldots,X_m)$. Then $\Omega = \gamma^0 d \gamma^1 \wedge \cdots \wedge d \gamma^m$.

Proof. Since from Lemma 3.6, Ω is decomposable, we may write $\Omega = \theta^1 \wedge$ $\cdots \wedge \theta^m$ for some linearly independent $\theta^1, \ldots, \theta^m \in \Lambda^1(U)$. Now ker (Ω) = $sp{Y_1,\ldots,Y_{n-m}}$ for some $Y_1,\ldots,Y_{n-m} \in \mathfrak{X}(U)$. From Lemma 3.6 and Proposition 2.7, we have that $A(\langle \Omega \rangle) = \ker(\Omega)$ is Frobenius integrable. Applying Theorem 3.14 with the linearly independent symmetries $X_1, \ldots, X_m \in$ $\mathfrak{X}(U)$ for $A(\langle \Omega \rangle)$ given in Theorem 3.15, we obtain that

$$
\{Y_1,\ldots,Y_{n-m},X_1,\ldots,X_m\}
$$

spans $\mathfrak{X}(U)$ and is dual to

$$
\left\{\phi^1,\ldots,\phi^{n-m},\omega^1,\ldots,\omega^m\right\},\,
$$

for some linearly independent $\phi^1, \ldots, \phi^{n-m} \in \Lambda^1(U)$ with $\omega^1, \ldots, \omega^m$ defined as in (4). Since $Y_j \square \Omega = 0$ for all $1 \le j \le n - m$, it follows that

$$
\Omega = \Omega(X_1, \dots, X_m) \omega^1 \wedge \dots \wedge \omega^m. \tag{6}
$$

Now Theorem 3.14 implies the equations in (5) , so (6) simplifies to give

$$
\Omega = \Omega(X_1, \ldots, X_m) d\gamma^1 \wedge \cdots \wedge d\gamma^m.
$$

Remark 1. The fact that the symmetries in Theorem 3.15 are non-trivial means that the denominator is non-zero in each of the definitions for ω^{i} .

Remark 2. The expression for γ^0 is easily derived since $\Omega = \Omega(X_1, \ldots, X_m)$ $\omega^1 \wedge \cdots \wedge \omega^m$ as $Y_j \Omega = 0$ for all $1 \leq j \leq (n-m)$ linearly independent vector fields Y_1, \ldots, Y_{n-m} in $A(\langle \Omega \rangle)$ that are used with X_1, \ldots, X_m to span $\mathfrak{X}(U)$.

Theorem 3.15, for a given Ω and solvable symmetry structure of vector fields, gives us explicit expressions for the relations described in Proposition 3 in [9].

In later sections, we will illustrate Theorem 3.15 with some applications. For now though, we have the following consequence of Theorem 3.15 regarding the its second remark:

Theorem 3.16. Given some Frobenius integrable $\Omega \in \Lambda^m(U)$ and a solvable structure $X_1, \ldots, X_m \in \mathfrak{X}(U)$ for $A(\langle \Omega \rangle)$ as in Theorem 3.15, then

$$
\mathcal{L}_{X_m}\Omega = \{X_m \cup d(\ln |\Omega(X_1, \dots, X_m)|)\}\Omega,
$$

$$
\mathcal{L}_{X_{m-1}}(X_m \cup \Omega) = \{X_{m-1} \cup d(\ln |\Omega(X_1, \dots, X_m)|)\}(X_m \cup \Omega),
$$

$$
\vdots
$$

$$
\mathcal{L}_{X_1}(X_2 \cup \dots \cup X_m \cup \Omega) = \{X_1 \cup d(\ln |\Omega(X_1, \dots, X_m)|)\}(X_2 \cup \dots \cup X_m \cup \Omega).
$$

Proof. First we will show that for all $1 \leq i \leq m$, $a(\omega^2 \wedge \cdots \wedge \omega^2) = 0$. From Theorem 3.15 it is obvious that $d\omega^1 = 0$ and for each $1 < i \leq m$ that $d\omega^i = 0$ $\mod \omega^-, \ldots, \omega^+$. Thus for all $i > 1$,

$$
d(\omega^1 \wedge \cdots \wedge \omega^i) = 0. \tag{7}
$$

From Theorem 3.15 it is lear that

$$
\Omega = \Omega(X_1, \dots, X_m) \omega^1 \wedge \dots \wedge \omega^m.
$$
 (8)

Hen
e

$$
d\left(\frac{\Omega}{\Omega(X_1,\ldots,X_m)}\right) = 0.
$$
\n(9)

Using that $\{\omega^*, \ldots, \omega^{\dots}\}\$ is dual to $\{\Lambda_1, \ldots, \Lambda_m\}\$ and contracting (8) with X_m , we obtain

$$
\omega^1 \wedge \cdots \wedge \omega^{m-1} = \frac{X_{m1} \Omega}{(-1)^{m-1} \Omega(X_1, \ldots, X_m)}
$$

From repeating this contraction with X_{m-1} and so on down to X_1 , we obtain for all $1 \leq i \leq m - 1$,

$$
\omega^1 \wedge \cdots \wedge \omega^i = \frac{X_{i+1} \cup \ldots \cup X_m \cup \Omega}{(-1)^{((m-1)+\cdots+i)} \Omega(X_1, \ldots, X_m)}.
$$

Hen
e from (7),

$$
d\left(\frac{X_{i+1}\mathbf{J}\dots\mathbf{J}X_m\mathbf{J}\Omega}{(-1)^{((m-1)+\dots+i)}\Omega(X_1,\dots,X_m)}\right) = 0. \tag{10}
$$

Equation (9) implies

$$
d\Omega = d\left(\ln|\Omega(X_1,\ldots,X_m)|\right) \wedge \Omega,\tag{11}
$$

while equation (10) means

$$
d(X_{i+1},\ldots,X_m)\Omega)=d(\ln |\Omega(X_1,\ldots,X_m)|)\wedge (X_{i+1},\ldots,X_m)\Omega), \quad (12)
$$

for all $1 \leq i \leq (m - 1)$. Now

$$
\mathcal{L}_{X_m}\Omega = X_m \mathbf{1} d\Omega + d\left(X_m \mathbf{1}\Omega\right),
$$

= $X_{m}\mathbf{1} \{d\left(\ln |\Omega(X_1, \ldots, X_m)|\right) \wedge \Omega\} + d\left(\ln |\Omega(X_1, \ldots, X_m)|\right) \wedge \left(X_m \mathbf{1}\Omega\right),$
= $\{X_m \mathbf{1} d(\ln |\Omega(X_1, \ldots, X_m)|)\}\Omega,$

where in the second line we have inserted equations (11) and (12) . To obtain the third line we used the identity Λ ($\omega \wedge \theta$) = (Λ ω) \wedge θ + (-1) σ $(X \cup \sigma)$ for differential forms σ, ω .

Finally, let $1 \leq i \leq (m-1)$. Then in a similar fashion to before, we get

$$
\mathcal{L}_{X_i}(X_{i+1}\mathbf{I}\ldots\mathbf{I} X_m\mathbf{I}\Omega)=X_i\mathbf{I}\{d(\ln|\Omega(X_1,\ldots,X_m)|)\wedge(X_{i+1}\mathbf{I}\ldots\mathbf{I} X_m\mathbf{I}\Omega)\}+d(\ln|\Omega(X_1,\ldots,X_m)|)\wedge(X_i\mathbf{I}\ldots\mathbf{I} X_m\mathbf{I}\Omega),
$$

which simplifies to

$$
\mathcal{L}_{X_i}(X_{i+1}\cup\ldots\cup X_m\cup\Omega)=\{X_i\cup d(\ln|\Omega(X_1,\ldots,X_m)|)\}(X_{i+1}\cup\ldots\cup X_m\cup\Omega).
$$

In general, each $\omega^2, \ldots, \omega^m$ in Theorem 3.15 is not exact. Our final results for this section examine some conditions on the symmetries X_1, \ldots, X_m in Theorem 3.15 that force at least one of $\omega^2, \ldots, \omega^m$ to be exact.

Theorem 3.17. Let $\Omega \in \Lambda^m(U)$ for some $m \geq 3$ such that Ω is Frobenius integrable. Let there exist a solvable structure of linearly independent symmetries $X_3, \ldots, X_m \in \mathfrak{X}(U)$ such that X_m is a non-trivial symmetry of $A(\langle\Omega\rangle)$, and that for all $3\leq i < m$, X_i is a non-trivial symmetry of $A(\langle\Omega\rangle) \oplus$ $sp{X_{i+1},...,X_m}$. Also, let there exist two linearly independent vector fields $X_1, X_2 \in \mathfrak{X}(U)$ that are non-trivial symmetries of $A(\langle \Omega \rangle) \oplus sp\{X_3, \ldots, X_m\}$ such that

$$
[X_1, X_2] = 0 \mod A(\langle \Omega \rangle) \oplus sp\{X_3, \dots, X_m\}.
$$
 (13)

For all $1 \leq i \leq m$, define ω^i by

$$
\omega^{i} := \frac{X_{1} \cup \ldots \cup X_{i-1} \cup X_{i+1} \cup \ldots \cup X_{m} \cup \Omega}{X_{i} \cup X_{1} \cup \ldots \cup X_{i-1} \cup X_{i+1} \cup \ldots \cup X_{m} \cup \Omega}
$$

Then $\{\omega^1,\ldots,\omega^m\}$ is dual to $\{X_1,\ldots,X_m\}$ and for all ω^i up to $i=m$,

$$
\omega^1 = d\gamma^1,
$$

\n
$$
\omega^2 = d\gamma^2,
$$

\n
$$
\omega^3 = d\gamma^3 - X_2(\gamma^3)d\gamma^2 - X_1(\gamma^3)d\gamma^1,
$$

\n
$$
\omega^4 = d\gamma^4 - X_3(\gamma^4)(d\gamma^3 - X_2(\gamma^3)d\gamma^2 - X_1(\gamma^3)d\gamma^1) - X_2(\gamma^4)d\gamma^2 - X_1(\gamma^4)d\gamma^1,
$$

\n
$$
\vdots
$$

\n
$$
\omega^m = d\gamma^m \mod d\gamma^1, \dots, d\gamma^{m-1},
$$

for some functionally independent $\gamma^1, \ldots, \gamma^m \in C^{\infty}(U)$. Finally, define $\gamma^0 :=$ $\Omega(X_1,\ldots,X_m)$. Then $\Omega = \gamma^0 d \gamma^1 \wedge \cdots \wedge d \gamma^m$.

Proof. We begin by showing that X_1 is a non-trivial symmetry of $A(\langle \Omega \rangle) \oplus$ $sp{X_2,\ldots,X_m}$. Since X_1 is a non-trivial symmetry of $A(\langle\Omega\rangle)\oplus sp{X_3,\ldots,X_m}$, we have from Corollary 3.13 that

$$
\mathcal{L}_{X_1}(X_3 \cup \ldots \cup X_m \cup \Omega) = \lambda(X_3 \cup \ldots \cup X_m \cup \Omega),
$$

for some $\lambda \in C^+(U)$. Using this fact and equation (13) then gives

$$
\mathcal{L}_{X_1}(X_2 \cup \ldots \cup X_m \cup \Omega) = [X_1, X_2] \cup X_3 \cup \ldots \cup X_m \cup \Omega + X_2 \cup \mathcal{L}_{X_1}(X_3 \cup \ldots \cup X_m \cup \Omega),
$$

= $\lambda (X_2 \cup \ldots \cup X_m \cup \Omega).$

From Theorem 3.13, our symmetries at this point satisfy Theorem 3.15. Therefore

$$
\omega^1 = d\gamma^1,
$$

\n
$$
\omega^2 = d\gamma^2 - X_1(\gamma^2)d\gamma^1,
$$

\n
$$
\omega^3 = d\gamma^3 - X_2(\gamma^3)(d\gamma^2 - X_1(\gamma^2)d\gamma^1) - X_1(\gamma^3)d\gamma^1,
$$

\n
$$
\vdots
$$

\n
$$
\omega^m = d\gamma^m \mod d\gamma^1, \dots, d\gamma^{m-1},
$$

for some functionally independent $\gamma^*,\ldots,\gamma^m\in C^\infty(U).$ To show that $\Lambda_1(\gamma^*)\equiv$ 0, we must show that

$$
d\omega^2 = d\left(\frac{X_1 \mathbf{1} X_3 \mathbf{1} \dots \mathbf{1} X_m \mathbf{1} \Omega}{X_2 \mathbf{1} X_1 \mathbf{1} X_3 \mathbf{1} \dots \mathbf{1} X_m \mathbf{1} \Omega}\right) = 0.
$$
 (14)

This case of proved by observing that since the χ - χ is the η if χ is χ is χ is χ if χ $sp{X_1, X_3, \ldots, X_m}$ is a Frobenius integral distribution, we therefore have that

$$
d(X_1 \mathbf{1} X_3 \mathbf{1} \dots X_m \mathbf{1} \Omega) = 0 \mod X_1 \mathbf{1} X_3 \mathbf{1} \dots \mathbf{1} X_m \mathbf{1} \Omega.
$$

Then to show that X2 is ^a non-trivial symmetry of A(h i)spfX1; X3; : : : ; Xmg we use the formula

$$
\mathcal{L}_{X_2}(X_1 \cup X_3 \cup \ldots \cup X_m \cup \Omega) = [X_2, X_1] \cup X_3 \cup \ldots \cup X_m \cup \Omega + X_1 \cup \mathcal{L}_{X_2}(X_3 \cup \ldots \cup X_m \cup \Omega).
$$

Now using equation (13) and that X_2 is a non-trivial symmetry of $X_3 \perp \ldots \perp$ \mathbf{w} , we deduce the desired result. Equation (14) and \mathbf{w} simple algebraic manipulation, or by applying Theorem 3.15. \Box

Remark. While Theorem 3.17 assumes $m \geq 3$, it is clear that is still holds when $m = 2$. In this situation, there is no need for symmetries other than ing to $\begin{array}{ccc} 1 & 1 & 1 \end{array}$ reduces to $\begin{array}{ccc} 1 & 1 & 1 \end{array}$ and the expression of the expre sions for ω in the conclusion of the theorem vanish for $i > 2$.

We can generalise Theorem 3.17 in the following way:

THEOREM 3.18. Let $\Omega \in \Lambda$ (U) for some $m > 3$, and suppose Ω is *Frobenius* integrable. For some $1 \leq l < m$, let there exist a solvable structure of $m - l$ linearly independent symmetries $X_{l+1}, \ldots, X_m \in \mathfrak{X}(U)$ such that X_m is a is a constant for all lines of A is and that for all lines of A is a constant for all lines of A

non-trivial symmetry of $A(\langle \Omega \rangle) \oplus sp\{X_{i+1},\ldots,X_m\}$. Also, let there exist *l* linearly independent vector fields $X_1, \ldots, X_l \in \mathfrak{X}(U)$ that are non-trivial symmetries of $A(\langle \Omega \rangle) \oplus sp\{X_{l+1},\ldots,X_m\}$ such that

$$
[X_u, X_v] = 0 \mod A(\langle \Omega \rangle) \oplus sp\{X_{l+1}, \dots, X_m\},\tag{15}
$$

for all $1 \leq u < v \leq l$. For all $1 \leq i \leq m$, define ω^{i} by

$$
\omega^i := \frac{X_{1\mathbf{\cup}} \dots \mathbf{\cup} X_{i-1\mathbf{\cup}} X_{i+1\mathbf{\cup}} \dots \mathbf{\cup} X_m \mathbf{\cup} \Omega}{X_{i\mathbf{\cup}} X_{1\mathbf{\cup}} \dots \mathbf{\cup} X_{i-1\mathbf{\cup}} X_{i+1\mathbf{\cup}} \dots \mathbf{\cup} X_m \mathbf{\cup} \Omega}
$$

Then $\{\omega^1,\ldots,\omega^m\}$ is dual to $\{X_1,\ldots,X_m\}$ and for all ω^i up to $i=l$,

$$
\omega^1 = d\gamma^1,
$$

$$
\vdots
$$

$$
\omega^l = d\gamma^l,
$$

with for each i greater than l up to $i = m$,

$$
\omega^{l+1} = d\gamma^{l+1} - X_l(\gamma^{l+1})d\gamma^l - X_{l-1}(\gamma^{l+1})d\gamma^{l-1} - \dots - X_1(\gamma^{l+1})d\gamma^1,
$$

\n
$$
\omega^{l+2} = d\gamma^{l+2} - X_{l+1}(\gamma^{l+2})\left(d\gamma^{l+1} - X_l(\gamma^{l+1})d\gamma^l - X_{l-1}(\gamma^{l+1})d\gamma^{l-1} - \dots - X_1(\gamma^{l+2})d\gamma^1\right) - X_l(\gamma^{l+2})d\gamma^l - \dots - X_1(\gamma^{l+2})d\gamma^1,
$$

\n
$$
\vdots
$$

\n
$$
\omega^m = d\gamma^m \mod d\gamma^1, \dots, d\gamma^{m-1},
$$

for some functionally independent $\gamma^1, \ldots, \gamma^m \in C^{\infty}(U)$. Finally, define $\gamma^0 :=$ $\Omega(X_1,\ldots,X_m)$. Then $\Omega = \gamma^0 d\gamma^1 \wedge \cdots \wedge d\gamma^m$.

Proof. (Outline) The proof is similar to that of Theorem 3.17, and essentially involves repeating the proof of Theorem 3.17 $l-1$ more times. To do this, from the fact that Ω is decomposable and $d\Omega = 0$ mod Ω , we can then apply Corollary 3.13 to obtain

$$
\mathcal{L}_{X_m} = \lambda_m \Omega,
$$

\n
$$
\mathcal{L}_{X_{m-1}} (X_m \cup \Omega) = \lambda_{m-1} (X_m \cup \Omega),
$$

\n
$$
\vdots
$$

\n
$$
\mathcal{L}_{X_{l+1}} (X_{l+2} \cup \ldots \cup X_m \cup \Omega) = \lambda_{l+1} (X_{l+2} \cup \ldots \cup X_m \cup \Omega),
$$

and also that

$$
\mathcal{L}_{X_{l}}(X_{l+1}\mathbf{1}X_{l+2}\mathbf{1}\ldots\mathbf{1}X_{m}\mathbf{1}\Omega)=\lambda_{l}(X_{l+1}\mathbf{1}X_{l+2}\mathbf{1}\ldots\mathbf{1}X_{m}\mathbf{1}\Omega),
$$
\n
$$
\mathcal{L}_{X_{l-1}}(X_{l+1}\mathbf{1}X_{l+2}\mathbf{1}\ldots\mathbf{1}X_{m}\mathbf{1}\Omega)=\lambda_{l-1}(X_{l+1}\mathbf{1}X_{l+2}\mathbf{1}\ldots\mathbf{1}X_{m}\mathbf{1}\Omega),
$$
\n
$$
\vdots
$$
\n
$$
\mathcal{L}_{X_{1}}(X_{l+1}\mathbf{1}X_{l+2}\mathbf{1}\ldots\mathbf{1}X_{m}\mathbf{1}\Omega)=\lambda_{1}(X_{l+1}\mathbf{1}X_{l+2}\mathbf{1}\ldots\mathbf{1}X_{m}\mathbf{1}\Omega),
$$

for some $\lambda_1, \ldots, \lambda_m \in C^{\sim}(U)$. Next, using (15), it is easy to show that

$$
\mathcal{L}_{X_m}\Omega = \lambda_m\Omega,
$$

\n
$$
\mathcal{L}_{X_{m-1}}(X_m\mathbf{1}\Omega) = \lambda_{m-1}(X_m\mathbf{1}\Omega),
$$

\n
$$
\vdots
$$

\n
$$
\mathcal{L}_{X_1}(X_2\mathbf{1} \dots \mathbf{1}X_m\mathbf{1}\Omega) = \lambda_1(X_2\mathbf{1} \dots \mathbf{1}X_m\mathbf{1}\Omega).
$$

Then we may apply Incorem 3.15 to give us that $\{\omega^*, \ldots, \omega^m\}$ is qual to $\{A_1,\ldots,A_m\}$, and that for all ω up to $i=m$,

$$
\omega^1 = d\gamma^1,
$$

\n
$$
\omega^2 = d\gamma^2 - X_1(\gamma^2)d\gamma^1,
$$

\n
$$
\omega^3 = d\gamma^3 - X_2(\gamma^3)(d\gamma^2 - X_1(\gamma^2)d\gamma^1) - X_1(\gamma^3)d\gamma^1
$$

\n:
\n:
\n
$$
\omega^m = d\gamma^m \mod d\gamma^1, \dots, d\gamma^{m-1},
$$

for some functionally independent $\gamma^*, \ldots, \gamma^* \in C^+(\mathcal{U})$. Now since we know already that $d\omega^1 = 0$, we only have to show that for each $1 < j < l$,

$$
d\omega^j = d\left(\frac{X_1 \mathbf{1} \dots \mathbf{1} X_{j-1} \mathbf{1} X_{j+1} \dots \mathbf{1} X_m \mathbf{1} \Omega}{X_j \mathbf{1} X_1 \mathbf{1} \dots \mathbf{1} X_{j-1} \mathbf{1} X_{j+1} \mathbf{1} \dots \mathbf{1} X_m \mathbf{1} \Omega}\right) = 0.
$$
 (16)

The original symmetry relations for X_1, \ldots, X_m tell us that for each j, . - *المحتود المحتويات المحتويات العبد المل*احية المعارفة المحاربين المحاربين المحاربين المحاربين المحاربين

$$
d(X_{1}\mathbf{1} \ldots \mathbf{1} X_{j-1}\mathbf{1} X_{j+1}\mathbf{1} \ldots \mathbf{1} X_m\mathbf{1} \Omega)
$$

= 0 mod $X_{1}\mathbf{1} \ldots \mathbf{1} X_{j-1}\mathbf{1} X_{j+1}\mathbf{1} \ldots \mathbf{1} X_m\mathbf{1} \Omega.$

Finally, using (15), and in similar fashion to the end of the proof of Theorem 3.17, we get that for ea
h j, Xj is ^a non-trivial symmetry of X1 : : : Xj1 \Box . Simple algebraic to the simple algebraic through the simple algebraic μ algebraic μ

Remark. As in Theorem 3.17, it is easy to see that Theorem 3.18 holds for all $m \geq 2$. However, here we can also say that the theorem holds if $l = m$. so (5) because $\vert \cdot \vert$ in this case $\vert \cdot \vert$, \vert is all $\vert \cdot \vert$ in this case $\vert \cdot \vert$ is also for $\vert \cdot \vert$ \sin uation, all ω^+ become exact, which is in accordance with the corollary to Proposition 2 given in $[1]$.

The next se
tion gives a simple appli
ation of some of the ideas presented above.

4 \blacksquare Dinerential forms in $\Lambda_1(\mathbb{R} \longrightarrow \mathbb{R}$

In this section we show that, provided we have enough symmetries, any differential form in Λ and ∞ is can be expressed locally in terms of m functionally independent fun
tions as in the on
lusion of Theorem 3.15. Further details will be given in Theorem 4.3 below, but first, consider the following result:

Lemma 4.1. Let $\Omega \in \Lambda^m(U)$ for some $m < n$ be non-zero, where U is defined as in previous sections (though the requirement that U be convex is not need to suppose the form of the fo

$$
\Omega := \gamma_1 \theta^2 \wedge \theta^3 \wedge \cdots \wedge \theta^{m+1} + \gamma_2 \theta^1 \wedge \theta^3 \wedge \cdots \wedge \theta^{m+1} + \cdots
$$

$$
+ \gamma_{m+1} \theta^1 \wedge \theta^2 \wedge \cdots \wedge \theta^m,
$$

for some unearly independent $\sigma^1, \ldots, \sigma^{m+1} \in \Lambda^1(U)$ and $\gamma_1, \ldots, \gamma_{m+1} \in C^{\infty}(U)$. Then is de
omposable.

Proof. Let $\Omega \in \Lambda^{\infty}(U)$ be as in the theorem. We can write

$$
\Omega = X \mathbf{J} \left(\theta^1 \wedge \cdots \wedge \theta^{m+1} \right),
$$

where

$$
X := \sum_{i=1}^{m+1} (-1)^{i-1} \gamma_i X_i,
$$

for some $A_1,\ldots,A_{m+1} \in \mathfrak{X}(U)$ qual to $\sigma^2,\ldots,\sigma^{m+1}$. Hence from Corollary 3.10 the result follows. \Box

From Lemma 4.1 we obtain the following useful result for m -forms in $(m+1)$ -dimensional spaces also found in [8] by Godbillon. Define W to be some open neighbourhood of ^Rm+1 .

Proposition 4.2. Let $\Omega \in \Lambda^m(W)$. Then Ω is probenius integrable.

Proof. Let $\Omega \in \Lambda^{\infty}(W)$. Lemma 4.1 implies

$$
\Omega = \theta^1 \wedge \cdots \wedge \theta^m,
$$

for some imearly independent $\sigma^*, \ldots, \sigma^* \in \Lambda^*(W)$. Now all is an $(m+1)$ form in Λ^{++} (W), so we may complete $\theta^+,\ldots,\theta^{++}$ to a basis by including some iniearly independent $\varphi \in \Lambda^-(W)$ with the property that

$$
d\Omega = \theta^1 \wedge \cdots \wedge \theta^m \wedge \phi.
$$

Theorem 4.3. Let $\Omega \in \Lambda^m(W)$, where W to be some open, convex neighbourhood of ^Rm+1 . If there exists a solvable stru
ture of m symmetries for $A(\Omega)$ as in Theorem 3.15, then we can compute functions γ ,..., γ \in $C^-(W)$ so that $\Omega = \gamma^* a \gamma^* \wedge \cdots \wedge a \gamma^*$.

Proof. We know from Proposition 4.2 and Proposition 3.6 respectively that is de omposable and that de omposable and that design and that design and the original contract of the social \Box algorithm for finding $\gamma^*, \ldots, \gamma^m$.

 \Box

$\overline{5}$ Some necessary conditions

For an arbitrary form differential form $\Delta \in \Lambda_+$ (U), we use ideas in the previous section to examine some necessary conditions for Δ to be decomposable and $d\Delta = 0$ mod Δ , so that we can apply Theorem 3.15. Of course if $m = n$, these two onditions trivially hold, and Proposition 4.2 and means they still hold if $m = n - 1$. In this section we examine the situation when $m < n - 1$. In what follows, we assume U is some open, convex neighbourhood of \mathbb{R}^n .

Theorem 5.1. Let $\Delta \in \Lambda^m(U)$ for some $m \leq n-1$. If there exist $n-m-1$ linearly independent vector fields $\Gamma_1, \ldots \Gamma_{n-m-1} \in \mathfrak{X}(U)$ in ker (Δ) , then Δ is decomposable. Moreover, if for each $1 \leq i \leq n - m - 1$,

$$
\mathcal{L}_{\Gamma_i} \Delta = \lambda_i \Delta, \tag{17}
$$

for some $\lambda_i \in C^{\infty}(U)$, then $a\Delta = 0$ mod Δ .

Proof. Let $\Delta \in \Lambda^m(U)$ with $m \leq n-1$, and let there exist intearly independent $\Gamma_1, \ldots \Gamma_{n-m-1} \in \mathfrak{X}(U)$ such that for all $1 \leq i \leq n-m-1$,

$$
\Gamma_i \mathbf{1} \Delta = 0. \tag{18}
$$

Now

$$
(sp\left\{\Gamma_{1},\ldots,\Gamma_{n-m-1}\right\})^{\perp}=sp\left\{\theta^{1},\ldots,\theta^{m+1}\right\},\,
$$

for some $\sigma^*, \ldots, \sigma^{n+1} \in \Lambda^*(U)$. Hence from (18), we must have

$$
\Delta = \Delta_{j_1...j_m} \theta^{j_1} \wedge \cdots \wedge \theta^{j_m},
$$

for some $\Delta_{j_1...j_m} \in C$ (*U*), with summation over $1 \leq j_1 < \cdots < j_m \leq m+1$. Therefore by Lemma 4.1, Δ is decomposable.

For the second part of the proof, we choose without loss,

 $\overline{}$ \wedge \cdots \wedge $\overline{}$

We can complete $\theta^-, \ldots, \theta^{n+1}$ to a basis for $\Lambda^-(U)$ by adding linearly independent $\varphi^*, \ldots, \varphi^*$ $\pi^- \in \Lambda^*(U)$ such that

$$
\{\phi^1, \dots, \phi^{n-m-1}, \theta^1, \dots, \theta^{m+1}\}\tag{19}
$$

is dual to

$$
\{\Gamma_1, \ldots \Gamma_{n-m-1}, Y_1, \ldots, Y_{m+1}\},\tag{20}
$$

for some linearly independent $Y_1, \ldots, Y_{m+1} \in \mathfrak{X}(U)$. Now with summation on k over $1 \leq k \leq m$, we can write

$$
d\Delta = \sigma_k \wedge \theta^1 \wedge \cdots \wedge \theta^{k-1} \wedge \theta^{k+1} \wedge \cdots \wedge \theta^m + \beta \wedge \Delta, \tag{21}
$$

for some $\sigma_1, \ldots, \sigma_m \in \Lambda^-(U)$ and $\rho \in \Lambda^-(U)$ with the property that each σ_k only depends on the basis vectors $\varphi^*, \ldots, \varphi^{n+m-1}$, φ^{n+m-1} . Hence from the dual basis property in (19) and (20), we have for each k ,

$$
Y_j \mathbf{1} \sigma_k = 0,\tag{22}
$$

for all $1 \leq j \leq m$. By combining the assumptions in (17) and (18), we have for all i ,

$$
\Gamma_{i\perp}d\Delta = \lambda_i \Delta. \tag{23}
$$

Using the dual basis property once more, we get that for each i and $1 \leq l \leq$ $m+1$, $1/7 = 0$. Hence substituting (21) into (25) gives (with sum),

$$
(\Gamma_{i}\mathbf{1}\sigma_{k})\wedge\theta^{1}\wedge\cdots\wedge\theta^{k-1}\wedge\theta^{k+1}\wedge\cdots\wedge\theta^{m}+(\Gamma_{i}\mathbf{1}\beta)\wedge\Delta=\lambda_{i}\Delta, \qquad(24)
$$

for each i . Since each $i_j \cup \sigma_k$ only depends on the basis vectors φ , . . . , φ , σ \cdots , for (24) to noid we must have

$$
\Gamma_i \mathbf{1} \sigma_k = 0,\tag{25}
$$

for each i and k. Hence from (22) and (25), $\ker(\sigma_k)$ is at least $(n-1)$ dimensional. This means $\mathbb{P}(X)=\mathbb{$ for each k. Therefore $d\Delta = \beta \wedge \Delta$. \Box

Theorem 5.1 has the following two corollaries:

Corollary 5.2. Let $\Delta \in \Lambda^m(U)$ such that $m \leq n-1$. If there exist $n-m$ 1 linearly independent Cauchy characteristic vector fields of the differential ideal $\langle \Delta, d\Delta \rangle$, then Δ is decomposable and $d\Delta = 0 \mod \Delta$.

Proof. Since the Cauchy characteristic vector fields are in the kernel of Δ . Theorem 5.1 implies Δ is decomposable. Now it is clear that (17) in Theorem 5.1 still notas for some $\lambda_1, \ldots, \lambda_{n-m-1} \in C^{\pm}(U)$. Hence from the theorem, $d\Delta = 0 \text{ mod } \Delta$ \Box

Corollary 5.3. Let $\Delta \in \Lambda^m(U)$ such that $m \leq n-1$. If there exist $n-m$ 1 linearly independent Cauchy characteristic vector fields of the differential ideal $\langle \Delta, d\Delta \rangle$, then the Cauchy characteristic space of $\langle \Delta, d\Delta \rangle$ is $(n - m)$. dimensional

Proof. From Corollary 5.2, Δ is decomposable, so ker(Δ) is $(n-m)$ -dimensional. The corollary also means Δ is closed modulo itself which implies $\langle \Delta \rangle$ = $\langle \Delta, d\Delta \rangle$, and hence their Cauchy characteristic spaces are equal. From Proposition 2.7 the result follows. \Box

Now the dimension of the Cauchy characteristic space of $\langle \Delta, d\Delta \rangle$ is always less than or equal to that of ker(Δ), and the maximum dimensional of ker(Δ) is $n - m$, which occurs when Δ is decomposable. Theorem 5.1 therefore

means that if ker(Δ) is at least $(n - m - 1)$ -dimensional, then it is $(n - m)$ dimensional. Similarly, Corollary 5.3 means that if the Cauchy characteristic space of $\langle \Delta, d\Delta \rangle$ is at least $(n - m - 1)$ -dimensional, then it is $(n - m)$ dimensional.

Next, we illustrate Corollary 5.2 with the following example:

Example 5.4. Suppose U^{\dagger} is some suitably chosen open, convex neighbourhood of \mathbb{R}^+ with coordinates $x^-, x^-, x^-, x^-,$ and

$$
\Delta := \frac{2x^2x^4}{x^3} dx^3 \wedge dx^2 - \left(\frac{x^4}{x^3}\right) dx^3 \wedge dx^1 - 2dx^4 \wedge dx^1 + \frac{1}{x^3x^4} dx^1 \wedge dx^2 + 4x^2 dx^4 \wedge dx^2.
$$

Now the vector field

$$
\Gamma := 4x^2 \frac{\partial}{\partial x^1} + 2 \frac{\partial}{\partial x^2} - \frac{1}{x^3 x^4} \frac{\partial}{\partial x^2},
$$

is a Cauchy characteristic of $\langle \Delta, d\Delta \rangle$. Hence from Corollary 5.2, Δ is decomposable and $d\Delta = 0$ mod Δ . Note from Corollary 5.3 that the Cauchy characteristic space of $\langle \Delta, d\Delta \rangle$ is two-dimensional.

We will now proceed to apply Theorem 3.15 to Δ . It is easy to see that $\overline{\partial x^1}$ is a non-trivial symmetry of Δ . With

$$
\frac{\partial}{\partial x^1} \Delta = \frac{1}{x^3 x^4} dx^2 + \frac{x^4}{x^3} dx^3 + 2dx^4,
$$

it is also easy to see that $\frac{\partial}{\partial x^2}$ is a non-trivial symmetry of $\frac{\partial}{\partial x^1}\mathsf{J}\,\Delta$. Now from Theorem 3.15 and Corollary 3.13,

$$
\omega^1 := \frac{\frac{\partial}{\partial x^1} \mathbf{1} \Delta}{\frac{\partial}{\partial x^2} \mathbf{1} \frac{\partial}{\partial x^1} \mathbf{1} \Delta} = dx^2 + (x^4)^2 dx^3 + 2x^3 x^4 dx^4 = d\left(x^2 + x^3 (x^4)^2\right).
$$

Also, it is not hard to show that

$$
\omega^{2} := \frac{\frac{\partial}{\partial x^{2}} \Delta}{\frac{\partial}{\partial x^{1}} \Delta \frac{\partial}{\partial x^{2}} \Delta} = dx^{1} + 2x^{2}(x^{4})^{2} dx^{3} + 4x^{2} x^{3} x^{4} dx^{4},
$$

$$
= d(x^{1} - (x^{2})^{2}) + 2x^{2} d(x^{2} + x^{3}(x^{4})^{2}).
$$

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e

$$
\Delta = \frac{1}{x^3 x^4} d \left(x^1 - (x^2)^2 \right) \wedge d \left(x^2 + x^3 (x^4)^2 \right).
$$

6Pfaffian equations

In this section we examine how symmetries may be used to express a differential one-form `normal form' given in (26). We begin with the following definition and theorem:

Definition 6.1. Let $\alpha \in \Lambda^1(U)$. The rankof the Pfaffian equation $\alpha = 0$ at the point $p \in U$ is the non-negative integer r such that $(d\alpha) \wedge \alpha \neq 0$ and α α α = 0 at p.

If a one-form α is exact, i.e. $\alpha = a$ for some $\gamma \in C^{\infty}(U)$, then it (and any linearly dependent one-form) has rank zero.

Theorem 0.2. Let $\alpha \in \Lambda$ (U) and suppose the equation $\alpha = 0$ is of constant Tank T on U . Then there exists a coordinate system $\gamma^*, \ldots, \gamma^* \in C^{\infty}(U)$, where $2r + 1 \leq n$, so that the equation becomes

$$
d\gamma^1 + \gamma^2 d\gamma^3 + \cdots + \gamma^{2r} d\gamma^{2r+1} = 0.
$$

Theorem 6.2 is known as the Pfaff problem. A proof of this theorem may be found in $[2]$.

It is easy to see that multiplying any one-form of constant rank on U by a nowhere zero smooth function f leaves the rank unchanged, using the fact that for any $m \in \mathbb{N}$, we have $(d(f\alpha))^m \wedge (f\alpha) = f^{m+1}(d\alpha)^m \wedge \alpha$. This allows us to express any $\alpha \in \Lambda$ (U) of constant rank r on U as

$$
\alpha = \gamma^0 (d\gamma^1 + \gamma^2 d\gamma^3 + \dots + \gamma^{2r} d\gamma^{2r+1}), \qquad (26)
$$

for some γ , \ldots , γ ⁻¹ \in C \lceil (U).

Theorem 6.3. Let $\alpha \in \Lambda^2(U)$. Suppose α is of constant rank r on U. Define $\Omega := (d\alpha) \wedge \alpha$. Then Ω is decomposable and $d\Omega = 0$ mod Ω .

Proof. Let $\alpha \in \Lambda^-(U)$ with α of constant rank r on U. Hence

$$
\alpha = \gamma^0 (d\gamma^1 + \gamma^2 d\gamma^3 + \dots + \gamma^{2r} d\gamma^{2r+1}),
$$

for some γ , \ldots , γ \cdots \in C \cdots (U). Define

$$
\overline{\alpha} := d\gamma^1 + \gamma^2 d\gamma^3 + \cdots + \gamma^{2r} d\gamma^{2r+1}.
$$

Further, let $\Omega := (d\alpha)^{\top} \wedge \alpha$. We will first show that $d\Omega = 0$. Simple computation yields

$$
(d\overline{\alpha})^r = r!d\gamma^2 \wedge \cdots \wedge d\gamma^{2r+1}.
$$

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e

$$
\overline{\Omega} = \overline{\alpha} \wedge (d\overline{\alpha})^r,
$$

= $r! d\gamma^1 \wedge d\gamma^2 \wedge \cdots \wedge d\gamma^{2r+1}$

$$
\Omega = (\gamma^0)^{r+1} (d\overline{\alpha})^r \wedge \overline{\alpha}.
$$

e die die ook die ook die die die die begin die die begin die ook die getal op die getal die die getal op die

$$
d\Omega = d((\gamma^0)^{r+1}) \wedge (d\overline{\alpha})^r \wedge \overline{\alpha}.
$$

But, $(d(\gamma^0 \overline{\alpha})) \wedge ((\gamma^0) \overline{\alpha}) = (\gamma^0)^{r+1}(d\overline{\alpha})^r \wedge \overline{\alpha}$. Hence $d\Omega = 0$ mod Ω as γ^0 is nowhere zero on U . Finally, since Ω is decomposable and $\Omega = (\gamma^2)^{-1}$ is therefore de
omposable. \Box

Our aim is to use Theorem 6.3 with Theorem 3.15 to ultimately find some coordinates for the Pfaff problem in Theorem 6.2. The next theorem illustrates how this may be done for one-forms that are of onstant rank one on U , which will be later extended to one-forms of any constant rank $r > 1$. The case $r = 0$ involves a trivial application of Theorem 3.15, and will therefore be ignored.

To assist in finding coordinates for the Pfaff problem, the following lemma will be needed:

Lemma 6.4. Let $\alpha \in \Lambda^2(U)$ and suppose α is of constant non-zero rank r on U. Let $\Omega := (d\alpha) \land \alpha$ and $X \in \mathfrak{X}(U)$ such that $X \cup \Omega = 0$. Then $X \cup \alpha = 0$.

Proof. Let $\alpha \in \Lambda$ (*U*). Suppose α is of constant non-zero rank r on U, and decrease the second communities are the sequence of the sequence of the second α

$$
0 = X \lrcorner \Omega = (X \lrcorner (d\alpha)^r) \wedge \alpha + (X \lrcorner \alpha)(d\alpha)^r
$$

By taking the exterior product with α , we obtain

$$
(X \lrcorner \alpha)(d\alpha)^r \wedge \alpha = 0.
$$

Since α is of rank r, (a α) \land $\alpha \neq 0$, and hence $\Lambda \lrcorner \alpha = 0$.

Theorem 6.5. Let $\alpha \in \Lambda^1(U)$ such that α is of constant rank one on U. er een die die die tot die gewone begin die die die die die die die die gewone die gewone bewerk die die begin $X_1, X_2, X_3 \in \mathfrak{X}(U)$ is a solvable structure of linearly independent symmetries sure that the symmetry of the extra symmetry \boldsymbol{y} is a non-trivial symmetry of the extra symmetry that X3 ⁼ 0, X2 is ^a non-trivial symmetry of A(h i)spfX3g, and X1 is ^a non-trivial symmetry of $A(\{Y\}) \oplus sp\{X_2, X_3\}$. Then with $\omega^2, \omega^2, \omega^2 \in \Lambda^2(U)$ defined by

$$
\omega^{1} := \frac{X_{2} \Box X_{3} \Box \Omega}{X_{1} \Box X_{2} \Box X_{3} \Box \Omega},
$$

$$
\omega^{2} := \frac{X_{1} \Box X_{3} \Box \Omega}{X_{2} \Box X_{1} \Box X_{3} \Box \Omega},
$$

$$
\omega^{3} := \frac{X_{1} \Box X_{2} \Box \Omega}{X_{3} \Box X_{1} \Box X_{1} \Box \Omega},
$$

we have

$$
\omega^1 = d\gamma^1,
$$

\n
$$
\omega^2 = d\gamma^2 - X_1(\gamma^2)d\gamma^1,
$$

\n
$$
\omega^3 = d\gamma^3 - X_2(\gamma^3)(d\gamma^2 - X_1(\gamma^2)d\gamma^1) - X_1(\gamma^3)d\gamma^1,
$$

for some functionally independent $\gamma^*, \gamma^*, \gamma^* \in C^{\infty}(U)$, and

$$
\alpha = (X_2 \mathbf{1} \alpha) \left(d\gamma^2 + \frac{(X_1 \mathbf{1} \alpha) - (X_2 \mathbf{1} \alpha) X_1(\gamma^2)}{(X_2 \mathbf{1} \alpha)} d\gamma^1 \right). \tag{27}
$$

 \Box

. ..., with a set of the state and the composition of the set of th $ax = 0$ mod st. Theorem 5.15 can be used to obtain $\{\omega^*, \omega^*, \omega^*\}$ qual to ${X_1, X_2, X_3}$, where

$$
\omega^1 = d\gamma^1,
$$

\n
$$
\omega^2 = d\gamma^2 - X_1(\gamma^2)d\gamma^1,
$$

\n
$$
\omega^3 = d\gamma^3 - X_2(\gamma^3)(d\gamma^2 - X_1(\gamma^2)d\gamma^1) - X_1(\gamma^3)d\gamma^1
$$

for some functionally independent $\gamma^*, \gamma^*, \gamma^* \in C^{\pm}(\mathcal{U})$. Now from Lemma 6.4, in all X and with the Ward (1-1). And since the α and α are left with α

$$
\alpha = (X_1 \mathsf{J} \alpha)\omega^1 + (X_2 \mathsf{J} \alpha)\omega^2.
$$

Now $X_2 \rvert \alpha \neq 0$ in the neighbourhood, since α is nowhere rank zero by assumption. Hen
e

$$
\alpha = (X_2 \mathbf{1} \alpha) \left(d\gamma^2 + \frac{(X_1 \mathbf{1} \alpha) - (X_2 \mathbf{1} \alpha) X_1(\gamma^2)}{(X_2 \mathbf{1} \alpha)} d\gamma^1 \right).
$$

Remark 1. The extra condition in Theorem 6.5 that the non-trivial symmetry X_3 satisfies $X_3 \rvert \alpha = 0$ implies from Proposition 2.6 that the symmetry is not a Cauchy characteristic vector field of $\langle \alpha, d\alpha \rangle$. Therefore $X_3 \perp d\alpha$ is not some multiple of α (as α is of rank one, it is impossible that $d\alpha = 0$ mod α). Such a symmetry exists since if γ ,..., γ are coordinates for U and $\alpha := \gamma^+(u\gamma^+ + \gamma^- u\gamma^+)$ is already in normal form for some $\gamma^+ \in C^+(U)$, then it is easy to show that Theorem 6.5 can be applied to such α with X_3 chosen as $\frac{1}{\partial \gamma^2}$ or $\frac{1}{\partial \gamma^3}$ – $\gamma^2 \frac{1}{\partial \gamma^1}$.

Remark. 2 In deriving our expression for α in (27), we do not need to calcurate γ . This significantly reduces the number of algebraic manipulations required.

We illustrate Theorem 6.5 with an example:

Example 6.6. Suppose we are in some open, convex heighbourhood of K³, denoted by U^* , with coordinates x^*, x^*, x^* . Define on some suitably chosen U^{\pm} ,

$$
\alpha := -\frac{x^2 x^3}{(x^1)^2} dx^1 + \left(\frac{x^1}{x^2} + \frac{x^3}{x^1}\right) dx^2 + \frac{x^1}{x^3} dx^3.
$$

By dimension, $(d\alpha)^T \wedge \alpha = 0$, and it is easy to show that $d\alpha \wedge \alpha \neq 0$ on some region of U^* . Suppose U^* is chosen such that $a\alpha \wedge \alpha \neq 0$ everywhere. Since *any* non-zero vector held is a non-trivial symmetry of $a\alpha \wedge \alpha \in \Lambda^*(U^*)$, we may choose any X_3 such that $X_3 \Box \alpha = 0$. So let

$$
X_3 := \frac{x^2 x^3}{(x^1)^2} \frac{\partial}{\partial x^3} + \frac{x^1}{x^3} \frac{\partial}{\partial x^1}
$$

 \Box

be the symmetry. Now

$$
X_2:=(x^3)^2\frac{\partial}{\partial x^3}
$$

is a non-trivial symmetry of $sp{X_3}$ $(A(\langle d\alpha \wedge \alpha \rangle)$ is zero-dimensional), and by inspection that

$$
X_1 := \frac{\partial}{\partial x^2}
$$

is a non-trivial symmetry of $sp{X_2, X_3}$. These yield

$$
\omega^1 := \frac{X_2 \mathbf{J} X_3 \mathbf{J} (d\alpha \wedge \alpha)}{X_1 \mathbf{J} X_2 \mathbf{J} X_3 \mathbf{J} (d\alpha \wedge \alpha)} = dx^2,
$$

and

$$
\omega^2 := \frac{X_{1\perp} X_{3\perp} (d\alpha \wedge \alpha)}{X_{2\perp} X_{1\perp} X_{3\perp} (d\alpha \wedge \alpha)} = -\frac{x^2}{(x^1)^3} dx^1 + \frac{dx^3}{(x^3)^2},
$$

$$
= d\left(\frac{x^2}{2(x^1)^2} - \frac{1}{x^3}\right) - \frac{1}{2(x^1)^2} dx^2
$$

Hence a simple calculation gives

$$
\alpha = x^1 x^3 \left(d \left(\frac{x^2}{2(x^1)^2} - \frac{1}{x^3} \right) + \left(\frac{1}{x^2 x^3} + \frac{1}{2(x^1)^2} \right) dx^2 \right).
$$

Such expressions for α are in general not unique, and may be found by choosing different symmetries. For example, we have also obtained

$$
\alpha = x^3 \left(d \left(\frac{x^2}{x^1} \right) + \frac{x^1}{x^3} d \left(\ln \left| x^2 x^3 \right| \right) \right).
$$

We now present a generalisation of Theorem 6.5:

Theorem 6.7. Let $\alpha \in \Lambda^1(U)$ have constant rank r on U, and define $\Omega := (d\alpha)^r \wedge \alpha$. Let $X_1, \ldots, X_{2r+1} \in \mathfrak{X}(U)$ be a solvable structure of linearly independent symmetries such that X_{2r+1} is a non-trivial symmetry of $A(\langle \Omega \rangle)$, and for each $1 \langle i \rangle \langle 2r+1, X_i \rangle$ is a non-trivial symmetry of $A(\langle \Omega \rangle) \oplus \{X_{i+1}, \ldots, X_{2r+1}\}.$ Suppose, in addition, that for the r vector fields X_{r+2},\ldots, X_{2r+1} , we have $X_{r+2}\square \alpha = 0,\ldots, X_{2r+1}\square \alpha = 0$. For all $1 \leq i \leq 2r+1$, define ω^i by

$$
\omega^i := \frac{X_{1\mathbf{J}} \dots \mathbf{J} X_{i-1\mathbf{J}} X_{i+1\mathbf{J}} \dots \mathbf{J} X_{2r+1\mathbf{J}} \Omega}{X_{i\mathbf{J}} X_{1\mathbf{J}} \dots \mathbf{J} X_{i-1\mathbf{J}} X_{i+1\mathbf{J}} \dots \mathbf{J} X_{2r+1\mathbf{J}} \Omega}
$$

Then for all ω^i up to $i = 2r + 1$,

$$
\omega^1 = d\gamma^1,
$$

\n
$$
\omega^2 = d\gamma^2 - X_1(\gamma^2)d\gamma^1,
$$

\n
$$
\omega^3 = d\gamma^3 - X_2(\gamma^3)(d\gamma^2 - X_1(\gamma^2)d\gamma^1) - X_1(\gamma^3)d\gamma^1,
$$

\n
$$
\vdots
$$

\n
$$
\omega^{2r+1} = d\gamma^{2r+1} \mod d\gamma^1, \dots, d\gamma^{2r},
$$

for some functionally independent $\gamma^*, \ldots, \gamma^{++} \in C^{++}(U)$, and $\alpha = (\Lambda_1 \Box \alpha) a \gamma^+ + (\Lambda_2 \Box \alpha) (a \gamma^- - \Lambda_1 (\gamma^-) a \gamma^+)$ $+ (\Lambda_3 \Box \alpha)(a\gamma) - \Lambda_2(\gamma)$ $(a\gamma) - \Lambda_1(\gamma) a\gamma$ $) - \Lambda_1(\gamma) a\gamma$ $) + \ldots$ $+ (\Lambda_{r+1} \Box \alpha)(a\gamma)^{-1} = \Lambda_r(\gamma^{r+1})(a\gamma)^{-1} = \cdots = \Lambda_1(\gamma^{r+1})a\gamma^{-1} = \ldots$ $-\Lambda_1(\gamma + \mu)\alpha\gamma$,

which when rearranged give α in the form of (26).

Proof. The proof follows in a similar fashion to Theorem 6.5. The conditions $X_{r+2}\mathbf{1}\alpha = 0,\ldots,X_{2r+1}\mathbf{1}\alpha = 0$ and Lemma 6.4 ensure that α is a imear combination of $a\gamma$,..., $a\gamma$. Further, since α is of constant rank r, $X_{r+1}\mathbf{I} \alpha \neq 0$, so we are permitted to divide by it, and hence express α in the form of (26) . \Box

Remark. Both remarks for Theorem 6.5 may be extended to Theorem 6.7 as follows: Firstly, from the proof of Theorem 6.3 it is clear that there exist r non-trivial symmetries A_{r+2},\ldots,A_{2r+1} of ($a\alpha$)' $\wedge\alpha$ in ker(α), and secondly, in deriving our expression for α , we do not need to calculate any γ . \cdot , \ldots , γ . \cdot .

7Darboux systems

This section gives an algorithm based on vector fields for generating a set of oordinates in Darboux's theorem given below in Theorem 7.4. To begin with, we present some preliminary material. In Bryant *et al.* [2] there is the following fundamental theorem:

Theorem 7.1. Let $\Omega \in \Lambda^2(U)$ and let r be the natural number such that Δt \neq 0 and Δt \cdot \cdot \equiv 0. Then there exist 2r unearly independent elements $\omega^-, \ldots, \omega^- \in \Lambda^-(U)$ such that

$$
\Omega = \omega^1 \wedge \omega^2 + \cdots + \omega^{2r-1} \wedge \omega^{2r}.
$$

In what follows, we will also make use of the following lemma:

Lemma 7.2. Let $\Omega \in \Lambda^2(U)$ and $T \in \mathbb{N}$ such that $\Omega^2 \neq 0$ and $\Omega^2 = 0$. Also let X ² X(U). Then X r = 0 if and only if X = 0.

Proof. Let $\Omega \in \Lambda^-(U)$ with $\Lambda \square \Omega^c = 0$ for some vector field $\Lambda \in \mathcal{K}(U)$. Then from Theorem 7.1 we have

$$
\Omega = \omega^1 \wedge \omega^2 + \dots + \omega^{2r-1} \wedge \omega^{2r},\tag{28}
$$

for some imearly independent $\omega^-, \ldots, \omega^- \in \Lambda^-(U)$. This implies

$$
\Omega^r = r! \omega^1 \wedge \cdots \wedge \omega^{2r}.
$$

Now $A \cup \Omega = 0$ implies that $A \cup \omega = 0$ for all $1 \leq i \leq 2r$. Hence using the expression for the converse is obvious for the converse in the converse is obvious for the converse in the con \Box since if Y is any vector field in $\mathfrak{X}(U)$, then Y $\exists M \equiv T(Y \sqcup M) \wedge M$.

Theorem 7.3. Let $\Omega \in \Lambda^2(U)$ be closed. Suppose r is the natural number such that $\alpha \neq 0$ and $\alpha \rightarrow \infty$. Further suppose we have a solvable structure of $2r$ linearly independent symmetries $X_1, \ldots, X_{2r} \in \mathfrak{X}(U)$ such that X_{2r} is a non-trivial symmetry of $A(\N^T)$, and for all $1 \leq i \leq 2r$, Λ_i is a non-trivial symmetry of $A(\lambda t) \oplus sp_{1}A_{i+1}, \ldots, A_{2r}$. Then Theorem 3.15 gives us an algorithm for expressing as existy in terms of the 2r functional ly independent functions γ ,..., $\gamma \in C$ (U) and their exterior derivatives

Proof. Let $\Omega \in \Lambda^1(U)$ be closed with $\Omega \neq 0$ and $\Omega^* \equiv 0$ for some $r \in \mathbb{N}$. Since $a_{32} = 0$ implies that a_{32} $\rangle = 0$, from Proposition 2.7, $\ker(M_{\perp}) = A(\langle M \rangle)$ is therefore Frobenius integrable. The fact that M is decomposable of degree 2r means that $A(\Omega)$ is generated by $n=2r$ linearly independent vector fields. Suppose we have a set of linearly independent symmetries $X_1, \ldots, X_{2r} \in \mathfrak{X}(U)$ such that X_{2r} is a non-trivial symmetry of $A(\Omega)$, and for all $1 \leq i \leq 2r$, A_i is a non-trivial symmetry of $A(\Omega)$ by $sp_1\Lambda_{i+1}, \ldots, \Lambda_{2r}$. Then by Theorem 3.15 we have on U, $\{\omega^-, \ldots, \omega^-\}$ dual to $\{\Lambda_1, \ldots, \Lambda_{2r}\}$, where for all $1 \leq j \leq 2r$,

$$
\omega^j := \frac{X_{1\mathbf{J}} \dots \mathbf{J} X_{j-1} \mathbf{J} X_{j+1} \mathbf{J} \dots \mathbf{J} X_{2r} \mathbf{J} \Omega^r}{X_{j\mathbf{J}} X_{1\mathbf{J}} \dots \mathbf{J} X_{j-1} \mathbf{J} X_{j+1} \mathbf{J} \dots \mathbf{J} X_{2r} \mathbf{J} \Omega^r},
$$

and

$$
\omega^1 = d\gamma^1,
$$

\n
$$
\omega^2 = d\gamma^2 - X_1(\gamma^2)d\gamma^1,
$$

\n
$$
\omega^3 = d\gamma^3 - X_2(\gamma^3)(d\gamma^2 - X_1(\gamma^2)d\gamma^1) - X_1(\gamma^3)d\gamma^1,
$$

\n
$$
\vdots
$$

\n
$$
\omega^{2r} = d\gamma^{2r} \mod d\gamma^1, \dots, d\gamma^{2r-1},
$$

for some functionally independent $\gamma^*, \ldots, \gamma^* \in C^{\{t\}}$). Then by Lemma 7.2, and using the fact that $\{X_1,\ldots,X_{2r}\}$ plus any set of generators of $A(\langle\Omega\rangle)$ spans $\mathfrak{X}(U)$, we can therefore write

$$
\Omega = \Omega(X_k, X_l)\omega^k \wedge \omega^l, \qquad 1 \le k < l \le 2r,
$$

where we are implying a double summation. This means that

$$
\Omega = \Omega_{kl} d\gamma^k \wedge d\gamma^l, \qquad 1 \le k < l \le 2r,\tag{29}
$$

for some functions $\Omega_{kl} \in C^{\infty}(U)$. But since Ω is closed, we must have for all $1 \in A(\Omega)$),

$$
\mathcal{L}_{\Gamma}\Omega = d(\Gamma \Box \Omega) = 0,
$$

also using Lemma 7.2. Since $\Gamma(\gamma) = 0$ for all i , it follows that (with sum)

$$
0 = \mathcal{L}_{\Gamma} \Omega = \Gamma(\Omega_{kl}) d\gamma^k \wedge d\gamma^l.
$$

 λ and l. Hence the 2r for each control on the 2r for eac \Box functions $\gamma^*, \ldots, \gamma^-$ and their exterior derivatives.

Remark. In applying Theorem 7.3, there will exist situations when it may be difficult to express each Ω_{kl} in terms of the known $\gamma^2, \ldots, \gamma^2$.

Next, consider Darboux's theorem proved in $[2, 5]$:

Theorem 7.4. (Darboux) Let $\Omega \in \Lambda^2(U)$ be closed so that $\Omega \neq 0$ and Ω Γ = 0 for some $r \in \mathbb{N}$. Then there exist coordinates γ ,..., γ such that

$$
\Omega = d\gamma^1 \wedge d\gamma^2 + \cdots + d\gamma^{2r-1} \wedge d\gamma^{2r}.
$$

Theorem 7.3 may be applied to Darboux's theorem; however, the difficulty is that Theorem 7.3 expresses Ω in terms of a sum of a maximum of $\binom{2r}{0}$ form omponents, whi
h must then be simplied to r omponents with unit one coefficients if we wish to find a set of coordinates in Darboux's theorem.

As an alternative approach extending work in [5] by Crampin and Pirani in their proof of Darboux's theorem (though similar proofs an be found in the literature), we now look to formulate an extraction process for generating a set of oordinates in the theorem using solvable symmetry stru
tures. The next three theorems will be useful in establishing this.

Theorem 7.5. Let $\Omega \in \Lambda^2(U)$ with $\Omega^2 \neq 0$ and $\Omega^2 = 0$ for some $T \geq 2$. Suppose there exist X1; X2 ² X(U) su
h that (X1; X2) ⁼ ¹ and (X1) ^ $(X_2\cup\Omega)\neq 0$. If Ω is defined by $\Omega := \Omega + (X_2\cup\Omega) \wedge (X_1\cup\Omega)$, then Ω $\neq 0$ and $\Omega = 0$.

Proof. Let $\Omega \in \Lambda^-(U)$ such that $\Omega \neq 0$ and $\Omega^{-+} = 0$ for some $r \geq 2$. Using t . The definition for the theorem gives \mathbf{q}

$$
\overline{\Omega}^r = \Omega^r + r\Omega^{r-1} \wedge (X_2 \cup \Omega) \wedge (X_1 \cup \Omega). \tag{30}
$$

 $\mathcal{N} = \mathcal{N}$ and $\mathcal{N} = \mathcal{N}$ and $\mathcal{N} = \mathcal{N}$ we have the set of $\mathcal{N} = \mathcal{N}$

$$
\Omega^r = \Omega^r(X_2 \cup X_1 \cup \Omega),
$$

= $X_2 \cup (\Omega^r \wedge (X_1 \cup \Omega)) - (X_2 \cup \Omega^r) \wedge (X_1 \cup \Omega),$
= $X_2 \cup (\Omega^r \wedge (X_1 \cup \Omega)) - (r(X_2 \cup \Omega) \wedge \Omega^{r-1}) \wedge (X_1 \cup \Omega).$ (31)

In the second line we have used the property A_2 (Ω) Λ (A_1) Ω) = (A_2) Ω) Λ $(\Lambda_1 \mathsf{J} \mathcal{U}) + (\Lambda_2 \mathsf{J} \Lambda_1 \mathsf{J} \mathcal{U}) \mathcal{U}$, and in the third, we have expanded $\Lambda_2 \mathsf{J} \mathcal{U}$. If we substitute the end result in (31) into the expression for Ω in (30), we obtain

$$
\overline{\Omega}^r = X_2 \mathbf{J} \left(\Omega^r \wedge (X_1 \mathbf{J} \Omega) \right). \tag{32}
$$

Dy Incorem 7.1, there exist imearly independent one-forms $\omega^*, \ldots, \omega^- \in$ Λ (U) such that

$$
\Omega = \omega^1 \wedge \omega^2 + \cdots + \omega^{2r-1} \wedge \omega^{2r}.
$$

Hence $A_1 \mathbf{J} u = a_1 \omega^2 + \cdots + a_{2r} \omega^{2r}$ for some $a_1, \ldots, a_{2r} \in C^{\infty}(U)$. Since

$$
\Omega^r = r! \omega^1 \wedge \cdots \wedge \omega^{2r}
$$

it follows that $\Omega^r \wedge (X_1 \mathbf{1} \Omega) = 0$. Thus from (32) we get $\Omega^r = 0$. Now suppose $\overline{\Omega}^{r-1} = 0$. Then

$$
0 = \overline{\Omega}^{r-1} = \Omega^{r-1} + (r-1)\Omega^{r-2} \wedge (X_2 \Box \Omega) \wedge (X_1 \Box \Omega).
$$

This implies

now support the support of the support of

$$
\Omega^{r-1} = (r-1)\Omega^{r-2} \wedge (X_1 \mathbf{1} \Omega) \wedge (X_2 \mathbf{1} \Omega). \tag{33}
$$

Taking the exterior produ
t with gives

$$
\Omega^r = (r-1)\Omega^{r-1} \wedge (X_1 \mathbf{J} \Omega) \wedge (X_2 \mathbf{J} \Omega) = 0,\tag{34}
$$

where the second equality comes from substituting Ω = 1n (34) with its expression in (33). The calculations still holds for $r = 2$, and hence we reach a contradiction for all $r \geq 2$. \Box

 R remark. Although Theorem 7.5 demands that Δ be such that Δ be such that Δ 1, we can relax this condition by saying that all we need is to find two vector elds Y1; Y2 ² X(U) su
h that (Y1; Y2) 6= 0. Then we an hoose X1; X2 as, respectively, section in that we have constructed to the μ substitution of μ

The second theorem we require concerns the foliated exterior derivative, as explained by Vaisman $[15]$:

THEOREM 1.0. Let $\omega \in \Lambda^2(U)$ and $\alpha^2, \ldots, \alpha^r \in \Lambda^2(U)$ be standary independent one-forms such that for all $1 \leq i \leq s$,

$$
d\alpha^i=0\mod \alpha^1,\ldots,\alpha^s
$$

(i.e. the Frobenius condition holds so that $\ker(\alpha^+ \wedge \cdots \wedge \alpha^+)$ is Frobenius integrable).

Then if

 $a\omega = 0 \mod \alpha^{\scriptscriptstyle\top}, \ldots, \alpha^{\scriptscriptstyle\top},$

then

 $\omega = a \cdot \mod \alpha \,$,..., α ,

for some $f \in C^{\infty}(U)$.

Using the foliated exterior derivative, we prove the following theorem:

Theorem 1.1. Let $\Omega \in \Lambda^2(U)$ be closed. If there exists a pair of vector field $X_1, X_2 \in \mathfrak{X}(U)$ such that

$$
1. \mathcal{L}_{X_1} \Omega = 0,
$$

- 2. LX2 ⁼ ⁰ mod X1 ;
- 3. (32 **1 3.**) 7. (32 **1 3.)** 7. (32) 7

then on U,

$$
(X_1 \cup \Omega) \wedge (X_2 \cup \Omega) = df \wedge dg,
$$

for some functionally independent smooth f and q.

Proof. Let $\Omega \in \Lambda$ (U) be closed and let there exist vector helds $\Lambda_1, \Lambda_2 \in$ $\mathcal{L}(\sigma)$ that satisfy the three conditions in the three three three three \mathcal{A}_1 . σ implies α (x1 d + d) and α and property α and α $\frac{1}{1}$, and for some smooth f $\frac{1}{1}$

where \mathbf{r} is the same argument to above, \mathbf{r} is above, \mathbf{r} and \mathbf{r} above, \mathbf{r} are above, \mathbf{r} and \mathbf{r} α and some smooth smooth g1. If, how assumption, LX2 if α , then by assumption by

$$
0 \neq \mathcal{L}_{X_2}\Omega = \alpha \wedge (X_1 \cup \Omega),
$$

for some $\alpha \in \Lambda^*(U)$. Therefore

$$
(\mathcal{L}_{X_2}\Omega) \wedge (X_1 \square \Omega) = 0.
$$

Using LX2 ⁼ X2 ^d ⁺ d(X2) and the fa
t that is losed gives

$$
d(X_{2}\lrcorner \Omega) \wedge (X_{1}\lrcorner \Omega) = 0.
$$

Hen
e

$$
d(X_2 \Box \Omega) = 0 \mod (X_1 \Box \Omega).
$$

Using Theorem 7.6, we then get

$$
X_2 \Box \Omega = dg_2 \mod df,
$$

for some smooth g_2 . Hence in both cases the result is proved.

 \Box

We now present the main result of this section:

Theorem 7.8. Let $\Omega \in \Lambda^-(U)$ be closed with $\Omega \neq 0$ and $\Omega^{-+}=0$ for some $r \in \mathbb{N}$. Then the following algorithm explicitly computes a set of $2r$ functions for Ω described in Darboux's theorem:

- 1. Find vector fields $X_1, X_2 \in \mathfrak{X}(U)$ such that:
	- $\left(\begin{array}{cc} -1 & -1 \\ 1 & -1 \end{array} \right)$
	- $\mathcal{L} \setminus \mathcal{L}$ and $\mathcal{L} \setminus \mathcal{L}$ is the set of $\mathcal{L} = \mathcal{L}$
	- $(1 2)$ $(2 1)$ $(3 2)$ $(3 2)$ $(2 1)$ $(3 2)$
	- $(1, 1, 1, 2)$
- $\mathbf{1} \bullet \mathbf{1} \bullet \mathbf{$
- β . Repeat steps 1 and 2 a further $r 2$ more times until $\alpha = 0$,
- 4. Apply Theorem 3.15 with a solvable structure of two symmetries $X_3, X_4 \in$ $\mathbf{x} \cdot \mathbf{y}$, and $\mathbf{x} \cdot \mathbf{y}$ is a non-trivial symmetry \mathbf{y} is an operator \mathbf{y} non-trivial symmetry of X3 with the property that (X3; X4) = 1.

Proof. Let $\Omega \in \Lambda^1(U)$ be closed with $\Omega \neq 0$ and $\Omega^{-1} = 0$ for some $r \in \mathbb{N}$. From Theorem 7.7 and then Theorem 7.5, we can compute $\Omega_1 \in \Lambda^2(U)$, where

$$
\Omega_1 = \Omega + dg_1 \wedge df_1,
$$

for some smooth f_1 and g_1 , with Ω_1^+ = \neq 0 and Ω_1^+ = 0. Then once again from Theorem 7.7 followed by Theorem 7.5, $\Omega_2 \in \Lambda$ (U) can be computed so that

$$
\Omega_2 = \Omega + dg_1 \wedge df_1 + dg_2 \wedge df_2,
$$

for some smooth f_2 and g_2 , with $\Omega_2^- \neq 0$ and $\Omega_2^- \equiv 0$. Continuing in this when r is of the form r

$$
\Omega_{r-1} = \Omega + dg_1 \wedge df_1 + dg_2 \wedge df_2 + \cdots + dg_{r-1} \wedge df_{r-1},
$$

such that $\Omega_{r-1} \neq 0$ and $\Omega_{r-1} = 0$. Applying step 4, Ω_{r-1} is closed, and from Theorem 7.1, r1 is also de
omposable. From Theorem 3.15 and Corollary 3.13, with X3 as ^a non-trivial symmetry of r1 and X4 as ^a nonrivial symmetry of $\mathbf{u} = \mathbf{u} + \mathbf{u}$ is a subset of $\mathbf{u} = \mathbf{u} + \mathbf{u}$ is a subset of $\mathbf{u} = \mathbf{u} + \mathbf{u}$

$$
\frac{X_{3}\Box \Omega_{r-1}}{X_{4}\Box X_{3}\Box \Omega_{r-1}} = dg_r,
$$

$$
\frac{X_{4}\Box \Omega_{r-1}}{X_{3}\Box X_{4}\Box \Omega_{r-1}} = df_r + \lambda dg_r,
$$

for some some smooth f , g , where f , g and f

$$
\Omega_{r-1} = \Omega_{r-1}(X_3, X_4)df_r \wedge dg_r = df_r \wedge dg_r.
$$

Therefore

$$
\Omega = df_1 \wedge dg_1 + df_2 \wedge dg_2 + \cdots + df_{r-1} \wedge dg_{r-1} + df_r \wedge dg_r.
$$

 \Box

Remark 1. In looking for two symmetries that satisfy the four conditions in Theorem 7.8, condition (d) can be relaxed a little by only requiring that onstants and XI or X2 may be seen and the seasons of the second computer of the second computer of the second while still satisfying the other three onditions. The same holds true for the two symmetries in step 4.

Remark 2. Conditions (a) and (b) are strong requirements, and may be difto the protection in the they imply the to satisfy. Since μ is the property imply μ imply μ implies the s hosen such and X2 is a losen such a losen the result in Theorem 7.8 is of more theoretical significance than practical use.

We an provide an alternative to the requirement in step 4 in Theorem 7.8 as follows:

Lemma 1.9. Let $\Omega \in \Lambda^2(U)$ be some arbitrary closed two-form. Suppose there exists some $\mathbf{1}_{\mathbf{0}}$ $\mathbf{1}_{\mathbf{0}}$ is that in that $\mathbf{1}_{\mathbf{0}}$ is that in that in that

$$
\mathcal{L}_{X_3}\Omega = 0,\tag{35}
$$

 \Box

and α is the set of α

$$
\mathcal{L}_{X_4}(X_3 \lrcorner \Omega) = 0.
$$

Proof.

$$
\mathcal{L}_{X_4}(X_3 \cup \Omega) = d(X_4 \cup X_3 \cup \Omega) + X_4 \cup d(X_3 \cup \Omega) = X_4 \cup (\mathcal{L}_{X_3} \Omega) = 0,
$$

using that I also also and the 1, equation (35), and the that is an extended.

We now apply the algorithm in Theorem 7.8 and the modification of Step 4 in Lemma 7.9 to an example. It is important to realise that the difficult part in applying Theorem 7.8 is in finding the first $r-1$ pairs of symmetries X_1, X_2 . Nevertheless, the main purposes of this example are to illustrate: i) the crucial role Theorem 7.5 plays in reducing the number of terms in a two-form by one; and ii) the flexibility in choosing X_4 in Lemma 7.9.

Example 7.10. Consider the following two-form $\Omega \in \Lambda$ (U), where U is some suitably chosen four-dimensional, open, convex neighbourhood of \mathbb{R}^4 with coordinates x_1, x_2, x_3 :

$$
\Omega := \left(\frac{x^1}{x^2}\right) \left(\frac{x^3}{x^2} - 2\right) dx^1 \wedge dx^2 + \frac{x^1}{x^2} dx^1 \wedge dx^3 - \frac{2x^1}{x^4} dx^1 \wedge dx^4 - \left(\frac{x^1}{x^2}\right)^2 dx^2 \wedge dx^3.
$$

Now it is easy to show that $a_2 z = 0$, $\Omega^+ \neq 0$ and $\Omega^+ \equiv 0$. We may then pro
eed to apply Theorem 7.8. Let

$$
X_1 := -\frac{1}{x^3} \left(\frac{x^2}{x^1}\right)^2 \frac{\partial}{\partial x^2} + \frac{x^2 x^4}{(x^1)^2 x^3} \frac{\partial}{\partial x^4}.
$$

Now

$$
\mathcal{L}_{X_1}\Omega = d\left(X_1 \Box \Omega\right),
$$

= $d\left(\frac{1}{x^3}dx^3 + \frac{2x^2}{x^1x^3}dx^1 + \frac{1}{x^3}\left(\frac{x^2}{x^1}\right)\left(\frac{x^3}{x^2} - 2\right)dx^1\right),$
= $d\left(\frac{1}{x^3}dx^3 + \frac{1}{x^1}dx^1\right) = 0,$

so condition (a) of step 1 in Theorem 7.8 is met. Hence

$$
X_1 \Box \Omega = d \left(\ln |x^1 x^3| \right).
$$

Let

$$
X_2 := x^3 \frac{\partial}{\partial x^3}
$$

where α is satisfied. The isometries of α is satisfied to the satisfied of α is satisfied.

$$
X_2 \Box \Omega = \frac{x^1 x^3}{x^2} dx^1 + x^3 \left(\frac{x^1}{x^2}\right)^2 dx^2,
$$

it is not hard to show the strip $\{1,2,3\}$, $\{2,3,4\}$, $\{3,4,5\}$, $\{4,7,8\}$, $\{5,7,8\}$, $\{6,7,8\}$ Also,

$$
\begin{aligned} \left(\mathcal{L}_{X_2}\Omega\right) \wedge \left(X_{1\Box}\Omega\right) &= d\left(X_{2\Box}\Omega\right) \wedge \left(X_{1\Box}\Omega\right), \\ &= \left(\frac{x^1}{x^2}dx^1 \wedge dx^3 - \frac{x^1x^3}{(x^2)^2}dx^1 \wedge dx^2 - \left(\frac{x^1}{x^2}\right)^2 dx^2 \wedge dx^3 \right) \\ &+ \frac{2x^1x^3}{(x^2)^2}dx^1 \wedge dx^2\right) \wedge \left(\frac{1}{x^1}dx^1 + \frac{1}{x^3}dx^3\right), \\ &= 0, \end{aligned}
$$

so condition (b) is met. Now

$$
d(X_{2}\mathbf{1}\Omega) = 0 \mod X_{1}\mathbf{1}\Omega.
$$

Using the foliated derivative, this implies

$$
X_2 \Box \Omega = dg_1 + \lambda_1 d \left(\ln |x^1 x^3| \right), \tag{36}
$$

for some $g_1, \lambda_1 \in C^{\infty}(U^{\perp})$. Performing a coordinate substitution gives

$$
X_2 \Box \Omega = -d \left(\frac{(x^1)^2 x^3}{x^2} \right) + \frac{(x^1)^2 x^3}{x^2} d \left(\ln |x^1 x^3| \right).
$$

Therefore

$$
(X_2 \Box \Omega) \wedge (X_1 \Box \Omega) = -d\left(\frac{(x^1)^2 x^3}{x^2}\right) \wedge d\left(\ln|x^1 x^3|\right) = -d\left(\frac{x^1}{x^2}\right) \wedge d(x^1 x^3).
$$

For other choice of X_1, X_2 , we may obtain an expression for the other twoform omponent of .

 $\frac{1}{1}$: $\frac{1}{2}$, $\frac{1}{1}$, $\frac{1}{1}$

$$
\Omega_1 = -\frac{2x^1}{x^2} dx^1 \wedge dx^2 - \frac{2x^1}{x^4} dx^1 \wedge dx^4.
$$

It is clear that $a_1 = 0$ and $a_1 = 0$ as expected, so we may proceed to apply the nal step in Theorem 7.8 on 1. Dening

$$
X_3 := x^1 x^4 \frac{\partial}{\partial x^4}
$$

we have

$$
\mathcal{L}_{X_3}\Omega_1 = d\left(X_3 \Box \Omega_1\right) = d\left(2(x^1)^2 dx^1\right) = 0.
$$

This implies

$$
X_3 \Box \Omega_1 = d \left(\frac{2(x^1)^3}{3} \right). \tag{37}
$$

Now hoose

$$
X_4 := \frac{1}{2(x^1)^2} \frac{\partial}{\partial x^1},
$$

so that \mathcal{L}_{4} is that \mathcal{L}_{1} . From Lemma 7.19, lemma \mathcal{L}_{3} (x3 \mathcal{L}_{4} (x3 \mathcal{L}_{5}), and hence from \mathcal{L}_{5} Theorem 3.15,

$$
X_4 \Box \Omega_1 = df_2 + \lambda_2 d \left(\frac{2(x^1)^3}{3} \right), \tag{38}
$$

for some $j_2, \lambda_2 \in C^\infty(U^*)$. To find j_2 , it is easy to show that

$$
X_4 \Box \Omega_1 = -d\left(\frac{1}{x^1} \ln|x^2 x^4|\right) \mod dx^1,
$$

and hen
e

$$
\Omega_1 = d\left(\frac{1}{x^1} \ln|x^2 x^4|\right) \wedge d\left(\frac{2(x^1)^3}{3}\right).
$$

On
e again we may simplify this:

$$
d\left(\frac{1}{x^1}\ln|x^2x^4|\right) \wedge d\left(\frac{2(x^1)^3}{3}\right) = 2(x^1)^2 d\left(\frac{1}{x^1}\ln|x^2x^4|\right) \wedge dx^1,
$$

= $2x^1 d\left(\ln|x^2x^4|\right) \wedge dx^1,$
= $d\left(\ln|x^2x^4|\right) \wedge d\left(\left(x^1\right)^2\right).$

Thus

$$
\Omega = d\left(\frac{x^1}{x^2}\right) \wedge d\left(x^1 x^3\right) + d\left(\ln|x^2 x^4|\right) \wedge d\left(\left(x^1\right)^2\right).
$$

8Summary

Using the idea of a solvable symmetry structure we presented various algorithms for expressing certain classes of differential forms in terms of simplified oordinate systems. We began by reviewing Lie's symmetry approa
h and then showed that it may applied to simplify differential forms which are deomposable and losed modulo themselves. We then gave a result showing that certain types of symmetry structures in Theorem 3.15 forced more than one of the ω to become closed, and looked at under what conditions a given differential form was decomposable and closed modulo itself.

Next, we examined the problem of finding simplifying coordinates for the Pfaffian problem. This was treated by imposing a special condition on the solvable symmetry structure applied to the Cauchy characteristic space of the differential ideal generated by the differential form $(d\alpha)^r \wedge \alpha$, where α was the Pfaffian form, and r was its rank.

Finally, we looked at differential two-forms where the main result there was an algorithm for finding the coordinates in Darboux's theorem, derived from the well-known iterative scheme, where a pair of new coordinates is extracted each time.

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