

Exterior Differential Systems with Symmetry

Mark Fels, Utah State University

July 25, 2006

Definition. *An exterior differential system is a subset of $\mathcal{I} \subset \Omega^*(M)$, which is closed under $\Omega^*(M) \wedge$ and exterior differentiation*

$$\mathcal{I} = \langle \theta^i, d\theta^i, \beta^a, d\beta^a, \dots, \rangle$$

where $\langle \ \rangle$ means algebraically generated by. The EDS that will be of interested are mainly (but not exclusively) **Pfaffian systems** - those generated by one-forms and their derivatives.

$$\mathcal{I} = \langle \theta^i, d\theta^i \rangle$$

Definition. *An integral manifold of \mathcal{I} is an immersion $s : N \rightarrow M$ such that $s^*\mathcal{I} = 0$.*

Definition. *A symmetry of an exterior differential system is a diffeomorphism $\phi : M \rightarrow M$ such that $\phi^*\mathcal{I} = \mathcal{I}$. A symmetry group of \mathcal{I} will be a Lie group G acting smoothly on M where each diffeomorphism $g : M \rightarrow M$ is a symmetry of \mathcal{I} .*

If ϕ is a symmetry of the EDS \mathcal{I} and $s : N \rightarrow M$ is an integral manifold of \mathcal{I} , then

$$(\phi \circ s)^*\mathcal{I} = s^*\phi^*\mathcal{I} = s^*\mathcal{I} = 0.$$

Therefore symmetries map integral manifolds to integral manifolds.

From now on:

- 1) G is a Lie group acting smoothly on M
- 2) \mathfrak{g} Lie algebra of right invariant vector-fields
- 3) γ - Lie algebra of infinitesimal generators
- 4) $\rho : \mathfrak{g} \rightarrow \gamma$ the homomorphism
- 3) $\Gamma \subset TM$ is the point-wise span of γ

Differential Equations give rise to EDS, and solutions to the differential equations are integral manifolds. Next are some example....

Example 1. The Chazy equation

$$y_{xxxx} = 2yy_{xxx} - 3y_x^2$$

gives rise to the EDS on $M = (x, y, y_x, y_{xx})$

$$\mathcal{I} = \langle dy - y_x dx, dy_x - y_{xx} dx, dy_{xx} - (2yy_{xxx} - 3y_x^2) dx \rangle$$

Solutions to the Chazy equation are integral manifolds.

The EDS is invariant with respect to the (local) action of $SL(2)$ on M :

$$\begin{aligned} \hat{x} &= \frac{ax + b}{cx + d}, & \hat{y} &= (cx + d)^2 y + 6c(cx + d) \\ \hat{y}_{\hat{x}} &= \frac{d\hat{y}}{dx} \frac{dx}{d\hat{x}} & \hat{y}_{\hat{x}\hat{x}} &= \frac{d\hat{y}_{\hat{x}}}{dx} \frac{dx}{d\hat{x}} \end{aligned}$$

where $ad - bc = 1$.

Example 2. The geodesic equation for

$$\eta = e^{-\frac{4}{3}x_4}(dx_1dx_3 - dx_2dx_2) + e^{\frac{2}{3}x_4}dx_3dx_3 + cdx_4dx_4$$

where $c < 0$ – Lorentz and $c > 0$ – split sig..

The EDS on $M = (t, x_i, \dot{x}_i), 1 \leq i \leq 4$ is

$$\mathcal{I} = \langle dx_i - \dot{x}_i dt,$$

$$d\dot{x}_1 - \frac{2}{3}\dot{x}_4(2\dot{x}_1 - 3\dot{x}_3e^{2x_4})dt$$

$$d\dot{x}_2 - \frac{4}{3}\dot{x}_2\dot{x}_3dt, \quad d\dot{x}_3 - \frac{4}{3}\dot{x}_3\dot{x}_4dt$$

$$d\dot{x}_4 + \frac{1}{3c} \left(4e^{-\frac{4}{3}x_4}(\dot{x}_1\dot{x}_3 + 2\dot{x}_2^2) + e^{\frac{2}{3}x_4}\dot{x}_3^2 \right) dt \rangle$$

The geodesics are integral manifolds.

The EDS is invariant with respect to time translations and the isometry group (5-d solvable) whose Killing vector-fields are

$$X_1 = \partial_{x_1}, \quad X_2 = \partial_{x_2}, \quad X_3 = \partial_{x_3},$$

$$X_4 = x_2\partial_{x_1} + x_3\partial_{x_2}$$

$$X_5 = -5x_1\partial_{x_1} - 2x_2\partial_{x_2} + x_3\partial_{x_3} - 3\partial_{x_4}$$

Example 3. The contact system on $M = J^2(\mathbb{R}, \mathbb{R}^2)$ with coordinates $(t, x, y, \dot{x}, \dot{y}, \ddot{x}, \ddot{y})$:

$$\mathcal{I} = \langle dx - \dot{x}dt, dy - \dot{y}dy, d\dot{x} - \ddot{x}dt, d\dot{y} - \ddot{y}dt, \\ d\ddot{x} \wedge dt, d\ddot{y} \wedge dt \rangle$$

Any prolonged graph $(x(t), y(t))$ is an integral manifold.

The $E^+(2)$ action

$$\begin{pmatrix} \hat{x} \\ \hat{y} \end{pmatrix} = \begin{pmatrix} \cos \theta & -\sin \theta \\ \sin \theta & \cos \theta \end{pmatrix} \begin{pmatrix} x \\ y \end{pmatrix} + \begin{pmatrix} a \\ b \end{pmatrix}$$

$$\begin{pmatrix} \hat{\dot{x}} \\ \hat{\dot{y}} \end{pmatrix} = \begin{pmatrix} \cos \theta & -\sin \theta \\ \sin \theta & \cos \theta \end{pmatrix} \begin{pmatrix} \dot{x} \\ \dot{y} \end{pmatrix}$$

$$\begin{pmatrix} \hat{\ddot{x}} \\ \hat{\ddot{y}} \end{pmatrix} = \begin{pmatrix} \cos \theta & -\sin \theta \\ \sin \theta & \cos \theta \end{pmatrix} \begin{pmatrix} \ddot{x} \\ \ddot{y} \end{pmatrix}$$

is a symmetry group of \mathcal{I} .

Example 4. The contact system on $M = J^2(\mathbb{R}, \mathbb{R}) \times J^2(\mathbb{R}, \mathbb{R})$: $(x, u, u_x, u_{xx}, y, v, v_y, v_{yy})$

$$\mathcal{I} = \langle du - u_x dx, du_x - u_{xx} dx, du_{xx} \wedge dx, \\ dv - v_y dy, dv_y - v_{yy} dy, dv_{yy} \wedge dy \rangle$$

This is also the system of PDE $u_y = 0, v_x = 0$ prolonged. Any prolonged graph $u = f(x), v = g(y)$ is an integral manifold.

Consider local symmetries of $SL(2, \mathbb{R})$

$$\begin{aligned} \hat{u} &= \frac{au + b}{cu + d}, & \hat{v} &= \frac{av + b}{cv + d} \\ \hat{u}_x &= \frac{u_x}{(cu + d)^2}, & \hat{v}_y &= \frac{v_y}{(cv + d)^2} \\ \hat{u}_{xx} &= \frac{u_{xx}}{(cu + d)^2} - \frac{2cu_x^2}{(cu + d)^3} \\ \hat{v}_{yy} &= \frac{v_{yy}}{(cv + d)^2} - \frac{2cv_y^2}{(cv + d)^3} \end{aligned}$$

where $ad - bc = 1$.

Suppose $\mathbf{q} : M \rightarrow M/G = \bar{M}$ is a submersion.

Then $\Gamma \subset TM$ is a rank q -subbundle, where q is the dimension of the orbits. Also

$$\ker \mathbf{q}_* = \Gamma$$

Definition. The quotient $\bar{\mathcal{I}} \subset \Omega^*(\bar{M})$ of \mathcal{I} is

$$\bar{\mathcal{I}} = \{ \bar{\theta} \in \Omega^*(\bar{M}) \mid \mathbf{q}^*\bar{\theta} \in \mathcal{I} \}$$

This definition is difficult to work with without any assumptions on the EDS.

One invariant condition is constant rank.

If the Pfaffian system \mathcal{I} is constant rank r then $\mathcal{I}^1 = \mathcal{I} \cap \Omega^1(M)$ are sections of a rank r subbundle $I \subset T^*M$.

The symmetries of \mathcal{I} then preserve I : $g^*I = I$.

The quotient $\bar{I} \subset T^*\bar{M}$ defined point-wise is

$$\bar{I}_{\bar{x}} = \{ \bar{\theta} \in T_{\bar{x}}^*\bar{M} \mid \mathbf{q}^*\bar{\theta}_{\bar{x}} \in I_x \text{ where } \mathbf{q}(x) = \bar{x} \}$$

This can be computed using $I^{\mathbf{An}} \subset TM$

$$I_x^{\mathbf{An}} = \{ V \in T_x M \mid \theta(V) = 0 \forall \theta \in I_x \}$$

which is a rank $n - r$ sub-bundle. Then

$$\bar{I} = \left(\mathbf{q}_*(I^{\mathbf{An}}) \right)^{\mathbf{An}}$$

\bar{I} is a sub-bundle if and only if

$$\dim I_x^{\mathbf{An}} \cap \ker \mathbf{q}_* = \dim I_x^{\mathbf{An}} \cap \Gamma_x = k, \quad \forall x \in M,$$

in which case $\text{rank } \bar{I} = r + k - q$, with $\text{rank } \Gamma = q$.

Computing

$$I_x^{\mathbf{An}} \cap \Gamma_x = \{Z \in \Gamma_x \mid \theta(Z) = 0, \forall \theta \in I_x\}$$

in the basis θ^α for I_x , and Z_i a basis Γ_x is

$$I_x^{\mathbf{An}} \cap \Gamma_x = \ker \theta^\alpha(Z_i). \quad (1)$$

How to compute $\bar{\mathcal{I}}$ (or \bar{I})?

1) Choose $\sigma : \bar{M} \rightarrow M$ a cross-section

2) Compute the semi-basic forms in \mathcal{I} (or I):

$$I_{x, sb} = \{\theta \in I_x \mid \theta(Z) = 0 \forall Z \in \Gamma_x\}$$

(the kernel in (1) on the form side)

3) The pullback $\sigma^* \theta_{sb} \in \bar{\mathcal{I}}$, and generate $\bar{\mathcal{I}}$.

\bar{I} being constant rank is also equivalent to I_{sb} being a constant rank sub-bundle'.

Observations:

Theorem. *If \mathcal{I} is a rank r completely integrable Pfaffian system, and $\text{rank } I^{\mathbf{An}} \cap \Gamma = k$, then $\bar{\mathcal{I}}$ is a rank $r + k - q$ completely integrable Pfaffian system.*

If $\text{rank } I^{\mathbf{An}} \cap \Gamma = k$ but \mathcal{I} is not completely integrable, then $\bar{\mathcal{I}}$ is not necessarily a Pfaffian system. It is possible to give sufficient conditions so that $\bar{\mathcal{I}}$ is Pfaffian system.

Theorem. *Suppose \mathcal{I} is a constant rank Pfaffian system invariant with respect to G and $Z_{x_0} \in I_{x_0}^{\mathbf{An}} \cap \Gamma_{x_0}$ with $Z_{x_0} \neq 0$. Then $e^{t\mathbf{z}}x_0$, $\mathbf{z} \in \mathfrak{g}$ is a one-dimensional integral manifold, where $Z_{x_0} = \rho(\mathbf{z})_{x_0}$.*

Example 1. $SL(2, \mathbb{R})$, Chazy. The EDS is

$$\theta^1 = dy - y_x dx,$$

$$\theta^2 = dy_x - y_{xx} dx,$$

$$\theta^3 = dy_{xx} - (2yy_{xx} - 3y_x^2) dx.$$

A set of generators for γ are,

$$X_1 = \partial_x,$$

$$X_2 = 2x\partial_x - 2y\partial_y - 4y_x\partial_{y_x} - 6y_{xx}D_{y_{xx}} - 8y_{xxx}\partial_{y_{xxx}}$$

$$X_3 = -x^2\partial_x + 2(xy + 3)\partial_y + 2(2xy_x + y)\partial_{y_x} \\ + 6(y_x + xy_{xx})D_{y_{xx}} + (12y_{xx} + 8xy_{xxx})\partial_{y_{xxx}}$$

The determinant $\det(\theta^i(X_j)) = 0$ if and only if

$$y_{xx} = yy_x - \frac{1}{9}y^3 \pm (y^2 - 6y_x)^3.$$

For initial conditions $(x^0, y^0, y_x^0, y_{xx}^0)$ satisfying this constraint, the (unique) solution is the one-dimensional orbit:

$$x = x^0 + 2t \frac{(\delta^0 + 3y_x^0 + y^0 \sqrt{\delta^0})}{ty_x^0(y^0 + \sqrt{\delta^0} - 1)},$$

$$y = y^0 + 2ty_x^0 \left((3t(y_x^0)^2 - y^0 + ty_x^0 \delta) \sqrt{\delta} + ty_x^0 y^0 \delta \right)$$

where $\delta_0 = (y^0)^2 - 6y_x^0$.

There is a 2 parameter family of invariant solutions. (The solutions as a graph are easily found).

Example 2. Geodesics: The EDS is

$$\begin{aligned} \mathcal{I} = & \langle dx_i - \dot{x}_i dt, d\dot{x}_1 - \frac{2}{3}\dot{x}_4(2\dot{x}_1 - 3\dot{x}_3 e^{2x_4})dt \\ & d\dot{x}_2 - \frac{4}{3}\dot{x}_2\dot{x}_3 dt, \quad d\dot{x}_3 - \frac{4}{3}\dot{x}_3\dot{x}_4 dt \\ & d\dot{x}_4 + \frac{1}{3c} \left(4e^{-\frac{4}{3}x_4}(\dot{x}_1\dot{x}_3 + 2\dot{x}_2^2) + e^{\frac{2}{3}x_4}\dot{x}_3^2 \right) dt \rangle \end{aligned}$$

At the point

$$t = 0, \quad \mathbf{x} = (0, 0, 0, 0), \quad \dot{\mathbf{x}} = \left(\frac{c_0}{4}, 0, c_0, 0 \right), \quad (2)$$

where $c_0 \neq 0$ the vector-field $X \in \gamma$,

$$X = \partial_t + \frac{c_0}{4}X_1 + c_0X_3 = \partial_t + \frac{c_0}{4}\partial_{x_1} + c_0\partial_{x_3}$$

satisfies $\theta(X) = 0, \forall \theta \in \mathcal{I}$ at (2). The integral curve of X in M through the point (2)

$$x_1 = \frac{c_0}{4}t, \quad x_2 = 0, \quad x_3 = c_0t, \quad x_4 = 0$$

and is a geodesic which is the orbit of a one parameter sub-group corresponding to X .

Homogeneous geodesics always exist for homogeneous Riemannian manifolds.

Example 5. The standard EDS for the differential equation

$$u_{xxxxx} = \frac{5u_{xxx}(9u_{xxxx}u_{xx} - 8u_{xxx}^2)}{9u_{xx}^2}$$

Is invariant with respect to the five dimensional symmetry group

$$G = SA(2) = \{ (A, b) \mid A \in SL(2, \mathbb{R}), b \in \mathbb{R}^2 \},$$

acting on (x, u) and then prolonged. Every solution is the orbit of a one-parameter subgroup, giving the general solution,

$$u = c_0 + c_1x \pm \sqrt{c_3x^2 + c_x x + c_2(4c_3c_2 - c_x^2)}.$$

Example 6. Every geodesic on a Riemannian symmetric space is homogeneous.

Remark: It is sometimes possible to generate integral manifolds with dimension > 1 when $\text{rank } I_x^{\mathbf{An}} \cap \Gamma_x > 1$.

Why the quotient?

- 1) We want to split the problem of finding integral manifolds for \mathcal{I} into finding integral manifolds for $\bar{\mathcal{I}}$, and then build integral manifolds utilizing G . (This second part is sometimes called the reconstruction problem).
- 2) Classify the integral manifolds that are inequivalent with respect to G .
- 3) The inverse problem: Use \mathcal{I} to simplify finding integral manifolds to $\bar{\mathcal{I}}$.

Let's look at problem 1. In order to implement the decomposition we need to partition M into G -invariant subsets on which we can control the behavior of the quotient. The key is:

$$I_x^{\mathbf{An}} \cap \Gamma_x$$

Consider the two G -invariant subsets

$$K = \{ x \in M \mid \Gamma_x \subset I_x^{\mathbf{An}} \}$$

and the transverse subset

$$M^t = \{ x \in M \mid \Gamma_x \cap I_x^{\mathbf{An}} = 0 \}.$$

What can be said about \mathcal{I} and $\bar{\mathcal{I}}$ and the reconstruction problem on these subsets? (which can be empty)

We start with K . Assume

1) $\iota : K \rightarrow M$ is an embedded submanifold,

2) the action of G on K is regular: $\mathbf{q} : K \rightarrow \bar{K}$

Let $\mathcal{I}_K = \iota^*\mathcal{I}$, which under obvious conditions is a constant rank Pfaffian system (rank depends on the embedding). In this case $\bar{\mathcal{I}}$ is a constant rank Pfaffian system with the same rank as \mathcal{I}_K .

Note that at each point in K , all forms in \mathcal{I}_K are semi-basic!

An integral manifold $s : N \rightarrow M$ is G invariant if $gs(N) = N$. These can be found as integral manifolds to \mathcal{I}_K . Here's how:

Theorem. *Let $\bar{N} \subset M$ be an embedded integral manifold of $\bar{\mathcal{I}}_K$. Then $N = \mathbf{q}^{-1}(\bar{N}) \subset M$ is a G -invariant integral manifold of \mathcal{I} .*

Proof: N is clearly G -invariant. So need to show it is an integral manifold.

Let $x \in N$, $X \in T_x N$, and $\bar{x} = \mathbf{q}(x)$, $\bar{X} = \mathbf{q}_* X$. Note $\bar{x} \in \bar{N}$ and $\bar{X} \in T_{\bar{x}} \bar{N}$.

Choose an open set $\bar{U} \subset \bar{K}$ containing \bar{x} , and a cross-section $\sigma : \bar{M} \rightarrow M$ with $\sigma(\bar{x}) = x$. Then

$$X = \sigma_* \bar{X} + V$$

for some $V \in \Gamma_x$. Evaluating on $\theta \in \mathcal{I}$,

$$\begin{aligned} \theta(X) &= \theta(\sigma_* \bar{X}_{\bar{x}} + V) \\ &= \sigma^* \theta(\bar{X}) + \theta(V) \\ &= 0. \end{aligned}$$

The first term vanishes because, all one-forms in \mathcal{I}_K are semi-basic so pullback to $\bar{\mathcal{I}}_K$, and \bar{N} is an integral manifold. The second term vanishes because we are at point of K (so that $V \in I_x^{\mathbf{An}}$). QED.

Remarks:

- 1) The reconstruction problem is algebraic.
- 2) Integral manifolds of \mathcal{I}_K can always be enlarged (locally) to be invariant.
- 3) K is the subset of M on which Γ are Cauchy-characteristics for \mathcal{I}_K . This theorem is not so surprising.

The transverse subset M^t .

1) Integral manifolds in M^t don't have continuous symmetry.

2) For $s : N \rightarrow M$, an integral manifold $\mathbf{q} \circ s : N \rightarrow \bar{M}^t$ is an integral manifold. (Ie. immersion property still holds).

If M^t is non-empty then the restriction

$$\mathcal{I}_{M^t} = \mathcal{I}|_{M^t}$$

has the same rank as \mathcal{I} , but the quotient $\bar{\mathcal{I}}_{M^t}$ is not necessarily a Pfaffian system. (It is similar though.)

The reconstruction problem is:

Theorem. *Let $\bar{N} \rightarrow \bar{M}^t$ be an embedded integral manifold of $\bar{\mathcal{I}}_{M^t}$. Then $\mathcal{I}|_{\mathbf{q}^{-1}(\bar{N})}$ is completely integrable, and the leaves are integrable manifolds of \mathcal{I} .*

As a consequence of this theorem, the integral manifolds of \mathcal{I} are surjective via \mathbf{q} onto the integral manifolds of $\bar{\mathcal{I}}_{M^t}$.

If the action of G on M^t is free, there is a nice geometric way to think about the reconstruction problem.

Let $s : \bar{N} \subset \bar{M}^t$ be an integral manifold of $\bar{\mathcal{I}}_{M^t}$, and let $\hat{s} : \bar{N} \rightarrow M$ be any cover of \bar{N} . The integral manifold N of \mathcal{I} which projects to \bar{N} is of the form

$$s(t) = \mu(A(t), \hat{s}(t))$$

where $A : \bar{N} \rightarrow G$.

In order for $s(t)$ to be an integral manifold of \mathcal{I} , $A(t)$ satisfies a generalized equation of Lie type. These are integrable by quadratures for (s.c) solvable Lie groups.

Example 3. $E^+(2)$ symmetry on $J^2(\mathbb{R}, \mathbb{R}^2)$ with the standard contact structure:

$$\mathcal{I} = \langle dx - \dot{x}dt, dy - \dot{y}dy, d\dot{x} - \ddot{x}dt, d\dot{y} - \ddot{y}dt, \\ d\ddot{x} \wedge dt, \ddot{y} \wedge dt \rangle$$

(Curves $(x(t), y(t))$ are integral manifolds)

The infinitesimal generators are

$$\gamma = \text{span} \{ \partial_x, \partial_y, x\partial_y - y\partial_x + \dot{x}\partial_{\dot{y}} - \dot{y}\partial_{\dot{x}} + \ddot{x}\partial_{\ddot{y}} - \ddot{y}\partial_{\ddot{x}} \}$$

and $E^+(2)$ action is transverse at

$$M^t = \{ \sigma \in J^2(\mathbb{R}, \mathbb{R}^2) \mid (\dot{x}, \dot{y}) \neq (0, 0) \}$$

The quotient is 4 dimensional $\bar{M}^t = (t, v, k_1, k_2)$.

Let $\sigma : \bar{M}^t \rightarrow M$ be the cross-section,

$$(t, x = 0, y = 0, \dot{x} = 0, \dot{y} = v, \ddot{x} = k_1, \ddot{y} = k_2).$$

The quotient EDS is (pullback semi-basic forms)

$$\bar{\mathcal{I}}_{M^t} = \langle dv - k_2dt, dk_2 \wedge dt, (k_1dk_1 + k_2dk_2) \wedge dt \rangle$$

A typical integral manifold for $\bar{\mathcal{I}}_{M^t}$ is

$$\bar{s}(t) = (t, v = v(t), k_2 = \frac{dv}{dt}, k_1 = k(t)), v(t) \neq 0.$$

An integral manifold in M^t which projects to \bar{s} is of the form *

$$s(t) = \mu(A(t), \sigma \circ \bar{s}(t))$$

where $A : \mathbb{R} \rightarrow E^+(2)$ satisfies

$$\begin{aligned} \frac{da}{dt} &= -v(t) \sin \theta(t) = 0, & \frac{db}{dt} &= v(t) \cos \theta(t), \\ \frac{d\theta}{dt} &= -\frac{k_1(t)}{v(t)}. \end{aligned}$$

An equation of Lie type for $\alpha : \mathbb{R} \rightarrow \mathfrak{g}$,

$$\alpha(t) = \left(0, -v(t), \frac{k_1(t)}{v(t)} \right).$$

*Finding an integral manifold to \mathcal{I} projecting to \bar{s} is the prescribed "curvature" problem

Example 4. $SL(2)$ on two copies of jet-space.
 On the set $u_x v \neq 0$, the action satisfies

$$I^{\mathbf{An}} \cap \Gamma = 0.$$

Choosing the cross-section on $\bar{M}^t = (x, y, w, w_x, w_y)$

$$\begin{aligned} x &= x, y = x, u = 0, v = 1, u_x = w, v_y = 1, \\ u_{xx} &= w_x - 2w^2, v_{yy} = w_y/w + 2 \end{aligned}$$

The quotient EDS is

$$\begin{aligned} \bar{\mathcal{I}}_{M^t} = & \langle dw - w_x dx - w_y dy, dw_x dx + dw_y dy, \\ & \left(dw_x + \left(\frac{w_x w_y}{w} - w^2 \right) dy \right) \wedge dx \rangle. \end{aligned}$$

Project integral manifold $f(x), g(y)$ of \mathcal{I} to

$$w = \frac{u_x v_y}{(u - v)^2} = \frac{f'(x)g'(y)}{(f(x) - g(y))^2}$$

integral manifolds of $\bar{\mathcal{I}}$.

If $w = e^u$ - Monge Ampere form of Liouville.

Exercise: Compute the equation of Lie type for the Chazy equation on the set of transverse initial conditions.

Is there a way to do this all at once?

Yes: In fact it might look familiar:

Suppose \mathcal{I} is constant rank r Pfaffian and θ^i are a basis of sections.

Then $\mu : M \rightarrow \mathbb{R}^r \otimes \mathfrak{g}^*$ given by

$$\theta^i(X)$$

is the moment map.

K is the zero-set of the moment map.

M^t is the full rank set for the moment map.

Reduction is the same as symplectic or contact reduction - restriction then quotient.