

Symmetries of Vector Exterior Differential Systems and the Inverse Problem in Second-Order Ostrograds'kii Mechanics

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Abstract

Symmetries for variational problems are considered as symmetries of vector-valued exterior differential systems. This approach is applied to equations for the classical spinning particle.

It is common in the similarity theories of differential equations to consider transformations which leave invariant some prolonged manifolds in jet-spaces of appropriate order, within which the integral manifolds of a given system of partial differential equations lie [1]. From this invariance it follows that solutions of the system in charge transform but into some other solutions. However, we are not sure enough that the prolonged manifold is “densely” covered by the integral manifolds. If not, then some additional symmetries may occur. In this paper we consider a symmetry of a differential equation defined in general as a generator of such a transformation, which carries every solution into some other solution.

Then we reformulate this into a slightly different language of algebraic invariance of some vector-valued exterior differential system. This has an advantage that we can use the Lagrange multiplies method to solve the problem. In the case of ordinary differential equations the two approaches (algebraic and general ones) coincide due to completeness of Pfaff systems (that is every differential form, annulled by every solution of the system, expands by the Lagrange multiplies into elements of the differential ideal generated by the system itself).

The Euler-Lagrange equations of variational calculus (of arbitrary order) naturally fall into the framework of vector bundle valued exterior differential systems when the dependent variables are globally segregated from the independent ones [2]. This is the case of the field theory. In general, such a segregation depends on the local chart $\mu : M \rightarrow \mathbb{R}^p \times \mathbb{R}^q$ which constitutes a way of introducing a fibred structure in the manifold M . In applying infinitesimal considerations we can use this structure every time the manifold M is not endowed by a fibred structure in an intrinsic manner. If Y denotes a fibred manifold over a base Z , we take $Y_s \stackrel{\text{def}}{=} J_s Y$ to mean the space of all the jets of order s of cross-sections of the surmersion $\pi : Y \rightarrow Z$. Then $\mu_j : Y_s \rightarrow J_s(\mathbb{R}^p \times \mathbb{R}^q) \approx J_s(\mathbb{R}^p, \mathbb{R}^q)$ will denote a standard s -order prolongation of the local chart μ .

1. Vector bundle valued differential systems. Let $E \rightarrow B$, $E_1 \rightarrow B_1$, $F \rightarrow X$, $F_1 \rightarrow X_1$, and $E' \rightarrow B$ be some fibre bundles, let $\Gamma(E)$ denote the set of smooth cross-

sections of the fibred manifold E and let $\mathfrak{F}(B)$ stand for a ring of smooth functions defined on the manifold B . We shall denote the tensor product of the inverse images of vector bundles with respect to some manifold morphisms $f : B \rightarrow X$ and $f_1 : B \rightarrow X_1$ by \otimes_B . If the manifold B is fibred over some manifold N , semi-basic differential forms over the manifold B with values in the vector bundle F are defined as smooth cross-sections of the vector bundle $\wedge T^*N \otimes_B \mathfrak{F}$. They constitute an $\mathfrak{F}(B)$ -submodule, denoted by $\mathfrak{F}_B(N, F)$, of the graded module $\mathfrak{F}(B, F) = \Gamma(\wedge T^*B \otimes_B F)$. Every bundle homomorphism $h : F \rightarrow F_1$ over the base homomorphism $\underline{h} : X \rightarrow X_1$ satisfying $f_1 = \underline{h} \circ f$ obviously defines the moduli homomorphism $h_\# : \mathfrak{F}(B, F) \rightarrow \mathfrak{F}(B, F_1)$. The canonical vector bundle pairing $E \times \text{Hom}(E, E') \rightarrow E'$ defines an exterior product $\wedge : \mathfrak{F}(B, E) \times \mathfrak{F}(B, E^* \otimes E') \rightarrow \mathfrak{F}(B, E')$. If we think of a cross-section $\varphi \in \mathfrak{F}^0(B, E^* \otimes E')$ as a B -morphism $\hat{\varphi} : E \rightarrow E'$, then $\hat{\varphi}_\# \omega = \omega \wedge \varphi$ for any $\omega \in \mathfrak{F}(B, E)$.

The vector bundle $\text{End}(E)$ is a bundle of algebras with the composition rule that of the superposition of endomorphisms. This fibrewise composition defines the structure of a graded non-commutative algebra in the $\mathfrak{F}(B)$ module of $\text{End}(E)$ -valued differential forms which we denote by $\mathfrak{F}(E, \text{End } E)$. The vector bundle E may be viewed as a bundle of $\text{End}(E)$ -moduli with respect to the pairing $E \times \text{End}(E) \rightarrow E$. This pairing makes the $\mathfrak{F}(B)$ -module $\mathfrak{F}(B, E)$ into a graded module over the algebra $\mathfrak{F}(B, \text{End } E)$.

Now let $h : E_1 \rightarrow E$ be a vector bundle homomorphism over a base morphism $g : B_1 \rightarrow B$ and let h_1 denote the corresponding induced B_1 -homomorphism of vector bundles $h_1 : E_1 \rightarrow g^{-1}E$. The dual vector bundle homomorphism $h_1^* : (g^{-1}E)^* \rightarrow E_1^*$ denotes, by virtue of identification $(g^{-1}E)^* \approx g^{-1}(E^*)$, some g -comorphism $h^* : g^{-1}(E^*) \rightarrow E_1^*$. If one sets $E = TB$, $E_1 = TB_1$, one gets the action $h^\#$ of h on ordinary differential forms on B . If α_g stands for the inverse image of the cross-section $\alpha \in \mathfrak{F}(B, \mathbb{R})$, then $h^\# \alpha = (\wedge h)^* \circ \alpha_g \in \mathfrak{F}(B_1, \mathbb{R})$. Given some cross-section $\omega \in \mathfrak{F}(B, E)$ and recalling the notion of its inverse image $\omega_g \in \Gamma(\wedge T^*B \otimes_{B_1} E)$, the differential form $g^* \omega \in \mathfrak{F}(B_1, E)$ is being defined by the superposition of mappings: $g^* \omega = ((\wedge Tg)^* \otimes \text{id}) \circ \omega_g$. If the morphism g has a canonical prolongation to a vector bundle g -comorphism $k : g^{-1}E \rightarrow E_1$, then the notation $g^\# \omega$ will mean the differential form $k_\# g^* \omega = ((\wedge Tg)^* \otimes k) \circ \omega$.

The functor g^* is consistent with the tensor product and contraction operations. Let throughout this paragraph the manifold B be parallelizable and the module $\Gamma(E)$ be free. Then, if $\omega = \alpha \otimes \gamma \in \mathfrak{F}(B, E) \approx \mathfrak{F}(B, \mathbb{R}) \otimes \Gamma(E)$ and if $\Omega = \beta \otimes \varphi \in \mathfrak{F}(B, E^* \otimes E') \approx \mathfrak{F}(B, \mathbb{R}) \otimes \Gamma(E^* \otimes E')$, then for the inverse images of the cross-sections ω , Ω , and $\omega \wedge \Omega = (\alpha \wedge \beta) \otimes \langle \gamma_g, \varphi \rangle$, it is true that $(\omega \wedge \Omega)_g = (\alpha_g \wedge \beta_g) \otimes \langle \gamma_g, \varphi_g \rangle \in \Gamma(\wedge T^*B \otimes_{B_1} E')$. Composing with $(\wedge Tg)^*$ one gets $g^*(\omega \wedge \Omega) = g^* \omega \wedge g^* \Omega$.

A one-form $\vartheta \in \mathfrak{F}^1(B, TB)$ may be thought of as a B -homomorphism of vector bundles $\hat{\vartheta} : TB \rightarrow TB$; then the dual homomorphism $\hat{\vartheta}^* : T^*B \rightarrow TB$ acts over the cross-sections from $\Gamma(T^*B) \equiv \mathfrak{F}^1(B, \mathbb{R})$ in an obvious manner and is being extended to $\mathfrak{F}(B, \mathbb{R})$ as a differentiation of degree 0 acting trivially on the ring of functions. We denote this extended action by $\vartheta \bar{\wedge}$. Again within local considerations for a parallelizable B and $\Gamma(E)$ being free, we can identify $\mathfrak{F}(B, \text{End } E)$ with the tensor product $\mathfrak{F}(B, \mathbb{R}) \otimes_{\mathfrak{F}(B)} \Gamma(\text{End } E)$ and define $\vartheta \bar{\wedge}(\alpha \otimes \varphi) = (\vartheta \bar{\wedge} \alpha) \otimes \text{id}(\varphi)$. Consider the ideal $\mathfrak{J} \equiv \mathfrak{J}(B) \equiv \mathfrak{J}(\vartheta) = \sum_{d>0} \mathfrak{J}^{(d)} = \vartheta \bar{\wedge} \mathfrak{F}(B, \mathbb{R})$ in the algebra $\mathfrak{F}(B, \mathbb{R})$. Consider also the $\mathfrak{F}(B, \mathbb{R})$ -submodule $\mathfrak{J}(B, \text{End } E) = \vartheta \bar{\wedge} \mathfrak{F}(B, \text{End } E)$ in the $\mathfrak{F}(B, \mathbb{R})$ -module $\mathfrak{F}(B, \text{End } E)$. We denote by $\mathfrak{J}(B, E)$ the submodule $\mathfrak{F}(B, E) \wedge \mathfrak{J}$ in the module $\mathfrak{F}(B, E)$ over the algebra $\mathfrak{F}(B, \mathbb{R})$.

It is clear that $\mathcal{J}(B, E)$ is also a submodule over the algebra $\mathfrak{F}(B, \text{End } E)$ and that $\mathcal{J}(B, E) = \mathfrak{F}(B, E) \wedge \mathcal{J}(B, \text{End } E)$. In fact, $\mathcal{J}(B, E)$ is generated over $\mathfrak{F}(B, \mathbb{R})$ by the $\mathfrak{F}^0(B, \text{End } E)$ -submodule of one-forms

$$\mathfrak{P}(B, E) = \vartheta \wedge \mathfrak{F}^1(B, E) = \mathfrak{F}^0(B, E) \wedge \mathfrak{P},$$

where $\mathfrak{P} \equiv \mathfrak{P}(B) \equiv \mathcal{J}^{(1)} = \vartheta \wedge \mathfrak{F}^1(B, \mathbb{R})$ is an $\mathfrak{F}(B)$ -submodule, that is the ordinary Pfaff system.

Definition 1 *An exterior differential system $\mathfrak{S} \equiv \mathfrak{S}(F)$ (with values in the vector bundle F) is an $\mathfrak{F}(B, \text{End } F)$ -submodule of the module $\mathfrak{F}(B, F)$. A Pfaff system is an exterior differential system generated by one-forms. Let a manifold Z be given. We call the germ of an immersion $\sigma : Z \rightarrow B$ a solution of the exterior differential system \mathfrak{S} if $\sigma^* \mathfrak{S} = 0$. Two exterior differential systems (in general, over different manifolds) are called equivalent if the sets of their solutions coincide (in general, are isomorphic).*

Let over a manifold B be given an exterior differential system \mathfrak{S} with values in the vector bundle E and an exterior differential system \mathfrak{S}' with values in the vector bundle E' and let $\text{sol}(\mathfrak{S})$ denote the set of solutions of the system \mathfrak{S} . If $\mathfrak{S} \wedge \mathfrak{F}(B, E^* \otimes E') \supset \mathfrak{S}'$, then $\text{sol}(\mathfrak{S}) \subset \text{sol}(\mathfrak{S}')$.

Definition 2 *Exterior differential systems \mathfrak{S} and \mathfrak{S}' are algebraically equivalent if at the same time both $\mathfrak{S} \wedge \mathfrak{F}(B, E^* \otimes E') \supset \mathfrak{S}'$ and $\mathfrak{S}' \wedge \mathfrak{F}(B, E'^* \otimes E) \supset \mathfrak{S}$.*

Lemma 1 *Two Pfaff systems, which are algebraically equivalent, are equivalent.*

This follows from the completeness of Pfaff systems.

Let $\tilde{\mathfrak{S}}$ denote the image of the exterior differential system \mathfrak{S} under the projection $j : \mathfrak{F}(B, E) \rightarrow \mathfrak{F}(B, E)/\mathcal{J}(B, E)$. In view of the inclusion $\mathcal{J}(B, E) \wedge \mathfrak{F}(B, E^* \otimes E') \subset \mathcal{J}(B, E')$, the action \wedge of elements from $\mathfrak{F}(B, E^* \otimes E')$ over the quotient module is defined and if for this action $\mathfrak{S} \wedge \mathfrak{F}(B, E^* \otimes E') \supset \mathfrak{S}'$, then also $\tilde{\mathfrak{S}} \wedge \mathfrak{F}(B, E^* \otimes E') \supset \tilde{\mathfrak{S}'}$.

Definition 3 *A solution of the system $(\mathfrak{S}, \vartheta)$ is a germ of an immersion $\sigma : Z \rightarrow B$, such that $\sigma^* j^{-1}(\tilde{\mathfrak{S}}) = 0$.*

Let over the manifold B_1 be specified a differential form $\vartheta_1 \in \mathfrak{F}(B_1, TB_1)$ such that $g^* \vartheta = \vartheta_1 \bar{\wedge} \varrho \in \mathfrak{P}(B_1, g^{-1}TB)$ for some $\varrho \in \mathfrak{F}(B_1, TB)$. Taking into account the definitions of the actions g^* , $\vartheta_1 \bar{\wedge}$ and identifying the differential forms ϑ , ϑ_1 , and ϱ with the corresponding vector bundle homomorphisms $\hat{\theta}$, $\hat{\theta}_1$ and $\hat{\rho} : TB_1 \rightarrow g^*TB$, we have $\hat{\theta}_g \circ (Tg)_1 = \hat{\rho} \circ \hat{\theta}_1$, where $\hat{\theta}_g$ is the inverse image of the homomorphism $\hat{\theta}$ with respect to the morphism g . For conjugated homomorphisms we also have $(Tg)^* \circ \hat{\theta}_g^* = \hat{\theta}_1^* \circ \hat{\rho}^*$.

In view of the definition of the action $\hat{\rho}^\#$ on the differential form $\alpha \in \mathfrak{F}^1(B, \mathbb{R})$, it holds that

$$g^*(\vartheta \bar{\wedge} \alpha) = (Tg)^* \circ \hat{\theta}_g^* \circ \alpha_g = \hat{\theta}_1^* \circ \hat{\rho}^* \circ \alpha_g = \vartheta_1 \bar{\wedge} \hat{\rho}^* \alpha \in \mathfrak{P}_1 \equiv \mathfrak{P}(B_1).$$

It follows that $g^* \mathfrak{P} \subset \mathfrak{P}_1$, thus $g^* \mathcal{J} \subset \mathcal{J}_1$ and also

$$g^* \mathcal{J}(B, E) = g^* \mathfrak{F}(B, E) \wedge g^* \mathcal{J} \subset \mathfrak{F}(B_1, g^{-1}E) \wedge \mathcal{J}(B_1, g^{-1}E),$$

so the action g^* of the homomorphism g over the quotient module $\mathfrak{F}(B, E/\mathcal{J}(B, E))$ is defined.

2. The Lie derivative. Let $w : B \rightarrow W$ be a morphism of manifolds and let w_t denote its deformation. Then $\dot{w}_t(0) : B \rightarrow TW$ is a lift of the morphism w . If the manifold W is fibred over the manifold B and if every w_t is a cross-section, then $\dot{w}_t(0)$ belongs to the vertical tangent bundle $VW \subset TW$. Let ξ and $\bar{\eta}$ be some vector fields on the manifolds B and W , respectively. The Lie derivative [3] of the morphism w with respect to the pair of vector fields $\xi, \bar{\eta}$ is defined by the expression $\mathbf{L}_{\xi, \bar{\eta}} w = (Tw) \circ \xi - \bar{\eta} \circ w$. If w is a cross-section and if the vector field $\bar{\eta}$ is projectible onto the vector field ξ , then the Lie derivative $\mathbf{L}_{\xi, \bar{\eta}} w$ belongs to VW . Let under these assertions $\exp_B t\bar{\eta}$ denote a one-parameter induced family of diffeomorphism over B , corresponding to $\exp t\bar{\eta}$ and let

$$w_t = \exp_B(-t\bar{\eta}) \circ w_{\exp t\xi} = \exp(-t\bar{\eta}) \circ w \circ \exp t\xi.$$

Differentiating at $t = 0$ we obtain

$$\dot{w}_t(0) = -(\exp t\bar{\eta}) \circ w + (Tw) \circ (\exp t\xi)' = \mathbf{L}_{\xi, \bar{\eta}} w.$$

If the fibre bundle W is a vector bundle, then $VW \approx W \times_B W$, the first component of $\dot{w}_t(0)$ coincides with the initial cross-section w , and the second one is identified with the Lie derivative $\mathbf{L}_{\xi, \bar{\eta}} w$ itself. In a special case setting $W = \wedge T^*B \otimes E$ and taking $\bar{\eta}$ to be build up from the standard lift to the cotangent bundle of the vector field ξ together with some vector field η on E , projectible into ξ , we agree with truncated notation of $\mathbf{L}_{\xi, \eta} \stackrel{\text{def}}{=} \mathbf{L}_{\xi, \bar{\eta}}$.

Definition 4 A diffeomorphism $g : B \rightarrow B$ is called a symmetry (vs an algebraic symmetry) of the exterior differential system $\mathfrak{S}(E)$ if the exterior differential systems $g^*\mathfrak{S}$ and \mathfrak{S} itself are equivalent (vs algebraically equivalent). A vector field ξ on the manifold B is called an infinitesimal symmetry (vs an algebraic infinitesimal symmetry) of the exterior differential system \mathfrak{S} if for all t the transformation $\exp t\xi$ is a symmetry (vs an algebraic symmetry) of the system \mathfrak{S} .

That ξ is an algebraic infinitesimal symmetry of the exterior differential system \mathfrak{S} means that

$$(\exp t\xi)^*\mathfrak{S} \subset \mathfrak{S} \wedge \mathfrak{F}(B, E^* \otimes (\exp t\xi)^{-1}E).$$

In particular, whatever the vector field $\eta \in \Gamma(TE)$, projectible into the vector field ξ be, the inclusions

$$\begin{aligned} (\exp -t\eta)_{B\#} (\exp t\xi)^{-1}\mathfrak{S} \subset \mathfrak{S} \wedge \mathfrak{F}(B, E^* \otimes (\exp t\xi)^{-1}E \wedge \mathfrak{F}^0(B, (\exp t\xi)^{-1}\star(E^* \otimes E)) \\ \subset \mathfrak{S} \wedge \mathfrak{F}(B, E^* \otimes E) = \mathfrak{S} \end{aligned}$$

hold. By differentiating with respect to the parameter t , we conclude that if ξ is the infinitesimal symmetry of the system \mathfrak{S} , then

$$\mathbf{L}_{\xi, \eta}\mathfrak{S} = \frac{d}{dt}(\exp t\xi)^*\mathfrak{S}(0) \subset \mathfrak{S}$$

for every vector field η , projectible into the field ξ . In the case of a trivial vector bundle E , one can put $\eta = (\xi, 0)$ and use the notation $\mathbf{L}_{\xi, 0}$ in place of $\mathbf{L}_{\xi, (\xi, 0)}$.

Definition 5 A diffeomorphism $g : B \rightarrow B$ is called a symmetry (vs algebraic symmetry) of the system $(\mathfrak{S}, \vartheta)$, if $g^*\mathfrak{J} \subset \mathfrak{J}$ and if the exterior differential systems $g^*\mathfrak{S}$ and \mathfrak{S} itself are equivalent (vs algebraically equivalent).

An algebraic symmetry of the system $(\mathfrak{S}, \vartheta)$ is a symmetry (vs an algebraic symmetry) of the exterior differential system $\mathfrak{S} + \mathfrak{J}(B, E)$ and *vice versa*. In that case $\mathbf{L}_\xi \mathfrak{J} \subset \mathfrak{J}$ and $\mathbf{L}_{\xi, \eta} \tilde{\mathfrak{S}} \subset \tilde{\mathfrak{S}}$, that is

$$\mathbf{L}_{\xi, \eta} \mathfrak{S} \subset \mathfrak{S} + \mathfrak{J}(B, E). \quad (1)$$

3. Symmetries of the Euler-Lagrange equations. On the manifold Y_{s+1} of jets of order $s+1$ of cross-sections of the fibred manifold $\pi : Y \rightarrow Z$ there exists a canonical contact differential form $\vartheta_s \in \mathfrak{F}_{s+1}(Y_s, V_s)$, semi-basic with respect to the projection $\pi_s^{s+1} : Y_{s+1} \rightarrow Y_s$ which takes values in the vector bundle V_s of vertical tangent vectors to the surmersion $\pi^s : Y_s \rightarrow Z$ and such that the Pfaff system $\mathfrak{P}_s = \vartheta_s \wedge \mathfrak{F}_{s+1}^1(Y_s, \mathbb{R})$ is nothing but the Cartan co-distribution. (We use a subscript s instead of Y_s where possible). For an open set $U_s \subset Y_s$ and applying the projection $\pi_v^s : Y_s \rightarrow Y_v$ let us put $\mathfrak{J}_v(U_s) = \mathfrak{J}(\pi_v^s|_{U_s}^* \vartheta_v)$ and let $v_j : Y_s \rightarrow J(\mathbb{R}^p, \mathbb{R}^q)$ be the canonical lift of a local fibred chart v on the manifold Y . A Lagrangian is an element $\tilde{\lambda}$ of the quotient sheaf, generated by the quotient moduli $\mathfrak{F}_{U_r}^p(Y, \mathbb{R})/\mathfrak{I}_0^{(p)}(U_r)$, which we shall briefly denote by $\mathcal{F}_r^p(Y, \mathbb{R})/\mathcal{I}_0^{(p)}(Y_r)$. (We denote the corresponding sheaves by calligraphic characters). There can always be found, at least locally, a representative λ , semi-basic with respect to the projection $\pi^r : Y_r \rightarrow Z$, that is $\lambda \in \mathcal{F}_r^p(Z, \mathbb{R})$. The Euler-Lagrange expressions which correspond to the Lagrange density $v_j^{-1*} \lambda \in \mathcal{F}_r^p(\mathbb{R}^p, \mathbb{R})$ naturally take the shape of the components of the local expression of some differential p -form

$$\varepsilon \in \mathcal{F}_{2r}^p(Z, V_0^*) \quad (2)$$

corresponding to the Lagrangian λ [4].

The symmetries of the Euler-Lagrange equations are nothing else but the symmetries of the system $(\mathfrak{S}_\varepsilon, \vartheta_{2r-1})$, where the module \mathfrak{S}_ε is generated by the vector bundle valued differential form ε . From here on we shall deal with variational calculus in one independent variable. As an example, we consider a system of third-order ordinary differential equations

$$\mathbf{E}_a = 0. \quad (3)$$

Let $Y = \mathbb{R}^p \times \mathbb{R}^q$, and let the canonical coordinates in the manifold $J_r(\mathbb{R}^1; \mathbb{R}^q)$ be denoted by $t, \mathbf{x} = (x^a), \mathbf{v} = (v^a), \mathbf{v}' = (v'^a), \dots, \mathbf{v}^{(r-1)} = (v^{(r-1)a})$.

This corresponds to $p = 1, k = q$ in (2). We introduce a vector-valued differential one-form

$$\varepsilon = \mathbf{E}_a \mathbf{d}x^a \otimes \mathbf{d}t, \quad (4)$$

where the expressions \mathbf{E}_a are Euler-Poisson ones. Applying the general criterion of [5] for an arbitrary system of differential equations to be a system of Euler-Poisson equations, it was established in [6] that the vector expression $\mathbf{E} = \{\mathbf{E}_a\}$ in (3) must take the shape of

$$\mathbf{E} = \mathbb{A} \cdot \mathbf{v}'' + (\mathbf{v}' \cdot \partial_v) \mathbb{A} \cdot \mathbf{v}' + \mathbb{B} \cdot \mathbf{v}' + \mathbf{c}, \quad (5)$$

where the skew-symmetric matrix \mathbb{A} , the matrix \mathbb{B} , and the column vector \mathbf{c} depend on the variables t , \mathbf{x} , $\mathbf{v} = d\mathbf{x}/dt$, and satisfy the following system of partial differential equations in t , x^a , and v^a [6, 7]

$$\begin{aligned}
& \partial_{\mathbf{v}[a} \mathbb{A}_{bc]} = 0 \\
& 2\mathbb{B}_{[ab]} - 3\mathbf{D}_1 \mathbb{A}_{ab} = 0 \\
& 2\partial_{\mathbf{v}[a} \mathbb{B}_{b]c} - 4\partial_{\mathbf{x}[a} \mathbb{A}_{b]c} + \partial_{\mathbf{x}^c} \mathbb{A}_{ab} + 2\mathbf{D}_1 \partial_{\mathbf{v}^c} \mathbb{A}_{ab} = 0 \\
& \partial_{\mathbf{v}(a} \mathbf{c}_{b)} - \mathbf{D}_1 \mathbb{B}_{(ab)} = 0 \\
& 2\partial_{\mathbf{v}^c} \partial_{\mathbf{v}[a} \mathbf{c}_{b]} - 4\partial_{\mathbf{x}[a} \mathbb{B}_{b]c} + \mathbf{D}_1^2 \partial_{\mathbf{v}^c} \mathbb{A}_{ab} + 6\mathbf{D}_1 \partial_{\mathbf{x}[a} \mathbb{A}_{bc]} = 0 \\
& 4\partial_{\mathbf{x}[a} \mathbf{c}_{b]} - 2\mathbf{D}_1 \partial_{\mathbf{v}[a} \mathbf{c}_{b]} - \mathbf{D}_1^3 \mathbb{A}_{ab} = 0.
\end{aligned} \tag{6}$$

In (6) \mathbf{D}_s denotes generators of the Cartan distribution,

$$\mathbf{D}_2 = \mathbf{v}' \cdot \partial_{\mathbf{v}} + \mathbf{D}_1, \quad \mathbf{D}_1 = \partial_t + \mathbf{v} \cdot \partial_{\mathbf{x}}.$$

Let ϑ_2 , ϑ_3 denote the canonical contact forms

$$\begin{aligned}
\vartheta_3 &= \frac{\partial}{\partial v'^a} \otimes (dv'^a - v''^a dt) + \vartheta_2, \\
\vartheta_2 &= \frac{\partial}{\partial v^a} \otimes (dv^a - v'^a dt) + \frac{\partial}{\partial x^a} \otimes (dx^a - v^a dt).
\end{aligned}$$

Along with the differential form ϵ , we introduce another one,

$$\underline{\epsilon} = A_{ab} dx^a \otimes dv'^b + K_a dx^a \otimes dt, \quad \mathbf{K} = (\mathbf{v}' \cdot \partial_{\mathbf{v}}) \mathbb{A} \cdot \mathbf{v}' + \mathbb{B} \cdot \mathbf{v}' + \mathbf{c}.$$

Exterior differential systems, generated by the forms ϵ and $\underline{\epsilon}$, are equivalent:

$$\underline{\epsilon} - \epsilon = \vartheta_3 \bar{\wedge} (A_{ab} dx^a \otimes dv'^b). \tag{7}$$

Now it is time to put in the concept of *symmetry*. Let

$$\mathbf{r} = \mathbf{t} \frac{\partial}{\partial t} + \mathbf{r}^a \frac{\partial}{\partial x^a} \tag{8}$$

denote a generator of some local group of transformations of the manifold $\mathbb{R} \times \mathbb{R}^q$, its successive prolongations to the space $J_s(\mathbb{R}; \mathbb{R}^q)$ denoted by \mathbf{r}_s ,

$$\mathbf{r}_2 = v^a \frac{\partial}{\partial v'^a} + \mathbf{r}_1.$$

The demand that the exterior differential system, generated by the vector-valued differential form $\underline{\epsilon}$, be invariant under the infinitesimal transformation \mathbf{r} incarnates in consistency with (1) into the following equation

$$\mathbf{L}(\mathbf{r}_2)(\underline{\epsilon}) = \underline{\Xi} \cdot \underline{\epsilon} + \vartheta_2 \bar{\wedge} \omega, \tag{9}$$

where the elements of a matrix $\underline{\Xi} \in \mathfrak{F}^0(J_2(\mathbb{R}; \mathbb{R}^q); \text{Hom}(\mathbb{R}^{q*}, \mathbb{R}^{q*}))$ and the coefficients of a π_1^2 -horizontal \mathbb{R}^{q*} -valued one-form $\omega \in \mathfrak{F}^1(J_1(\mathbb{R}; \mathbb{R}^q); \mathbb{R}^{q*})$ depend upon the variables t , \mathbf{x} , \mathbf{v} , and \mathbf{v}' . Both $\underline{\Xi}$ and ω play the role of Lagrange multipliers. Splitting equation (9) with

respect to independent differentials $\mathbf{d}t$, $\mathbf{d}x^a$, $\mathbf{d}v^a$, and $\mathbf{d}v'^a$, gives the following system of partial differential equations

$$\mathbf{L}(\mathbf{x}_1)A_{ab} = \Xi_a^c A_{cb} - A_{ac} \frac{\partial}{\partial v'^b} v^c \quad (10)$$

$$\mathbf{L}(\mathbf{x}_2)K_a = \Xi_a^b K_b - A_{ab} \mathbf{D}_2 v^b - K_a \mathbf{D}_1 t. \quad (11)$$

Variational problems in a parametric form. A variational problem in a parametric form is one posed on the manifold $J_r(Z, M)$, so we have to put $Y = Z \times M$, the dimension of the manifold M equals to $p + q$ with $p = 1$, and $k = p + q = 1 + q$ in (2). Canonical coordinates in the manifold $J_r(\mathbb{R}; \mathbb{R}^{1+q})$ are denoted by τ , $\mathbf{x} = (x^\rho)$, $\mathbf{u} = (u^\rho)$, $\dot{\mathbf{u}} = (\dot{u}^\rho)$, \dots , $\overset{r-1}{\mathbf{u}} = (\overset{r-1}{u}^\rho)$. The manifold of r^{th} -order velocities, $T_r^1 M$, is defined as $T_r^1 M = J_r(\mathbb{R}; M)_0$. There exists an obvious isomorphism $J_r(\mathbb{R}; M) \approx \mathbb{R} \times T_r^1 M$. Coordinates in the manifold $T_r^1 M$ are denoted by $\mathbf{x}, \mathbf{u}, \dot{\mathbf{u}}, \dots, \overset{r-1}{\mathbf{u}}$. If M is a (pseudo-)Euclidean n -dimensional space (of an arbitrary signature), the Hodge operator “ $*$ ” is defined as $(*w)_{\rho_{k+1}\dots\rho_n} = \frac{1}{(n-k)!} e_{\rho_1\dots\rho_n} w^{\rho_1\dots\rho_k}$.

Consider now a variational problem in a parametric form, set by a Lagrangian

$$\ell(\tau, x^\rho, u^\rho, \dots, \overset{r-1}{u}^\rho) \mathbf{d}\tau$$

on the space $J_r(\mathbb{R}; M)$. As long as we constrain ourselves only to the case of autonomous Euler-Poisson equations,

$$\mathcal{E}_\rho = 0, \quad (12)$$

the differential form

$$\varepsilon = \mathcal{E}_\rho \mathbf{d}x^\rho \otimes \mathbf{d}\tau, \quad (13)$$

may *globally* be deprived of the factor $\mathbf{d}\tau$, constituting thus a *globally* defined T^*M -valued density

$$\mathbf{e} = \mathcal{E}_\rho \mathbf{d}x^\rho. \quad (14)$$

Let $C_r^1 M$ denote the manifold of the r^{th} order contact germs of one-dimensional submanifolds of M . The projection $\wp : T_r^1 M \setminus 0 \rightarrow C_r^1 M$ can be employed to generate an autonomous variational problem set over $T_r^1 M$ from every one variational problem over $C_r^1 M$.

Lemma 2 ([2]) *In terms of a local chart, if in (4) the local semi-basic differential form ε corresponds to the Lagrangian*

$$\lambda = L \mathbf{d}t,$$

then the vector-valued density

$$\mathbf{e} = x^a (\mathbf{E}_a \circ \wp) \mathbf{d}t + t (\mathbf{E}_a \circ \wp) \mathbf{d}x^a \quad (15)$$

corresponds to the Lagrangian

$$\ell(\tau, x^\rho, u^\rho, \dots, \overset{r-1}{u}^\rho) \mathbf{d}\tau = \mathcal{L}(x^\rho, u^\rho, \dots, \overset{r-1}{u}^\rho) \mathbf{d}\tau$$

with the Lagrange function

$$\mathcal{L} = u^0 L \circ \wp. \quad (16)$$

4. The classical spinning particle. Let in (8) the generator \mathfrak{r} correspond to (pseudo-)orthogonal transformations of a four-dimensional (pseudo-)Euclidean space. In this case, equations (10) have no solutions of the third order, that is to say, no invariant system of Euler-Poisson equations (3) consists of at least one equation of the third order.

Nevertheless, if we allow a vector parameter $\mathbf{s} = (s^0, \mathbf{s})$ (that is, a constant quantity transforming as a four-vector under the action of the (pseudo-)orthogonal group) to enter expressions (5), then equations (6&10) turn out to possess a family of solutions, which depend on \mathbf{s} and actually contain some third-order derivatives. We succeeded in obtaining the following one (under the unessential assumption $g^{00} = +1$)

$$\begin{aligned} \mathbf{E} = & \frac{\mathbf{v}'' \times (\mathbf{s} - s_0 \mathbf{v})}{[(1 + \mathbf{v} \cdot \mathbf{v})(s_0^2 + \mathbf{s} \cdot \mathbf{s}) - (s_0 + \mathbf{s} \cdot \mathbf{v})^2]^{3/2}} - \\ & 3 \frac{(s_0^2 + \mathbf{s} \cdot \mathbf{s}) \mathbf{v}' \cdot \mathbf{v} - (s_0 + \mathbf{s} \cdot \mathbf{v}) \mathbf{s} \cdot \mathbf{v}'}{[(1 + \mathbf{v} \cdot \mathbf{v})(s_0^2 + \mathbf{s} \cdot \mathbf{s}) - (s_0 + \mathbf{s} \cdot \mathbf{v})^2]^{5/2}} \mathbf{v}' \times (\mathbf{s} - s_0 \mathbf{v}) + \\ & m \frac{(1 + \mathbf{v} \cdot \mathbf{v}) \mathbf{v}' - (\mathbf{v}' \cdot \mathbf{v}) \mathbf{v}}{(1 + \mathbf{v} \cdot \mathbf{v})^{3/2} (s_0^2 + \mathbf{s} \cdot \mathbf{s})^{3/2}}. \end{aligned} \quad (17)$$

Applying the prescription (15), the above expression (17) produces the corresponding four-dimensional expression and thus the desired Euler-Poisson equations

$$\mathcal{E} = \frac{* \dot{\mathbf{u}} \wedge \mathbf{u} \wedge \mathbf{s}}{\|\mathbf{s} \wedge \mathbf{u}\|^3} - 3 \frac{* \dot{\mathbf{u}} \wedge \mathbf{u} \wedge \mathbf{s}}{\|\mathbf{s} \wedge \mathbf{u}\|^5} (\dot{\mathbf{u}} \wedge \mathbf{s}) \cdot (\mathbf{u} \wedge \mathbf{s}) + \frac{m}{\|\mathbf{s}\|^3} \left[\frac{\dot{\mathbf{u}}}{\|\mathbf{u}\|} - \frac{\dot{\mathbf{u}} \cdot \mathbf{u}}{\|\mathbf{u}\|^3} \mathbf{u} \right] = 0. \quad (18)$$

Assume \mathfrak{r}_3 denote a third-order prolongation of the generator \mathfrak{r} (8) of (pseudo-)orthogonal transformations with the group parameters \mathbf{n}, \mathbf{q} to the space $J_3(\mathbb{R}; \mathbb{R}^3)$,

$$\begin{aligned} \mathfrak{r}_3 = & -(\mathbf{q} \cdot \mathbf{s}) \partial_{s_0} + s^0 (\mathbf{q} \cdot \partial_{\mathbf{s}}) + [\mathbf{n}, \mathbf{s}, \partial_{\mathbf{s}}] - (\mathbf{q} \cdot \mathbf{x}) \partial_t + t (\mathbf{q} \cdot \partial_{\mathbf{x}}) + [\mathbf{n}, \mathbf{x}, \partial_{\mathbf{x}}] + \\ & (\mathbf{q} \cdot \partial_{\mathbf{v}}) + (\mathbf{q} \cdot \mathbf{v}) (\mathbf{v} \cdot \partial_{\mathbf{v}}) + [\mathbf{n}, \mathbf{v}, \partial_{\mathbf{v}}] + \\ & 2(\mathbf{q} \cdot \mathbf{v}) (\mathbf{v}' \cdot \partial_{\mathbf{v}'}) + (\mathbf{q} \cdot \mathbf{v}') (\mathbf{v} \cdot \partial_{\mathbf{v}'}) + [\mathbf{n}, \mathbf{v}', \partial_{\mathbf{v}'}]. \end{aligned} \quad (19)$$

The following assertions are true:

1. Vector quantity \mathcal{E} in (18) constitutes a system of Euler-Poisson expressions;
2. Let vector field \mathfrak{r}_3 be given by (19), and let \mathbf{E} be given by (17) (recall that (17) is nothing more but merely (18) parametrized with respect to the coordinate time $t = x^0$). The relation

$$\mathbf{L}(\mathfrak{r}_3)(\mathbf{E}) = \mathbf{n} \times \mathbf{E} + (\mathbf{q} \cdot \mathbf{v}) \mathbf{E} - (\mathbf{v} \cdot \mathbf{E}) \mathbf{q}$$

proves the invariance of the Euler-Poisson equations (18), and at the same time shows that whatever a possible Lagrange function for (18) might exist, it by no means will reveal invariance under (pseudo-)orthogonal transformations, even in the generalized sense (that is up to a total derivative term).

3. Equations (18) describe (in metric signature 2) the motion of a relativistic spinning free particle with the constant spin four-vector \mathbf{s} and rest mass

$$m_0 = m \left[1 - \frac{(\mathbf{s} \cdot \mathbf{u})^2}{(\mathbf{s} \cdot \mathbf{s})(\mathbf{u} \cdot \mathbf{u})} \right]^{3/2}.$$

In [7, 8] and [6] more details about the classical spinning particle equations are supplied.

A comment on the order of the Euler-Lagrange form. Euler-Lagrange equations are polynomial with respect to the derivatives of orders greater than the maximal order of the derivatives which enter the corresponding Lagrange function. It was this property that inspired us to try to diminish the order of the underlying manifold $J_s(\mathbb{R}, \mathbb{R}^q)$ from $s = 3$ to $s = 2$ by means of introducing the differential form $\underline{\epsilon}$ in (7) rather than ϵ .

References

- [1] Ovsyannikov L.V., Group Analysis of Differential Systems, Orlando, Academic Press, 1982, 416p.
- [2] Matsyuk R.Ya., Exterior differential equations in generalized mechanics and symmetry properties, In: Methods for Studying Differential and Integral Operators, Kyiv, Naukova Dumka, 1989, 153–160 (In Russian).
- [3] Kolář I., Lie derivatives and higher-order Lagrangians, In: Proceedings of the Conference (ČSSR–GDR–Poland) on Differential Geometry and its Applications, September 1980, Praha, Univ. Karlova, 1981, 117–123.
- [4] Matsyuk R.Ya., Geometrical meaning of the integration by parts, In: Lobachevskij and the Contemporary Geometry, The International Scientific Conference, Kazan', August 18–22, Thesis, Part I, Kazan': Kazan' State Univ., 1992, 60–61 (In Russian).
- [5] Tulczyjew W.M., The Euler-Lagrange resolution, In bk. Differential Geometrical Methods in Mathematical Physics, Proc. conf., 1979 [Lecture Notes in Mathematics, 836], Berlin e.a: Springer-Verlag, 1980, 22–48.
- [6] Matsyuk R.Ya., Poincaré-invariant Equations of Motion in Lagrangian Mechanics With Higher Derivatives, Thesis, Institute for Applied Problems in Mechanics and Mathematics, Academy of Science, Ukraine, L'viv, 1984, 140p. (In Russian).
- [7] Matsyuk R.Ya., Lagrangian analysis of the third-order invariant equations of motion in the relativistic mechanics of classical particles, *Sov. Phys.-Dokl.*, 1985, V.30, N 6, 458–460.
- [8] Matsyuk R.Ya., Lagrangian approach to spinning or radiating particle higher-order equations of motion in special relativity, In: 11th International Conference on General Relativity and Gravitation. Stockholm, Sweden, July 6–12, 1986. Abstracts of contributed papers, V.II, Stockholm, 1986, p.648.